

# PHOTOMETRY OF VARIABLE STARS FROM THE THU-NAOC TRANSIENT SURVEY. I. THE FIRST TWO YEARS

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## ABSTRACT

In this paper, we report the detections of stellar variabilities from the first two years of observations of a sky area of about 1300 square degrees from the Tsinghua University–NAOC Transient Survey. A total of 1237 variable stars (including 299 new ones) were detected with a brightness  $<18.0$  mag and a magnitude variation  $\gtrsim 0.1$  mag on a timescale from a few hours to a few hundred days. Among such detections, we tentatively identified 661 RR Lyrae stars, 431 binaries, 72 semi-regular pulsators, 29 Mira stars, 11 slow irregular variables, 11 RS Canum Venaticorum stars, 7 Gamma Doradus stars, 5 long-period variables, 3 W Virginis stars, 3 Delta Scuti stars, 2 Anomalous Cepheids, 1 Cepheid, and 1 nova-like star based on their time-series variability index  $J_s$  and their phased diagrams. Moreover, we found that 14 RR Lyrae stars show the Blazhko effect and 67 contact eclipsing binaries exhibit the O'Connell effect. Since the period and amplitude of light variations of RR Lyrae variables depend on their chemical compositions, their photometric observations can be used to investigate the distribution of metallicity along the direction perpendicular to the Galactic disk. We find that the metallicity of RR Lyrae stars shows large scatter at regions closer to the Galactic plane (e.g.,  $-3.0 < [\text{Fe}/\text{H}] < 0$ ) but tends to converge at  $[\text{Fe}/\text{H}] \sim -1.7$  at larger Galactic latitudes. This variation may be related to the fact that the RRAB Lyrae stars in the Galactic halo come from globular clusters with different metallicities and vertical distances, i.e., OoI and OoII populations, favoring the dual-halo model.

*Key words:* binaries: eclipsing – stars: evolution – stars: variables: general – stars: variables: RR Lyrae

*Supporting material:* machine-readable and VO table

## 1. INTRODUCTION

Time-domain astronomy (TDA) has emerged as a key field of current astronomy and astrophysics (Grindlay et al. 2012). This offers new routes to an astrophysical understanding of stellar and galactic activities such as variable stars, novae, supernovae, gamma-ray bursts, active galactic nuclei, quasars, etc. The Tsinghua University–NAOC Transient Survey (TNTS) is a four-year project designed to search for such transients (Zhang et al. 2015). The observations were carried out in a clear filter with a 0.6 m Schmidt telescope located at Xinglong station of the NAOC. Under a typical weather condition at Xinglong station (e.g., with a seeing of about 2 arcsec), this telescope and the CCD system can reach a detection limit of about 19.5 mag ( $3\sigma$ ) in the clear filter for an exposure of 60 s. The relatively short cadence (e.g., 3–4 days) and broad synoptic coverage of the sky areas provide an excellent chance to search for and study variable stars of different timescales. In this paper, we present the result of searching for variable stars from these survey data.

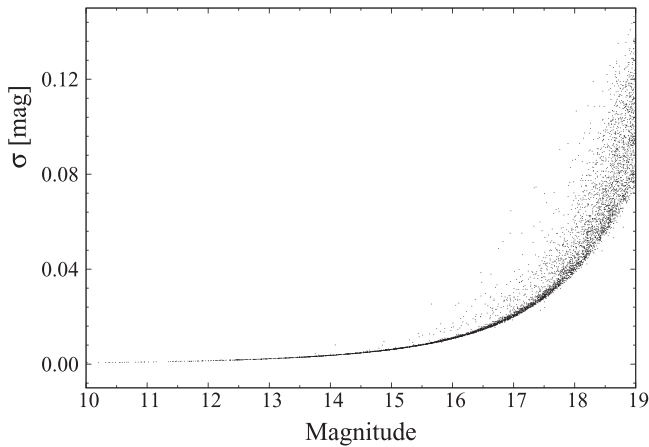
This work presents the result of searching for variable stars from the TNTS data accumulated during the period from 2012 October to 2014 August. The observations and data reductions are described in Section 2. The time-series photometry from TNTS images is presented in Section 3. The classifications of variable stars and analysis are given in Section 4. We summarize in Section 5.

## 2. OBSERVATIONS AND DATA REDUCTIONS

The TNTS is conducted with a 0.6 m Schmidt telescope equipped with a  $4k \times 4k$  CCD. The CCD has a pixel scale of  $1.3 \text{ arcsec pixel}^{-1}$  and a field of view (FOV) of  $1^\circ 5' \times 1^\circ 5'$ . It was designed to operate for four years starting from 2012 October. This survey covers a sky area of  $\sim 1800$  square degrees with a Galactic latitude  $|b| > 10^\circ$  and  $0^\circ < \delta < 60^\circ$ . It usually takes about 2 minutes to get an image for a specific sky zone, including a 60 s exposure for the selected target, 22 s readout time, and 30 s of movement and stabilization for the telescope. In order to efficiently rule out cosmic rays and moving objects, we took two exposures for each sky zone with a temporal interval of 1–1.5 hr and the whole survey area was visited every 3–4 days (see details in Zhang et al. 2015).

About 85,000 images were accumulated from our survey performed during the period from 2012 October to 2014 August, and these images cover about 1800 square degrees of the northern sky. In this study, however, we only chose fields with repeated observations  $>40$  times for detecting variable stars. This selection criterion results in observations of a sky area of  $\sim 1300$  square degrees used for analysis. The TNTS images were reduced using the IRAF<sup>7</sup> standard procedure,

<sup>7</sup> IRAF, the Image Reduction and Analysis Facility, is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy (AURA), Inc. under cooperative agreement with the National Science Foundation (NSF). <http://iraf.noao.edu/>



**Figure 1.** Photometric error vs. magnitude in a clear band for a typical image of the field SE65 from the TNTS.

including the corrections for bias, flat field, and removal of cosmic rays. The positions of those images were then registered to each other using the Guide Star Catalog (GSC; Jenkner et al. 1990), with an accuracy of 0.5 arcsec (Zhou et al. 2003). The SExtractor (Bertin & Arnouts 1996) was used to perform aperture photometry on the chosen images. The radius of aperture photometry was set as 5 pixels (6.8 arcsec), with a sky annulus size of 24 pixels (31.2 arcsec).

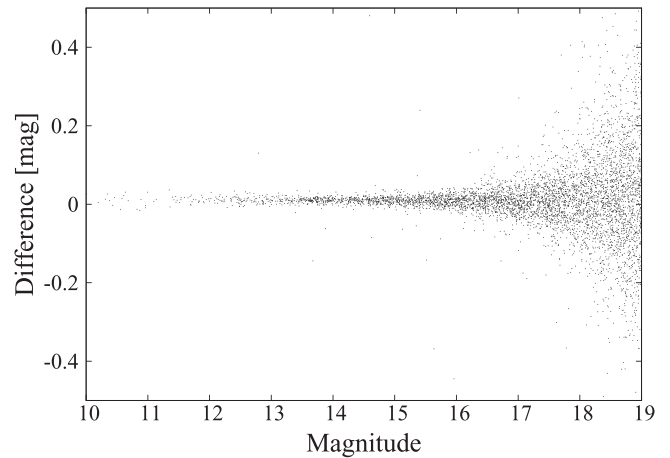
Since the stars on the TNTS images are not very crowded, we selected the parameter “BEST OF MAG” of SExtractor as output results. The detection threshold was set as  $2\sigma$  above the sky background. To alleviate the effect of nonuniform background on the photometry, e.g., moonlight and clouds, we set the BACKGROUND TYPE as LOCAL. A range of 5000–17,000 stars can be detected on the TNTS images. In our study we take the typical sky field SE65 as an example to illustrate the photometric error. For stars brighter than 17.0 mag, the photometric error is generally less than 0.02 mag according to the output result from SExtractor (see Figure 1).

### 3. TIME-SERIES PHOTOMETRY

To get the time-series photometry of each star, we feed the matching tool SRPMatch with photometry catalogs produced by SExtractor for two images of the sky field. SRPMatch is a command line tool embedded in the Swift Reduction Package (hereafter SRP<sup>8</sup>), which can be used to match the photometry catalog. We set the match tolerance as 1.5 arcsec ( $\sim 1$  pixel). For each field, one is the reference image with the maximum number of detected stars and the other is a random image from the same field. We calibrated our time-series photometry using the following two steps: magnitude calibration and rescaling of the photometric errors.

#### 3.1. Magnitude Calibration

To calibrate magnitudes of images taken at different epochs, it is necessary to calculate the magnitude offset of each image relative to the reference image. Figure 2 shows the magnitude offsets measured for the sky field SE65 (which is the same as that used in Figure 1) and its reference frame. One can see that the photometric error progressively increases for stars fainter



**Figure 2.** Magnitude offsets between the same image of the field SE65 used in Figure 1 and the reference frame of the same field.

than 14.0 mag. On the other hand, the stars brighter than 12.5 mag tend to be at the edge of saturation on the TNTS images. We thus chose those stars with magnitudes between 12.5 and 14.0 mag to determine the magnitude offset for each image. To mitigate the effect of nonuniform background due to moonlight or obscuration by clouds, we divided every image into 16 small subsections and then calculated the offset of each subsection to perform the photometric calibrations. Outlier rejections of  $3\sigma$  were applied iteratively to make sure that no variables were involved in the calculations. The median value of the magnitude offsets of 16 subsections is almost the same as that of the whole image, e.g.,  $\Delta M_{zp} = 0.010 \pm 0.001$  for the sky field SE65. The uncertainty in the magnitude offset is so small (relative to the photometric errors from SExtractor) that we neglected it in the following calculations.

Previous studies indicate that unfiltered magnitude is approximately equivalent to that in the *R* band of the PPMX (Position and Proper Motions eXtended) system (Röser et al. 2008) with an uncertainty of 0.2–0.3 mag depending on the number of flux-calibrated stars available in the field (e.g., Li et al. 2003, 2011). For each field of the TNTS, we match the reference image with the PPMX catalog to choose appropriate stars for calibrating the unfiltered magnitudes. Owing to the large FOV of the TNTS images, the number of stars that can be used for the flux calibration is relatively large. The median uncertainty in the photometric calibration for the 600 fields used in our study is  $\sim 0.010 \pm 0.005$  mag, which is much smaller than that from Li et al. (2003).

