

TRACE FISSION PRODUCT RATIOS FOR NUCLEAR FORENSICS
ATTRIBUTION OF WEAPONS-GRADE PLUTONIUM FROM FAST BREEDER
REACTOR BLANKETS

A Thesis
by
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ABSTRACT

A nuclear terrorist attack is one of the most serious threats to the national security of the United States, and in the wake of an attack, attribution of responsibility will be of the utmost importance. Plutonium, a by-product in spent nuclear reactor fuel, can be used in a nuclear weapon when obtained from reactor fuel discharged at a low burnup (1 MWd/kg). Characteristics of plutonium reprocessed from reactor fuel depend on factors such as the reactor type (thermal or fast reactor), fuel burnup, production history and the plutonium separation process used. Detailed understanding of the plutonium isotopic composition and fission product contaminant concentrations in separated plutonium would aid nuclear forensics activities aimed at source attribution in the case of interdicted smuggled plutonium, bolstering nuclear deterrence. The study presented here shows that trace fission product to plutonium ratios are amenable for nuclear forensics attribution. Through computational reactor core physics simulations, results are obtained for weapons-grade plutonium produced in a Fast Breeder Reactor (FBR). These fission product to plutonium ratios for the FBR are further compared with results reported elsewhere for a thermal Pressurized Heavy Water Reactor. This comparison of isotopic ratios results in substantial differences between fast and thermal neutron reactor systems, leading to the determination that a suite of selected isotopic ratios can attribute separated weapons-grade plutonium to a fast or thermal neutron source reactor system.

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NOMENCLATURE

CANDU	–	Canada Deuterium Uranium reactor
CSR	–	Control Safety Rod
DF	–	Decontamination Factor
DSR	–	Diverse Safety Rod
FBR	–	Fast Breeder Reactor
HFIR	–	High Flux Isotope Reactor
IAEA	–	International Atomic Energy Agency
ITWG	–	International Technical Working Group
LANL	–	Los Alamos National Laboratory
MCNP	–	Monte Carlo N-Particle
MOX	–	Mixed Oxide
MWd/kg	–	Megawatt-day per kilogram of heavy metal
MWe	–	Megawatt electric
NPT	–	Treaty on the Non-Proliferation of Nuclear Weapons
ORNL	–	Oak Ridge National Laboratory
PFBR	–	Prototype Fast Breeder Reactor
PHWR	–	Pressurized Heavy Water Reactor
PuO ₂	–	Plutonium Dioxide
PUREX	–	Plutonium Uranium Recovery by Extraction

PWR – Pressurized Water Reactor
RB – Radial Blanket
RDD – Radiological Dispersal Device
 UO_2 – Uranium Dioxide

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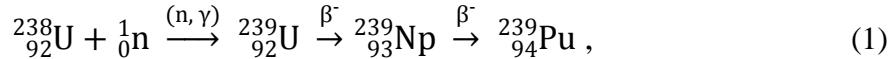
I. INTRODUCTION

The Nuclear Forensics and Attribution Act, signed into law by President Barack Obama in 2010, states that a nuclear terrorist attack is one of the most serious threats to the national security of the United States, and that in the wake of an attack, attribution of responsibility would be of the utmost importance.¹ Attribution begins with technical nuclear forensics, the process by which intercepted, illicit nuclear material is analyzed in order to identify the origin and source of the material.² Recognizing a threat illustrates the value of a robust forensics and attribution capability that results in establishing a credible nuclear deterrent.

There are two basic techniques for attributing reactor characteristics from results of analyzed material. These consist of a database method and an inverse analysis method. With the database method the measured results from material analyses are compared to a database of isotopic results to find the closest solution. In an inverse analysis method the data obtained from the measured material is used to develop model parameters for the forward model.³ The forward model being a reactor physics code. Both methods have their disadvantages. For example, a downside of the database method is that the database is never certain to be complete. Therefore, forensic analysis for attribution using a database method requires that data concerning foreign-origin materials be available and studied.⁴ However, it is possible that well validated computational models may be able to substitute data for cases where known material samples are not available.