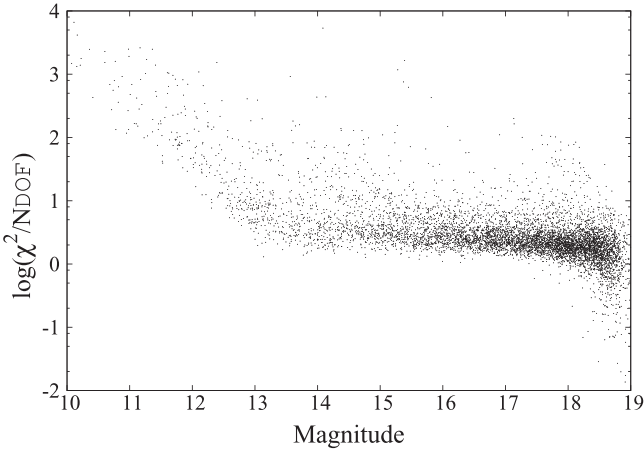
#### 3.2. Rescaling the Measured Photometric Errors

For some reason, the photometric errors given by SExtractor might be overestimated or underestimated. We would either miss real variables or select bogus variables as real ones (Kaluzny et al. 1998). Assuming that the photometric errors roughly follow the Gaussian distribution, the logarithm of the  $\chi^2/N_{\text{DOF}}$  of stars with constant luminosities should be close to zero (see the right panel of Figure 6 from Kaluzny et al. 1998). The  $\chi^2/N_{\text{DOF}}$  is defined as

$$\chi^2 = \left( \frac{m - \bar{m}}{\sigma_m} \right)^2 \quad (1)$$

$$N_{\text{DOF}} = N - M, \quad (2)$$

<sup>8</sup> <https://pypi.python.org/pypi/SRPAstro>



**Figure 3.**  $\log(\chi^2/N_{\text{DOF}})$  distribution as a function of the weighted mean magnitude for stars in the field SE65 before rescaling their photometric errors.

where  $\bar{m}$  is the weighted mean value (the weighting factor is the inverse square of the error  $\sigma_m$ ) of the time-series magnitude  $m$  of a star,  $N$  is the number of measurements, and  $M$  is the number of parameters (here we let it be one). Take the field SE65 as an example. Figure 3 shows the plot of the logarithm of the  $\chi^2/N_{\text{DOF}}$  against weighted mean magnitude for all of the stars. One can see that the brightest stars ( $\sim 12.0$  mag) have  $\chi^2/N_{\text{DOF}} \sim 100$ , and their photometric errors are thus seriously underestimated. The possible reasons for this discrepancy include underestimation of the flat-fielding error, imperfect aperture photometry, etc. To eliminate this inconsistency, the magnitude error  $\sigma$  is multiplied by an appropriate scaling factor  $F$  following the method proposed by Kaluzny et al. (1998).

In detail, we first binned the data  $\chi^2/N_{\text{DOF}}$  at an interval of 0.1 mag. In the logarithm plot, we used a second-order polynomial to fit the data for stars with brightness  $< 13.5$  mag. For stars with brightness between 13.5 and 18.0 mag, we used a linear function for the fit, while those stars with weighted mean magnitudes fainter than 18.0 mag were discarded in the fit due to relatively larger photometric errors. We then applied the scaling factor  $F$  to rescale the photometric error of each star. The distribution of the rescaled photometric errors for stars in the field SE65 is shown in Figure 4.

#### 4. VARIABLES FROM THE TNTS

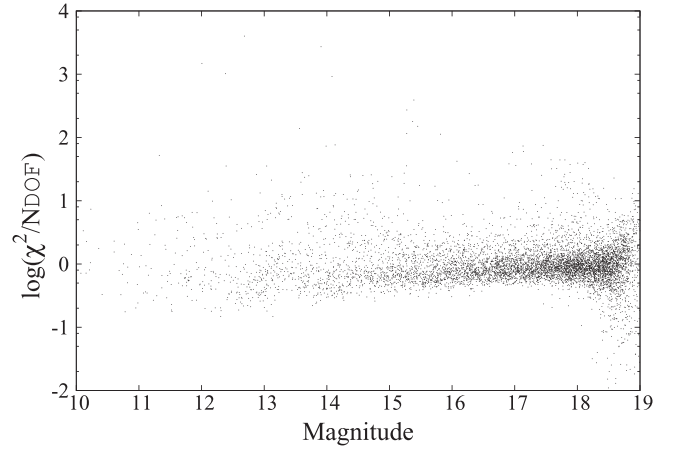
##### 4.1. Searching for Variable Stars

The variable candidates were initially selected using the Welch–Stetson variability index  $J_s$  (Stetson 1996). We used VARTOOLS (Hartman et al. 2008) to compute the  $J_s$  index based on the rescaled errors of photometry. For each star one can calculate the variability index  $J$  as

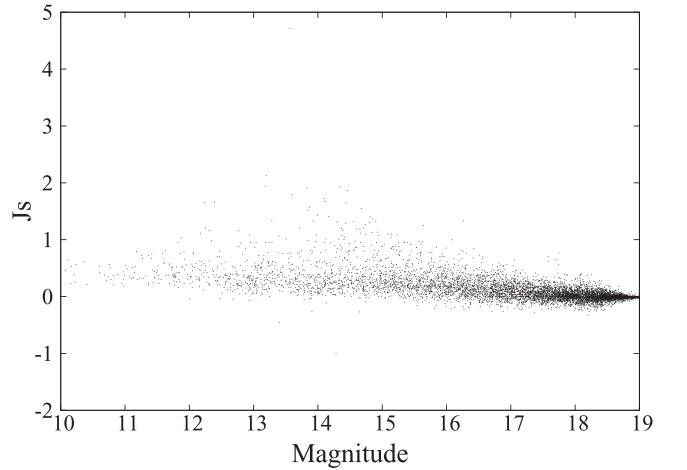
$$J = \frac{\sum_{k=1}^n w_k \operatorname{sgn}(P_k) \sqrt{|P_k|}}{\sum_{k=1}^n w_k}, \quad (3)$$

where  $\operatorname{sgn}()$  is the sign function. We considered  $n$  pairs of observations and each has a weight  $w_k$ ,

$$P_k = \begin{cases} \delta_{i(k)} \delta_{j(k)}, & \text{if } i(k) \neq j(k), \\ \delta_{i(k)}^2 - 1, & \text{if } i(k) = j(k) \end{cases} \quad (4)$$



**Figure 4.** Same diagram as in Figure 3 but after rescaling for photometric errors.



**Figure 5.** Distribution of the  $J_s$  index calculated for stars in the sky field SE65.

where  $P_k$  is the product of the normalized residuals of the two paired observations  $i$  and  $j$  and

$$\delta = \sqrt{\frac{n}{n-1} \frac{m - \bar{m}}{\sigma_m}} \quad (5)$$

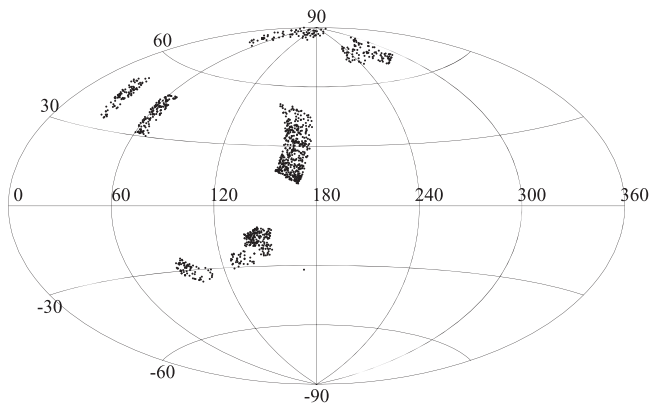
represents the magnitude residual from the mean value scaled by the standard error (Stetson 1996). In detail,  $n$  is the total number of observations contributing to the mean. The final variability index is

$$J_s = J \frac{\sum w_k}{w_{\max}}. \quad (6)$$

In our calculations, we considered the two observations taken on the same night as a pair and give each of them a weight of 0.5. Otherwise a weight of 1.0 was given if there is only a single observation.

As an example, we show in Figure 5 the distribution of  $J_s$  obtained for stars in the field SE65. As expected, most stars have  $J_s$  values close to zero. The mean  $J_s$  value of 600 sky fields involved in this study is 0.07. For each field, we chose those stars with  $J_s$  larger than the mean value by  $3\sigma$  as candidates of variable stars.

We searched for periodic variabilities among candidates using the Lomb–Scargle (L–S) method (Lomb 1976;



**Figure 6.** Distribution of the TNTS variable stars in the Galactic coordinates.

Scargle 1982; Press & Rybicki 1989; Press et al. 1992) and Phase Dispersion Minimization (PDM) method (Stellingwerf 1978). The L–S method applies the statistical properties of least-square frequency analysis of unequally spaced data on a series of test periods. We set test periods ranging from 0.1 to 600 days with a bin size of 0.1 day. We determined the most likely period for each star in terms of the power value given by the L–S method. Then we performed the PDM method to double check the periods detected from the L–S method since the PDM method is useful for data sets with gaps, non-sinusoidal variations, poor time coverage, or other problems that would make the Fourier technique unusable. The folded light curves were then established using the period from the L–S/PDM or twice that for a binary if necessary. The period derived from the PDM method was only adopted if the phased light curve folded by the PDM has smaller scatter than that from the LS method. We examined the light curve and phase diagram for each candidate to pick out variables and throw out those showing alias variations (e.g., one day, a half day, one-third day) and having large dispersions. Finally we found 1237 variables with magnitude variations  $\gtrsim 0.1$  mag and brightness  $< 18.0$  mag based on the L–S method.

All of these candidates had been cross-checked with those listed in the VSX database (the International Variable Star Index, Watson et al. 2015), which is created by amateur astronomers of the American Association of Variable Star Observers (AAVSO).<sup>9</sup> We found that 938 out of these 1237 objects have been included in this database with 438 from the Catalina Sky Survey (CSS) sample (Drake et al. 2014), 106 from the Lincoln Near-Earth Asteroid Research (LINEAR) survey sample (Palaversa et al. 2013), and the remaining 299 can be considered new variables. Figure 6 shows the distribution of these variables, in which they are displayed according to the Galactic coordinates. The parameters of the first 20 variable stars from the variable star candidates table are listed in Table 1. Columns 1 and 2 list the survey field and the ID number of the field, Columns 3 and 4 give the Right Ascension and declination of the relevant field, Column 5 gives the mean unfiltered magnitude (calibrated with the *R*-band magnitudes of the PPMX catalog), Column 6 gives the line of sight extinction based on Schlafly & Finkbeiner (2011), Column 7 gives the amplitude of variation (magnitude from maximum to minimum), Column 8 gives the variability index, Columns 9 and 10 give the periods of the variables from our

measurements and the corresponding method, Column 11 gives the period from the VSX, Column 12 gives the minimum light of the variables detected in the observation season, Column 13 gives a tentative classification of the variable star whenever possible, and Column 14 shows the additional information about the variable stars, including previous identifications from the VSX database (e.g., CSS or LINEAR). Note that for those known variables, the periods estimated from the TNTS are highly consistent with those given by the VSX database except for about 30 (see Figure 7). The phased light curves of new periodic variable stars are shown in the Figure 8.

#### 4.2. Classification of Variables

Most of the variables exhibit regular variations in periods and amplitudes, i.e., RR Lyrae and eclipsing binaries. Table 2 lists the types of the 1237 variables from the first two years of survey data, which are classified based on the shapes of their light curves, amplitudes, and periods of phase diagrams. Half of the detected variables are found to be RR Lyrae according to their specific periods and light curve shapes. The remaining variables include 431 eclipsing binaries, 72 semi-regular pulsators, 29 Mira stars, 11 slow irregular variables, 11 RS Canum Venaticorum stars, 7 Gamma Doradus stars, 5 long-period variables, 3 W Virginis stars, 3 Delta Scuti stars, 2 anomalous Cepheids, 1 Cepheid, and 1 nova-like star. Note that our tentative classifications of variable stars are based only on the unfiltered photometry.

##### 4.2.1. RR Lyrae Stars

RR Lyrae stars, located at intersections of the instability strip and horizontal branch, belong to a subclass of pulsating variable stars. They are old, low-mass, and metal-poor “Population II” stars which are usually found in globular clusters. According to their pulsation modes, the RR Lyrae stars can be divided into three subtypes: RRAB-type (pulsating in the fundamental mode), RRC-type (pulsating in the first overtone), and RRD-type (pulsating simultaneously in two modes; Smith 1995). Meanwhile, different subtypes of variable stars have different periods and amplitudes. In our sample, there are 526 RRAB-type stars, 126 RRC-type stars, and 9 RRD-type stars in light of the periods and amplitudes of variations, which are shown in Figure 9. For our sample, the mean period of the RRAB-type variables is 0.57 day, while the RRC-type variables have a mean period of 0.34 day. The mean amplitude of the light variation for the RRAB subgroup (0.84 mag) is relatively larger than that of the RRC subgroup (0.44 mag).

The light curves of RRAB Lyrae stars are characterized by variations with a fast rising followed by a slow declining (Cacciari et al. 2005; Jurcsik et al. 2009). The variable star S100-3820 (see Figure 8) was originally detected by AASVO VSX and identified as an RRAB-type star. However, our observation shows that it has a period of  $\sim 0.5$  day and a variation amplitude of 0.37 mag, lying at the boundary between RRAB-type and RRC-type (see Figure 9). According to our result, its rise time is equal to the decline time, which likely just evolved off the zero-age horizontal branch (ZAHB) and exhibits light variation with a fixed amplitude and a longer period (Cacciari et al. 2005, Section 3.2). Another two interesting RR Lyrae stars in our sample are S037-10947 (see Figure 8) and S585-9662 (see Figure 8), which lie at a similar

<sup>9</sup> AAVSO: <http://www.aavso.org/vsx/>

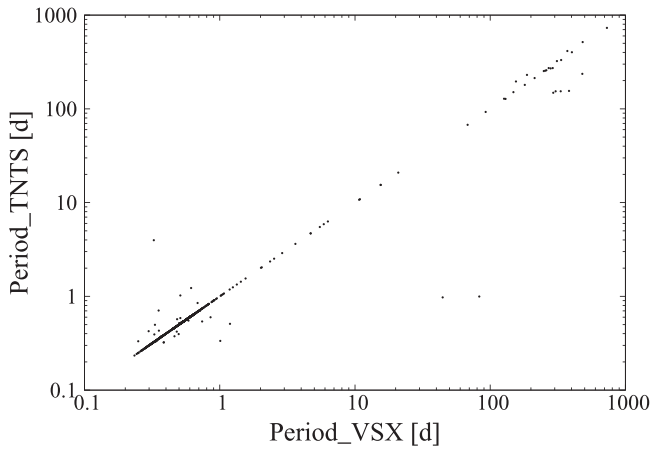


**Table 1**  
Variable Stars

(1) Field	(2) ID	(3) R.A. (J2000.0) <sup>a</sup>	(4) decl.	(5) Unfilter	(6) $A_v$ (mag)	(7) Amplitude	(8) $J_s$	(9) Period (TNTS) (d)	(10) Src <sup>b</sup>	(11) Period (VSX) (d)	(12) $T_0^c$ (HJD-2456000)	(13) Type <sup>d</sup>	(14) Note <sup>e</sup>
S007	424	22:47:55.16	+25:46:59.60	13.88	0.19	0.33	0.89	0.382994	(LS)	0.382989	196.824	EC	V, C
S008	1839	22:51:42.89	+25:32:21.01	14.59	0.35	0.45	1.59	0.354250	(LS)	0.354265	196.650	RRC	V, C
S008	542	22:53:00.98	+25:43:48.03	14.84	0.34	0.23	1.23	46.094477	(LS)	...	214.880	SR	N
S008	729	22:53:18.25	+25:42:40.00	14.35	0.35	1.11	8.48	0.471269	(LS)	0.471272	196.318	RRAB	V
S008	8407	22:54:12.33	+24:38:07.03	15.41	0.26	1.02	2.46	0.608931	(LS)	0.608926	...	RRAB	V, C
S008	8461	22:54:18.92	+24:37:59.08	13.59	0.26	0.25	0.77	13.345449	(LS)	...	190.193	SR	N
S009	9362	23:01:39.04	+24:26:04.04	14.97	0.43	0.41	1.20	0.359450	(PDM)	0.359444	196.675	RRC	V, C
S011	4138	23:11:27.60	+25:08:42.01	17.18	0.18	1.04	0.85	0.546095	(LS)	0.546052	196.395	RRAB	V, C
S012	9389	23:18:25.72	+24:25:54.09	15.57	0.20	0.46	1.35	0.305437	(LS)	0.305435	196.489	RRC	V, C
S012	9446	23:21:42.07	+24:28:40.02	14.52	0.14	0.86	5.58	0.563126	(LS)	0.563094	196.381	RRAB	V, C
S013	7367	23:24:04.12	+24:32:33.09	16.61	0.11	1.00	1.57	0.509679	(PDM)	0.509708	196.495	RRAB	V, C
S014	4043	23:35:53.93	+25:08:56.05	16.41	0.14	1.00	1.03	0.549363	(LS)	0.549271	196.643	RRAB	V, C
S014	6689	23:32:02.64	+24:38:44.07	14.74	0.21	1.83	11.25	...	...	165.000000	229.217	L	V
S014	6967	23:32:26.03	+24:38:28.04	10.47	0.21	2.00	0.97	...	...	194.000000	...	L	V
S015	7103	23:40:47.85	+24:25:50.06	17.01	0.12	1.23	1.26	0.478438	(PDM)	0.478459	...	RRAB	V, C
S037	10947	22:45:05.01	+26:04:52.02	15.31	0.17	0.39	1.33	0.497227	(LS)	0.331645	196.565	RRAB	V, C
S037	8742	22:43:33.57	+26:18:47.05	17.93	0.14	1.31	0.59	0.526378	(LS)	0.526361	196.609	RRAB	V, C
S038	10429	22:52:06.31	+25:59:50.02	16.34	0.37	1.18	1.13	0.489781	(PDM)	0.489789	196.636	RRAB	V, C
S038	12180	22:53:18.19	+25:42:47.00	14.44	0.35	1.10	4.87	0.471307	(LS)	0.471272	196.295	RRAB	V
S039	10942	22:59:08.64	+25:49:50.09	14.78	0.36	0.90	5.48	0.497052	(LS)	...	196.866	RRAB	N

**Note.**<sup>a</sup> Based on the TNTS reference images.<sup>b</sup> L–S: Lomb–Scargle method; PDM: Phase Dispersion Minimization method; VSX: from VSX catalogs.<sup>c</sup> Epoch when primary eclipse or minimum light occurs.<sup>d</sup> ACEP—anomalous Cepheid; RS—RS Canum Venaticorum; EC—contact binary; ES—semi-detached binary; ED—detached binary; CWA—W Virginis; RR—RR Lyrae variable; BL—“Blazhko effect”; CEP—Cepheid; GDOR—Gamma Doradus star; DSCT—Delta Scuti star; SR—semi-regular pulsator; M—Mira variable; LPV—long-period variable; L—slow irregular variable.<sup>e</sup> V—AAVSO VSX variable; C—CSS variable; L—LINEAR variable; N—New variable.

(This table is available in its entirety in machine-readable and Virtual Observatory (VO) forms.)



**Figure 7.** Measured periods from the TNTS data vs. those from the AAVSO VSX for the variables in common.

**Table 2**  
Distribution of Types of Variables from the TNTS

Variable Type	$N$	(%)	$N$ (New)
RR lyr	661	53.44	59
Eclipsing binaries	431	34.84	157
Semiregular pulsators	72	5.82	68
Mira (LPV)	34	2.75	0
Irregular	11	0.89	0
RS Canum Venaticorum	11	0.89	2
Gamma Doradus stars	7	0.57	7
W Virginis	3	0.24	0
Delta Scuti Stars	3	0.24	3
Anomalous Cepheids	2	0.16	1
Cepheids	1	0.08	1
Nova	1	0.08	1

position as S100-3820 in Figure 9. Although they have typical features of this type of variable star, i.e., fast rising and slow declining, they show a little bump close to the minimum light. More data are needed to determine the properties of these two interesting RR Lyrae stars.