I.A. Background and Motivation

Plutonium, a byproduct in spent nuclear fuel, is bred in nuclear reactor fuel by the conversion of uranium. The production of ^{239}Pu is through the neutron capture of ^{238}U :



where the half-lives for the beta decays of ^{239}U and ^{239}Np are 23.47 minutes and 2.355 days, respectively.⁵ As the fuel burnup increases, neutron capture reactions in ^{239}Pu and successive isotopes of plutonium lead to the production of higher mass number plutonium isotopes. Eventually a full composition of plutonium isotopes (^{238}Pu , ^{239}Pu , ^{240}Pu , ^{241}Pu , and ^{242}Pu) exist in the irradiated fuel. Fuel burnup is defined as the thermal energy produced per unit mass of fuel, expressed in Megawatt-day per kilogram of heavy metal (MWd/kg). The isotopic composition of the plutonium is dependent upon the fuel enrichment, the amount of fuel burnup, the nature of the reactor neutron energy spectrum to which the fuel is exposed, and the cooling time after the irradiation has occurred. Thus, the resulting composition of the discharged fuel may be able to provide information on the reactor system which produced the plutonium.

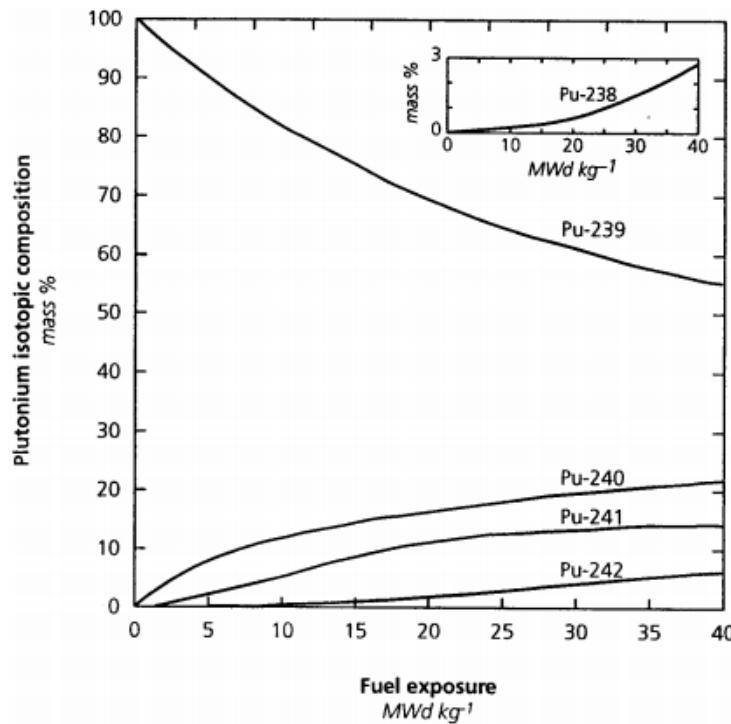


Figure 1. Plutonium isotopic composition as a function of fuel burnup in a PWR.

Fuel burnup is the major contributing factor to the plutonium isotopic composition. Figure 1, taken from reference 6, illustrates the change in plutonium composition as a function of fuel burnup in the case of a typical pressurized water reactor (PWR).⁶ It can be seen from Figure 1 that as the fuel burnup increases, the concentration of ²³⁹Pu, and thus the quality of the plutonium with respect to weapons usability, decreases. For levels of fuel burnup (7 to 45 MWd/kg) commonly achieved in power reactors, the resulting plutonium is reactor-grade rather than weapons-grade.⁶ However, a reactor could, in theory, be designed or misused to discharge fuel at a low burnup (1 to 2 MWd/kg) for the purpose of obtaining weapons-grade plutonium. If a reactor is to discharge uranium fuel subject to a low burnup of around 1 MWd/kg, the

plutonium produced will be weapons-grade, irrespective of the reactor type in which the uranium is irradiated. The characteristics of the plutonium composition defining plutonium grade as per reference 6 are listed in Table I.⁶

Table I

Approximate isotopic composition of various grades of plutonium.