Some RR Lyrae stars are known to have the “Blazhko effect” in which they exhibit modulations in their periods, pulsation amplitudes, light curve shapes, and radial velocities (Blazhko 1907; Buchler & Kolláth 2011). Previous surveys have found that 10%–30% of the RRAB-type variables have such an effect (Alcock et al. 2003) and this ratio can reach up to 50% according to a recent study by Jurcsik et al. (2009). Fourteen RRAB-type Lyrae stars of our sample are found to have phase modulations and can be tentatively classified as RRAB stars with the Blazhko effect. Obviously, we did not detect such a high ratio from our sample when compared with previous surveys. It is likely that the current TNTS data may not be qualified enough to see most of the phase modulations. Increasing the cadence of the survey may help to discover more RR Lyrae stars with the Blazhko effect.

The period–amplitude diagram for RR Lyrae stars, also termed the Bailey diagram, was used to investigate their properties in globular clusters more than a century ago (Bailey 1902). Inspecting our RRAB stars shown in Figure 9, we find that these variables have different locations in this diagram. For example the periods could differ by  $\sim 0.1$ – $0.2$  day at a given amplitude. There has been debate over whether

Oosterhoff-type or metallicity of the clusters accounts for these differences. The Oosterhoff-type classification of the Galactic globular clusters was based on the frequency distribution of their RR Lyrae variables’ periods (Oosterhoff 1939). For example, RRAB stars in Oosterhoff I (OoI) clusters have mean periods of  $\sim 0.55$  day, while those in Oosterhoff II (OoII) clusters have mean periods of  $\sim 0.64$  day (i.e., Smith 1995; Clement et al. 2001; Catelan 2009). On the other hand, the metallicity has been found to correlate with both period and amplitude for the RR Lyrae stars (i.e., Marconi et al. 2015 and references therein). It has been found that the observed period of RR Lyrae stars tends to become shorter with increasing metallicity (i.e., Preston 1959; Kunder & Chaboyer 2009; Fiorentino et al. 2015), while the pulsation amplitude decreases as the metallicity increases (i.e., Bono et al. 1996). Such an amplitude–period–metallicity relation can be used to estimate metallicity for different RRAB Lyrae stars (Sandage 1982, 2004, Alcock et al. 2000, Kinemuchi et al. 2006). The metallicity dependence of the period–amplitude relation can be interpreted as an opacity effect in stellar evolution (i.e., Bono et al. 1996), but a detailed explanation is beyond the scope of this paper. Note that the effect of Oosterhoff classification on the period–amplitude relation may be in part coupled with that of metallicity, as OoI clusters tend to have higher metallicity than clusters of OoII (Catelan 2009).

Note that the amplitudes of RR Lyraes shown in Figure 9 are measured in the unfiltered light curves, which are found to be systematically smaller by 0.15 mag than those measured in the  $V$  band (i.e., Zorotovic et al. 2010). A similar offset was noticed in the unfiltered data from other surveys such as the CSS (Drake et al. 2013) and SSS (the Siding Spring Survey, Torrealba et al. 2015). We thus convert the unfiltered amplitude from the TNTS to the  $V$ -band value by adding an offset of 0.15 mag for each star of our sample. The corrected period–amplitude diagram of the RRAB-type stars is shown in the inset of Figure 9, where we add the reference line of OoI (solid line) based on Equation (11) in Zorotovic et al. (2010). Overplotted is the OoII line (dashed line) which follows the same trend as OoI but shifted in period by  $\Delta \log P = +0.06$ . Following Miceli et al. (2008), a period shift of 0.05 day relative to the OoI line was adopted to divide our RRAB Lyrae stars into the OoI and OoII types. One can see that the majority of the TNTS RRAB Lyrae stars cluster around the OoI line, which is consistent with that inferred from the CSS sample extending up to  $\sim 60$  kpc above the Galactic plane (see Figure 5 of Drake et al. 2013) and also consistent with that from the Siding Spring Survey (see Figure 11 of Torrealba et al. 2015).

Since we have a relatively large sample of RRAB Lyrae stars located from the disk to the halo of the Galaxy, we can carry on studies on the spatial distributions of their metallicities and this helps trace history of the Galaxy formation (Catelan 2009; Pietrukowicz et al. 2014). The metallicity is estimated using the empirical period–amplitude–metallicity relation given by Sandage (2004), which follows as

$$[\text{Fe}/\text{H}] = - (1.453 \pm 0.027)A_V - (7.990 \pm 0.091)\log P - (2.145 \pm 0.025), \quad (7)$$

where  $A_V$  is the  $V$ -band amplitude corrected from the unfiltered value and  $P$  represents the period. The amplitude–period–metallicity relation gives an estimate of the metallicity that is

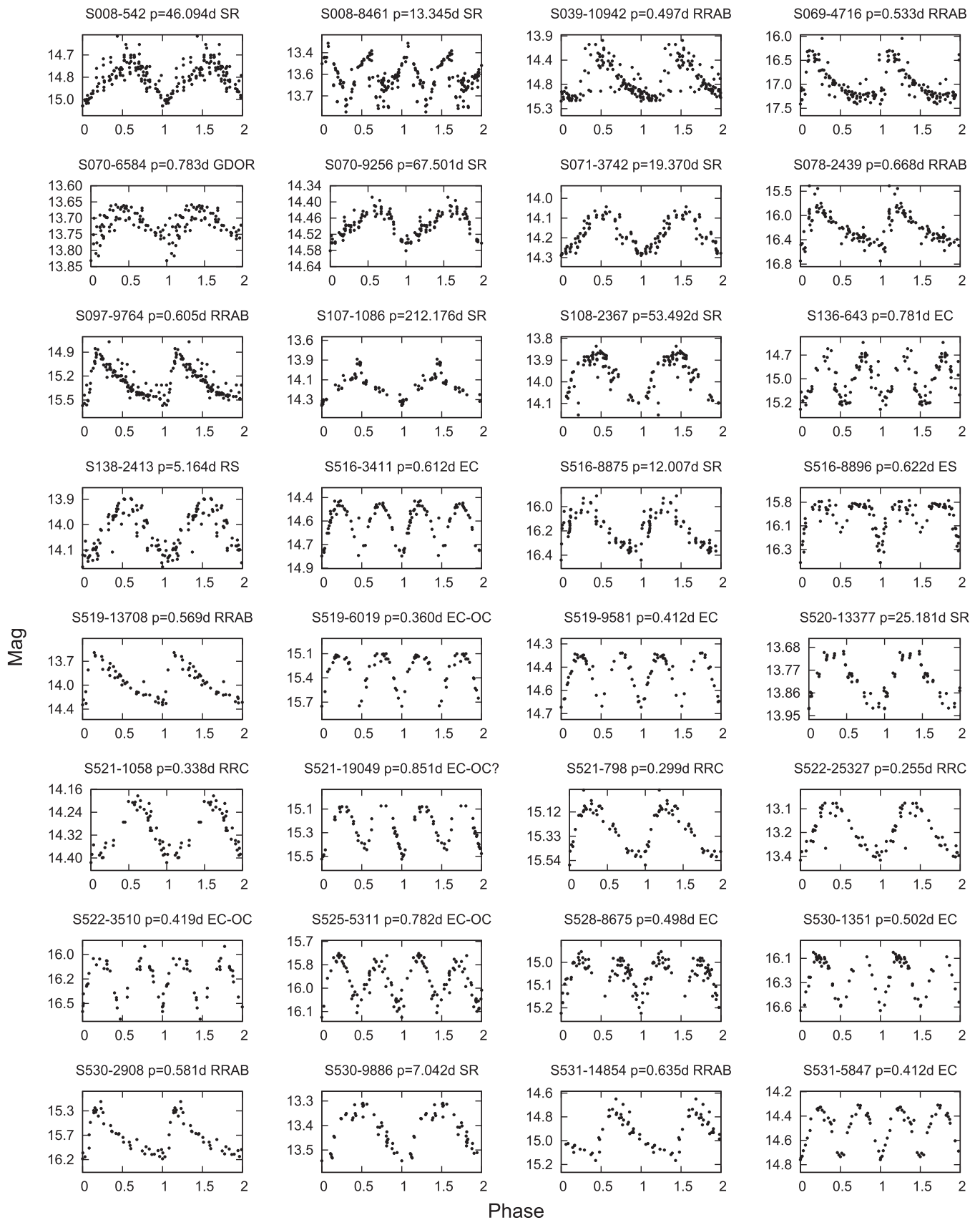


Figure 8. Phase diagrams of new variables, including two special RRAB stars that are not new.

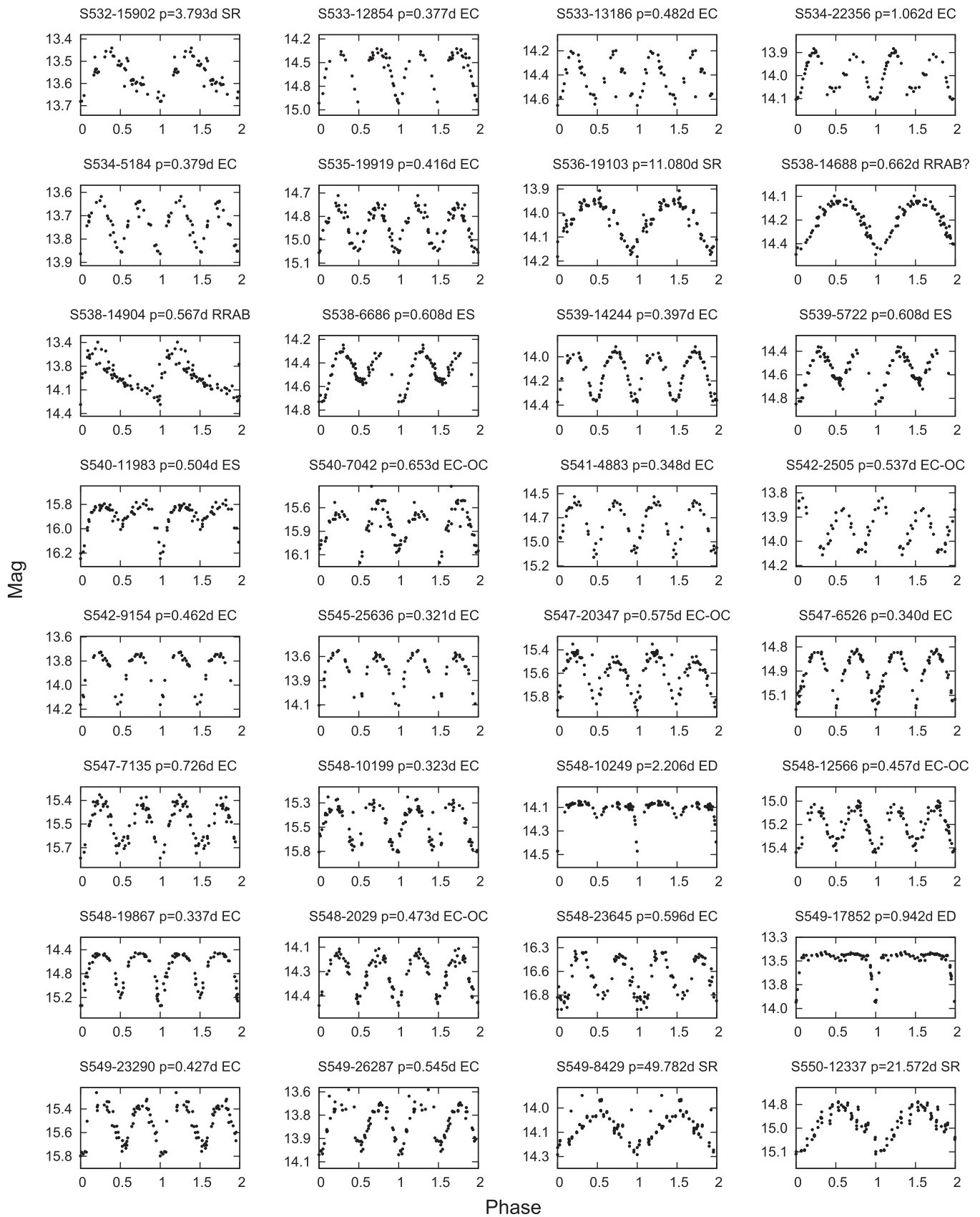


Figure 8. (Continued.)



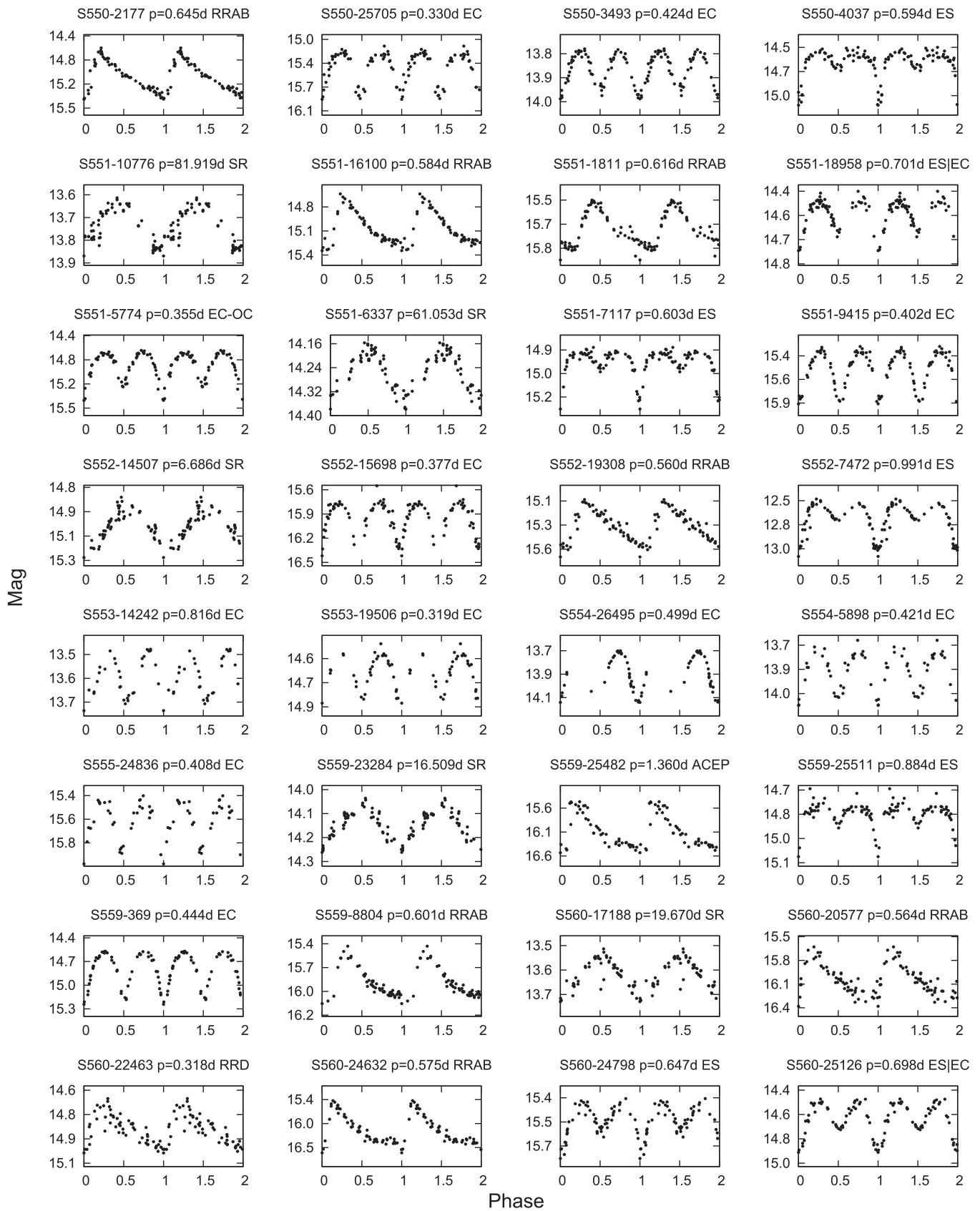


Figure 8. (Continued.)

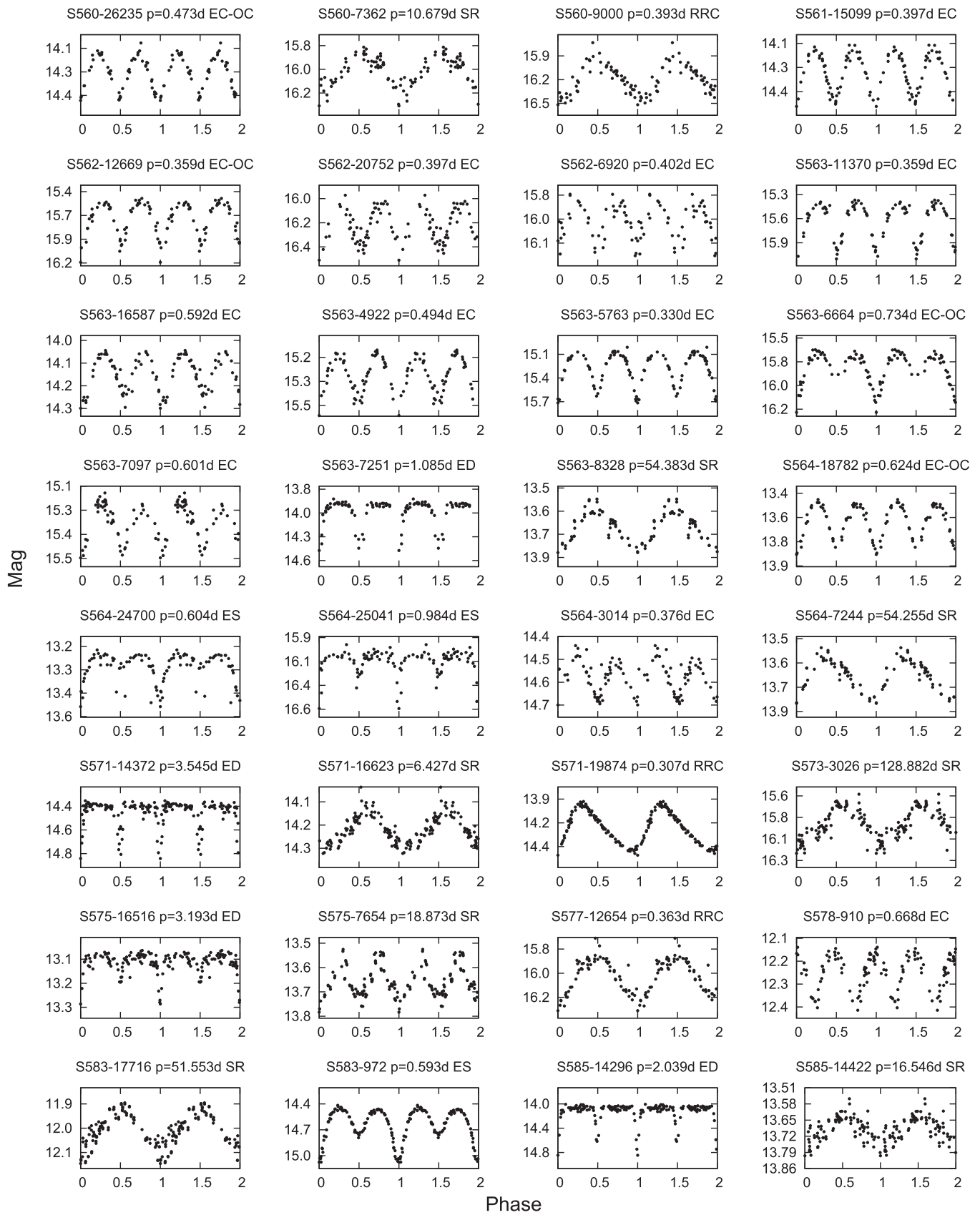


Figure 8. (Continued.)

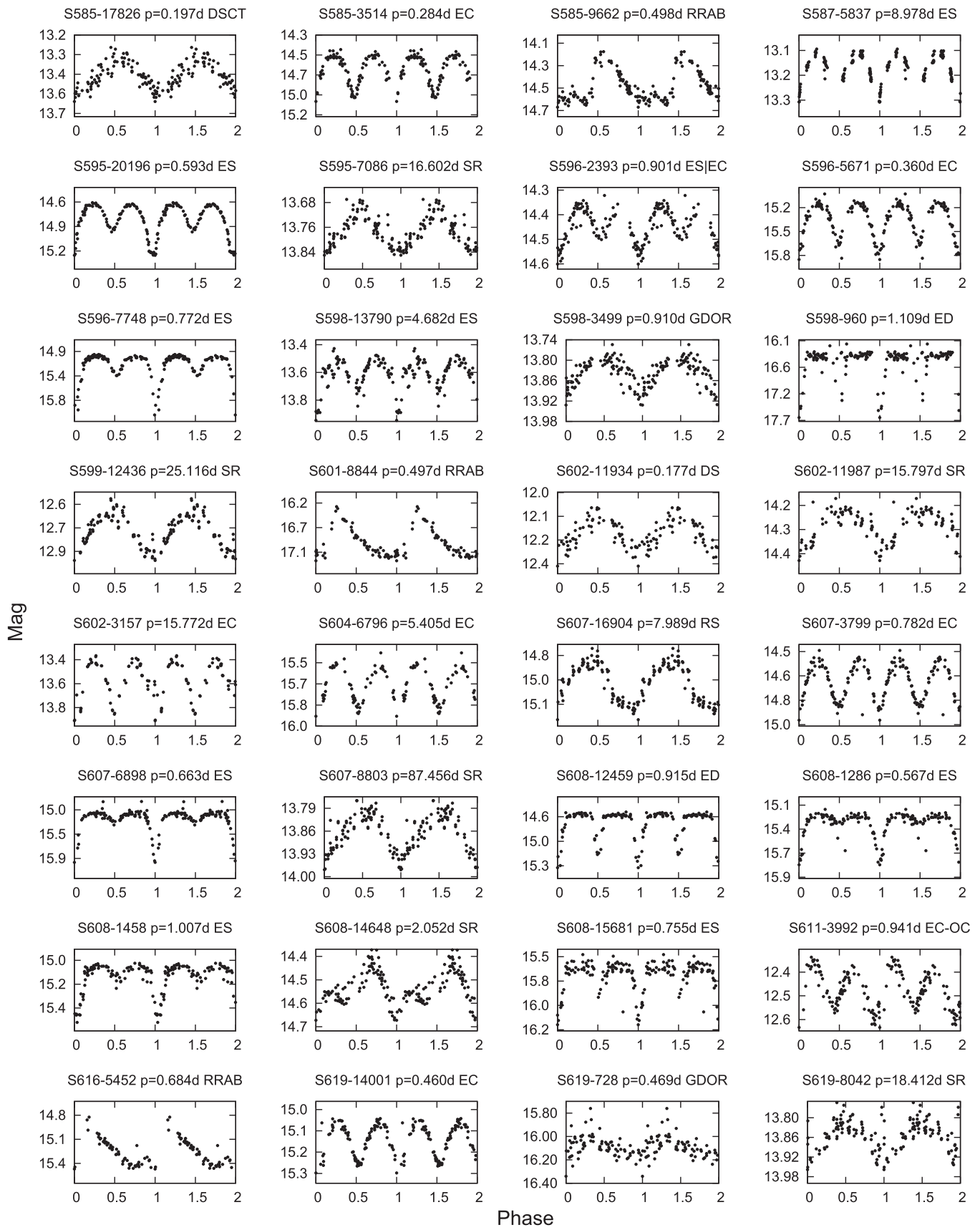


Figure 8. (Continued.)

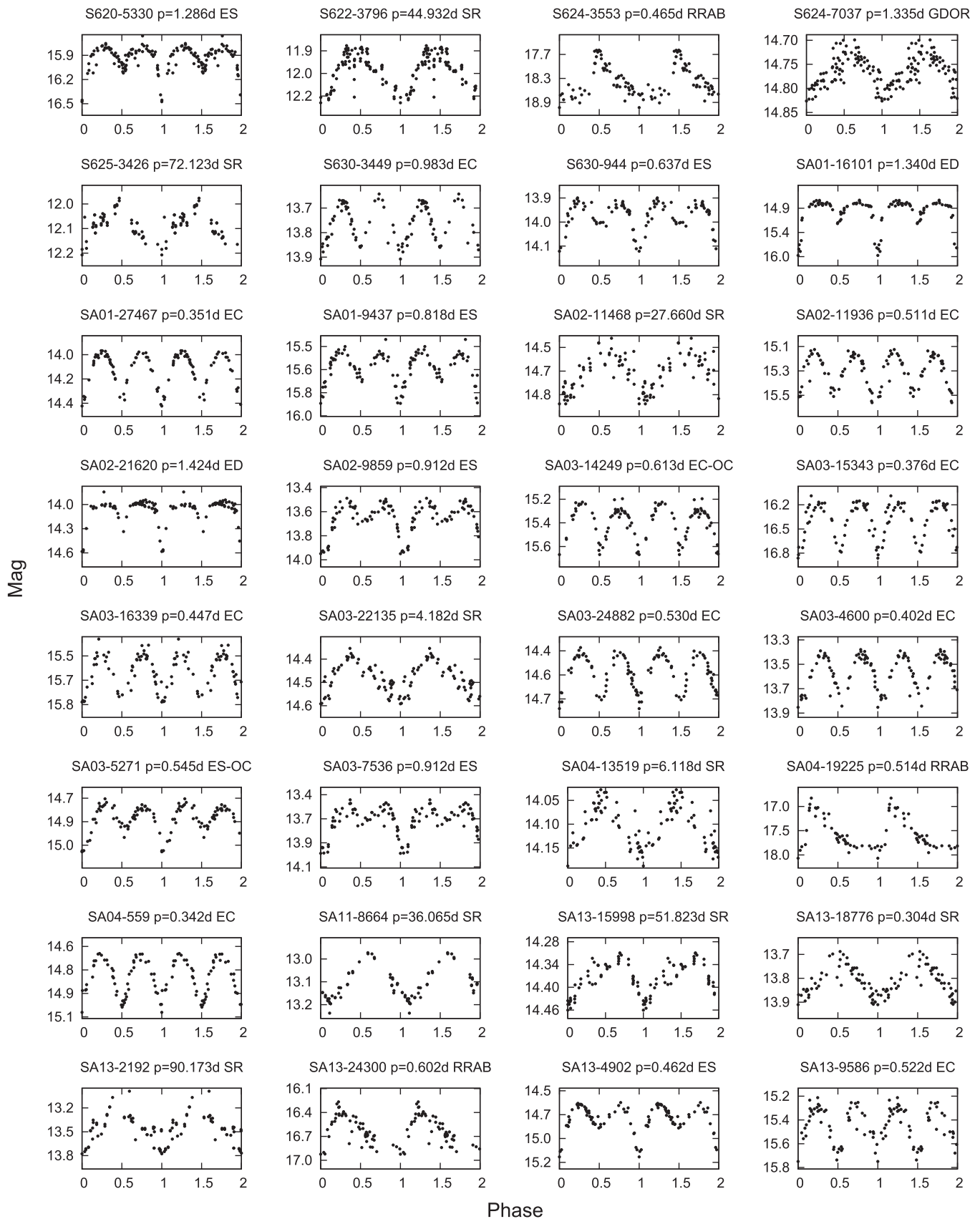
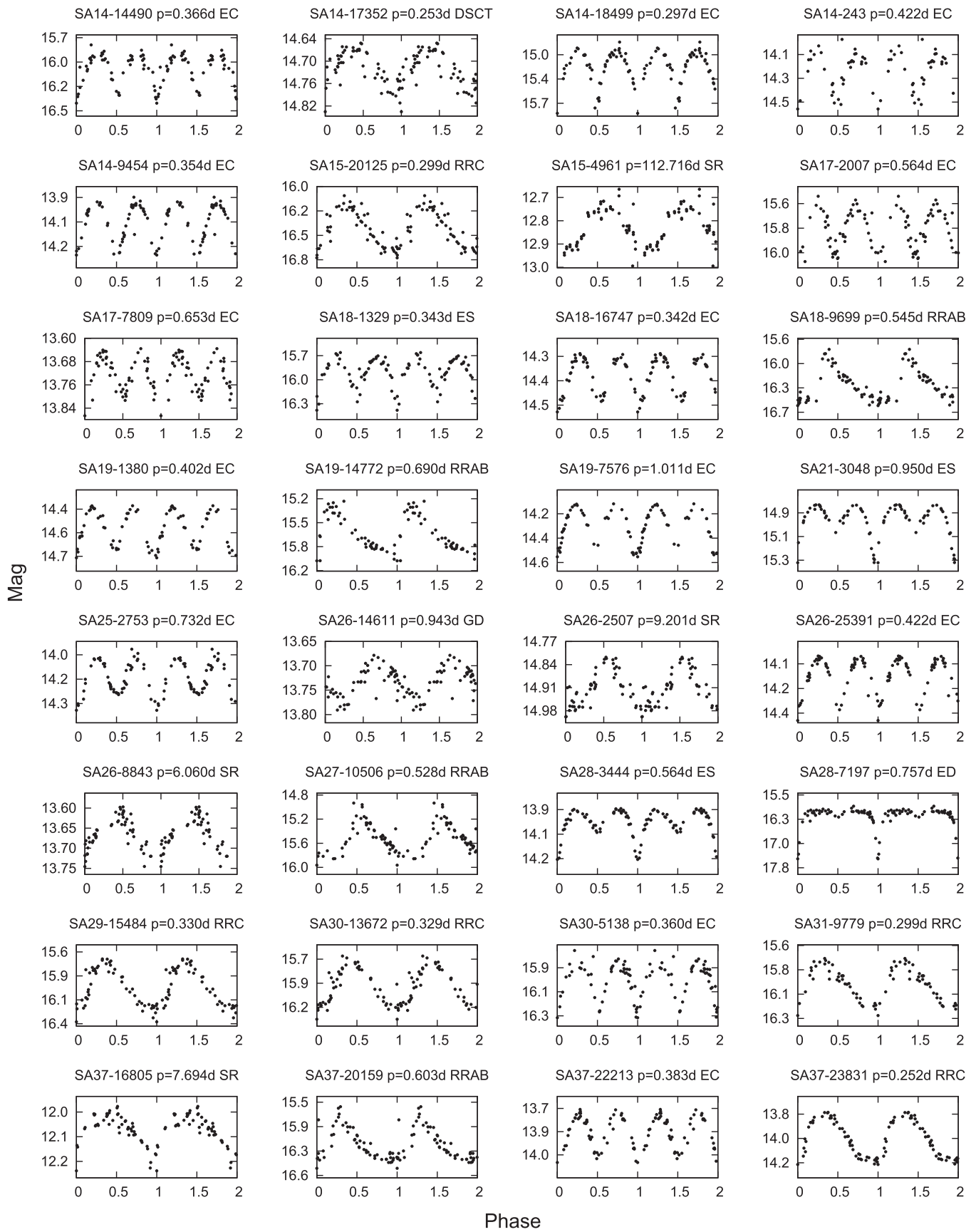


Figure 8. (Continued.)



Phase  
Figure 8. (Continued.)



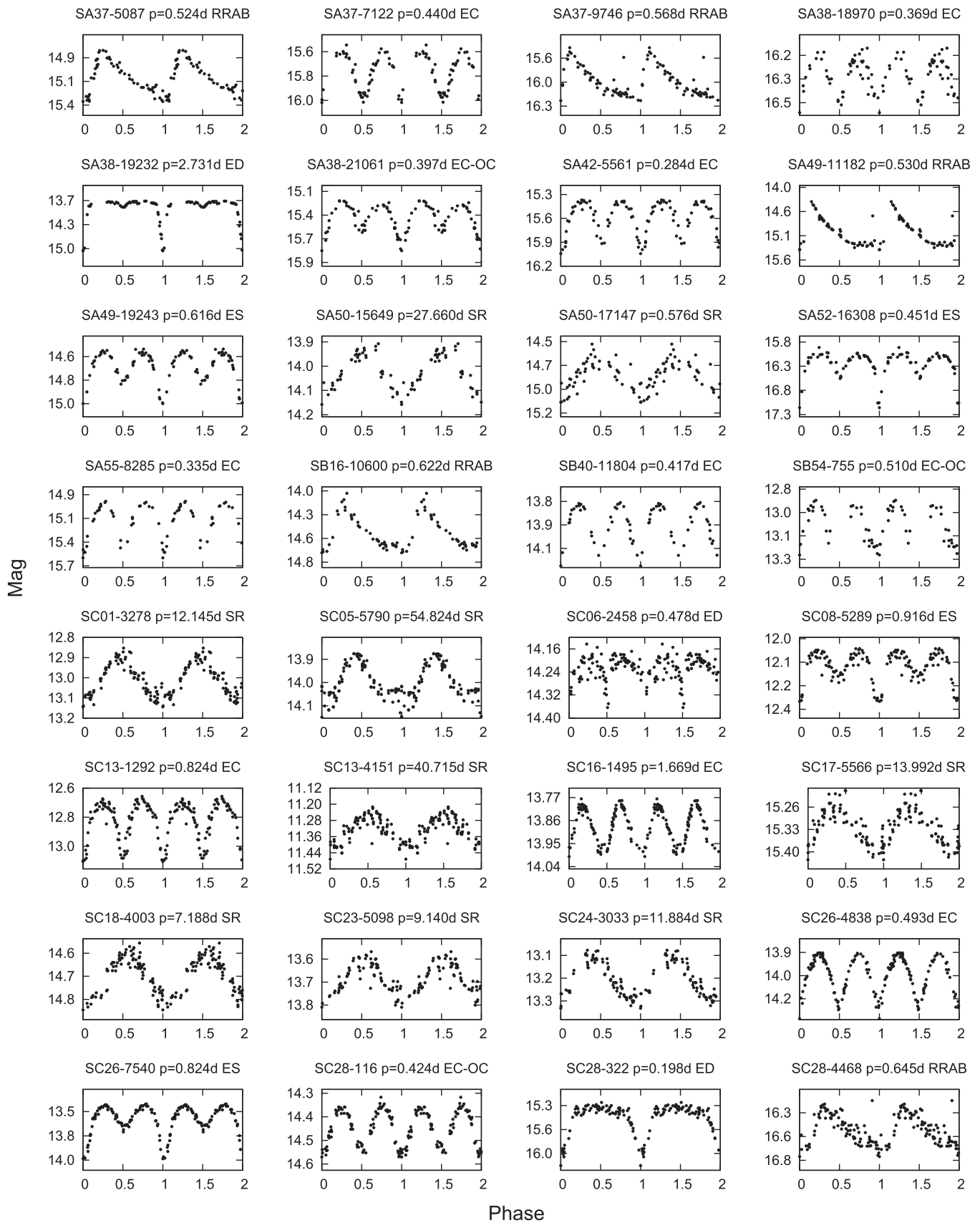


Figure 8. (Continued.)

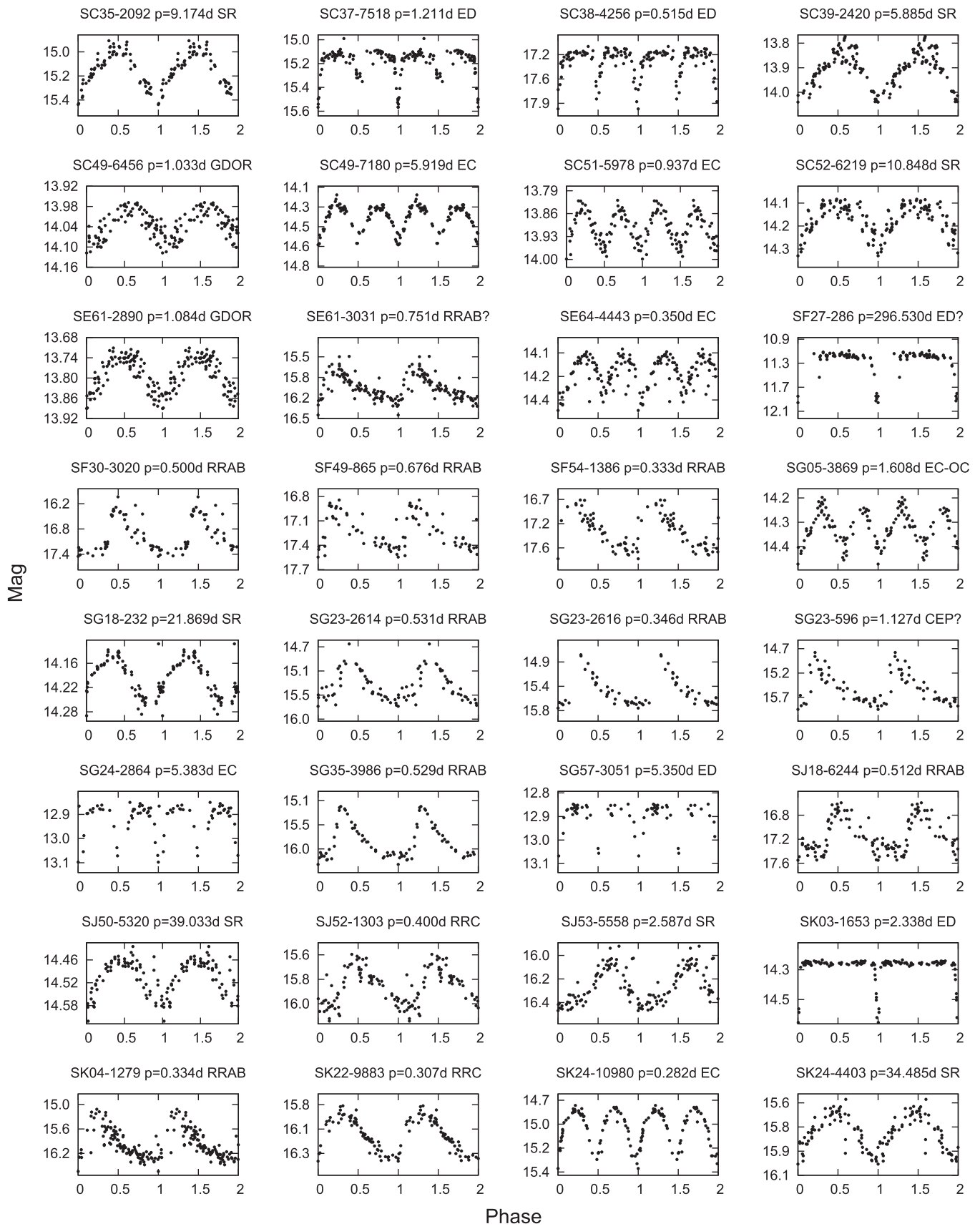


Figure 8. (Continued.)

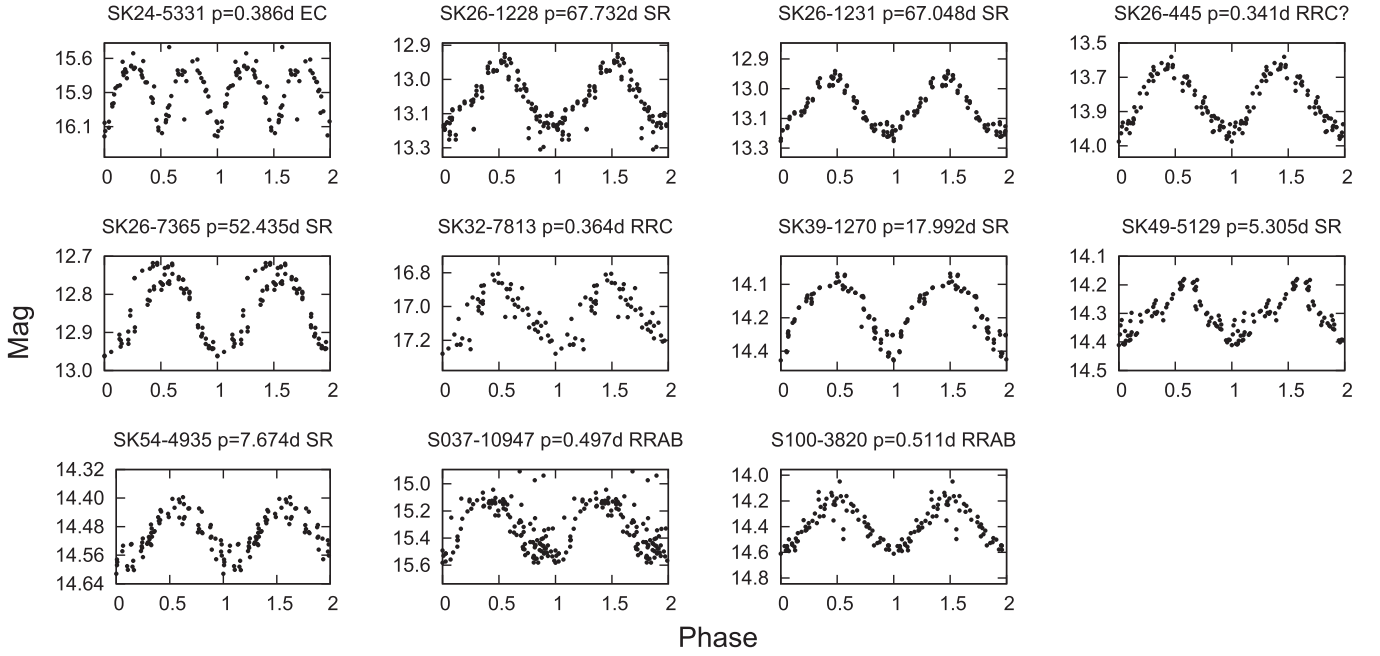
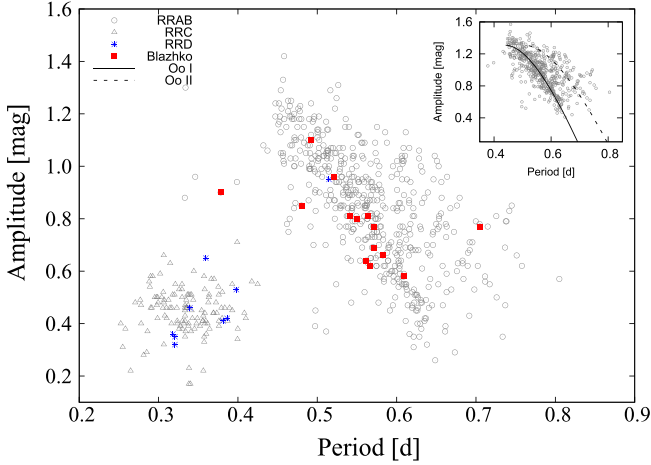


Figure 8. (Continued.)



**Figure 9.** Amplitude of light variation as a function of the period for the TNTS RR Lyrae stars. The circles and triangles represent the RRAB-type stars (fundamental overtone) and RRC-type stars (first overtone), respectively. The blue points and red points represent RRD-type stars and RRAB stars with the Blazhko effect. The inset panel shows the Bailey diagram of the RRAB-type stars, with the unfiltered amplitudes converted into the  $V$ -band values. The solid and dashed lines represent the locations of the Oosterhoff I and Oosterhoff II subtypes (Zorotovic et al. 2010).

generally consistent with the spectroscopic method, with an uncertainty of 0.3 dex (see Figure 9 in Kinemuchi et al. 2006).

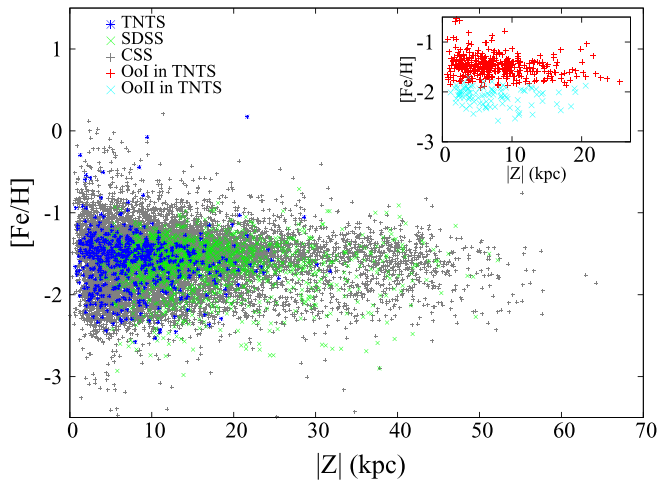
The distance to the RRAB Lyrae stars can be calculated using the relation given as:

$$d = 10[(V_0)_s - M_V + 5]/5 \quad (8)$$

where  $(V_0)_s$  represents the static average magnitudes in the  $V$  band. The Galactic extinctions of the TNTS unfiltered magnitudes were corrected from the recalibrated infrared dust maps (Schlafly & Finkbeiner 2011). For the RR Lyrae sample

in common, we found that the TNTS unfiltered magnitudes have an offset of  $\sim 0.1$  mag relative to the  $V$ -band magnitudes from Drake et al. (2013). This difference was corrected in calculating the distances to the TNTS RRAB Lyrae stars. Assuming an averaged absolute magnitude  $M_V = 0.6$  mag for the RRAB Lyrae stars (Keller et al. 2008), we can derive their heliocentric distances from the TNTS data. Then we converted these distances into those relative to the Galactic plane of the Milky Way (i.e.,  $|Z|$ ).

Figure 10 shows the spatial distribution of the metallicity derived for the TNTS RRAB Lyrae. The measurements from the CSS photometry and the SDSS spectra are overplotted for comparison. As can be seen, the photometric metallicities from the TNTS and CSS observations are consistent with the spectroscopic estimates from the SDSS sample. The mean metallicity from the TNTS, the CSS, and the SDSS is  $-1.59$  dex,  $-1.57$  dex, and  $-1.64$  dex, respectively. The inset of Figure 10 shows the TNTS RRAB Lyrae distinguished by the Oosterhoff types. From Figures 9 and 10, one can see that the OoI population can occur at larger distances perpendicular to the Galactic plane (i.e.,  $|Z| \sim 45$  kpc), and they have a higher mean metallicity (i.e.,  $[\text{Fe}/\text{H}] \sim -1.5$  dex) than the OoII population (i.e.,  $[\text{Fe}/\text{H}] \sim -2.1$  dex). The metallicity of RRAB Lyrae tends to have a narrower distribution along the vertical distances away from the Galactic plane. This tendency can be explained in part by the fact that RRAB Lyrae in OoI clusters dominate at larger vertical distances, while a considerable fraction of these stars come from the OoII clusters (i.e.,  $\sim 22\%$  for the TNTS data) located at relatively smaller vertical distances (i.e.,  $|Z| \leq 20$  kpc). This dichotomy in metallicity distribution/Oosterhoff classification of RRAB Lyrae stars suggests that the Galactic stellar halo was perhaps formed by two distinct populations, favoring the dual-halo models (i.e., Miceli et al. 2008; An et al. 2013).



**Figure 10.** Distribution of the metallicity measured for the RRAB stars vs. vertical distances from the Galactic plane of the Milky Way. The blue, green, and gray points represent the TNTS, SDSS, and CSS data, respectively. The red and blue points in the inset panel represent the OoI and OoII subsamples of RRAB Lyrae from the TNTS.

#### 4.2.2. Eclipsing Binary Stars

Based on the variations of light curves, the eclipsing binaries can be classified into three broad categories: contact (EC), semi-detached (ES), and detached (ED) (Paczynski et al. 2006). The EC binary represents a typical W UMa system (EW), which shows a continuous change in brightness and has a depth between the primary and secondary minima that is almost equal or does not obviously vary. The ES binary represents the Beta Lyrae-type system (EB), which shows a shallower change in brightness during the out-of-eclipse phase than the EC binary and has a deeper primary depth than that of the secondary. The ED binary represents an Algol-type system (EA), which shows almost constant brightness during the out-of-eclipse phase. Among our sample, there are 313 ECs, 88 ESs, and 30 EDs. For the 299 new variables, one-third belong to the eclipsing binary stars.

There are 67 interesting EC or ES binaries that exhibit the O’Connell effect where the two successive out-of-eclipse maxima are unequally high in the light curves (O’Connell 1951; Nataf et al. 2010). Models such as captured circumstellar material (Liu & Yang 2003) and starspot activity (Bell et al. 1990) have been proposed to explain the O’Connell effect, but these explanations are still controversial. More photometric and spectroscopic observations are needed to understand the origin of such an effect seen in the eclipse binaries.

#### 4.2.3. Other Types of Variables

The long periodic variable stars include Miras, semi-regular pulsators, and slow irregular variables. Miras show large-amplitude variabilities (i.e.,  $>2.5$  mag in  $V$  band) with periods up to several hundred days. Semi-regular pulsators show diverse periodicities in a range from several days to several thousand days and the amplitude in light variations can vary from 0.01 mag to several magnitudes.<sup>10</sup> In our sample, 29 stars are Mira stars, 62 are semi-regular pulsators, and 11 are slow irregular variables.

<sup>10</sup> <http://www.sai.msu.su/gcvs/gcvs/iii/vartype.txt>

Observationally, Gamma Doradus stars show multiple periods ranging from 0.3 to 3 days with a typical variation amplitude of less than 0.1 mag (Cuypers et al. 2009). Based on our determinations of the periodicity, five new stars from our discovery (S070-6584, S598-3499, S619-728, SC49-6456, and SE61-2890) perhaps belong to this class. Delta Scuti variables are late A- or early F-type pulsating stars. Their pulsation periods are found to be in the range of 0.02–0.25 days (Breger 2000), such as star S585-17826 listed in Table 1. W Virginis (CWA) is a type II Cepheid; this type of star is regarded as the low-mass analogue of a Cepheid (Sterken & Jaschek 1996). SK39-3665 may be such a star in our sample.

## 5. CONCLUSION

We have presented the analysis of the unfiltered CCD images from the Tsinghua–NAOC transient survey. A total of 600 sky fields ( $\sim 1300$  square degrees), with repeated observations larger than 40 times during the period from 2012 October to 2014 August, were selected for the detection of variable stars. We carried out photometry using SExtractor. Fluxes of all other images are normalized to the reference ones; and the resulting unfiltered magnitudes were further calibrated using the  $R$ -band magnitudes of the PPMX catalog. We used SRPAstro, VARTOOLS, and PDM to create light curves and detect variable stars.

We detected a total of 1237 variables (including 299 new ones) with magnitude variations of  $\geq 0.1$  mag and periods ranging from 0.1 to 500 days. These variables include RR Lyrae, eclipsing binaries, semi-regular stars, etc. Among them 10 RRAB stars show the Blazhko effect and 64 eclipse binaries show the O’Connell effect.

We further examined the spatial distribution of the photometric metallicity for the RRAB stars and found that it shows a larger scatter near the Galactic plane (i.e.,  $-3 < [\text{Fe}/\text{H}] < 0$ ) but tends to converge at larger distances (i.e.,  $-2 < [\text{Fe}/\text{H}] < -1$ ). This variation can be explained by the fact that the OoI type of RRAB Lyrae stars has a higher metallicity and can extend out to a larger vertical distance (i.e.,  $|Z| \sim 45$  kpc), while the OoII type of RRAB stars has a lower metallicity and is relatively abundant at a vertical distance of  $|Z| \leq 20$  kpc. This result favors the idea that the Galactic halo was formed by two components with distinct properties.

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