Grade	Isotope (%)				
	²³⁸ Pu	²³⁹ Pu	²⁴⁰ Pu	²⁴¹ Pu	²⁴² Pu
Super-grade	-	98	2	-	-
Weapons-grade	0.12	93.8	5.8	0.25	0.022
Reactor-grade	1.3	60.3	24.3	9.1	5
MOX-grade	1.9	40.4	32.1	17.8	7.8
FBR blanket	-	96	4	-	-

From Table I, FBR (fast breeder reactor) blanket plutonium is estimated to be weapons-grade. This is because during the standard burn cycle of an FBR the blanket material situated at the periphery of the reactor core is exposed to a burnup around 1 MWd/kg, whereas it is unusual for a power reactor to discharge regular uranium fuel at that low of a burnup. Therefore, any country operating an FBR will be generating significant quantities of weapons-grade plutonium in the fuel blankets. Countries

including India and China currently have FBR programs where they are actively developing and operating FBRs for research and power production purposes.⁷

India's 500 MWe Prototype Fast Breeder Reactor (PFBR)⁸ is in the advanced stages of its construction, at the time of this writing, and will produce significant quantities of weapons-grade plutonium during normal operation. Previous work has estimated that about 140 kg of weapons-grade plutonium will be produced in the blankets of the PFBR each year.⁹ This capability is significant as India is a non-signatory of the Treaty on the Non-Proliferation of Nuclear Weapons (NPT). If India uses the PFBR and other reactors, only to meet civilian energy needs, international safeguards would provide confidence that plutonium is not being diverted from their power reactors. In 2010, the United States and India enacted a nuclear deal based on India's civilian and military nuclear separation plan which would allow access for international safeguards on Indian power reactors in exchange for India's ability to receive full civil nuclear energy cooperation, including fuel supplies for safeguarded nuclear power reactors.¹⁰ In accordance with the Indo-US 123 agreement, civilian power reactors will be placed under permanent International Atomic Energy Agency (IAEA) safeguards. However, India retains the sole right to determine whether nuclear facilities are civilian or military. Thus per the agreement, the PFBR as well as eight of their twenty-two, CANDU derived, Indian Pressurized Heavy Water Reactors (PHWRs) are exempt from IAEA safeguards.¹⁰ Both of these reactor types have the ability to produce weapons-grade plutonium, which can be used for the purpose of weapons production.

Operating in a non-safeguarded manner, the PHWRs are of particular interest because of their unique online refueling capabilities. Light water reactors undergo batch refueling which require a reactor shutdown to be performed, whereas heavy water reactors are frequently refueled while online and in operation. Their use of natural uranium fuel (0.72% ^{235}U), having a low reactivity, leads to typically refueling one fuel channel (eight fuel bundles ~100 kg U) per day. The PHWRs usually reach an average fuel burnup of about 6.7 MWd/kg.¹¹ However, the online refueling makes the reactors more susceptible to the diversion of material from the core, and this allows the potential for the fuel to be intentionally exposed to a low burnup and then removed from the core for purposes outside civilian energy production.

Due to the prevalent risk of misuse of these two types of reactors and the ability to produce weapons-grade plutonium from both, it is of interest to have detailed characterizations of weapons-grade plutonium produced by FBR and PHWR types. This research develops a suite of fission product to plutonium ratios for separated plutonium produced in the PFBR radial blanket, followed by a comparison of those ratios to material from a PHWR source reported elsewhere.¹² Detailed understanding of the characteristics, such as plutonium composition and fission product contaminant concentrations in separated plutonium would aid nuclear forensics activities which are aimed at source attribution in the case of interdicted smuggled plutonium in a pre-detonation state as well as to some extent for post-detonation analyses, bolstering nuclear deterrence.

I.B. Previous Work

Multiple studies on technical nuclear forensics research have been published which demonstrate the ability of isotopic data to retain information about the source of the produced special nuclear material. An environmental monitoring system was developed at Los Alamos National Laboratory (LANL) using fissionogenic noble gases, namely xenon and krypton, as a verification technique for reprocessing facilities.¹³ The relative concentrations of stable xenon and krypton isotopes depend on several reactor operating parameters. Measurements of isotopic ratios of these noble fission gases could thus be used to verify operator declarations or determine fuel parameters such as fuel type, burnup, and reactor type. The technique developed uses a high-precision mass spectrometer to measure stable noble gas isotope compositions in samples taken from a reprocessing plant exhaust stack. Selected isotopic ratios were determined and using sophisticated data analysis the ratios were compared to a database of isotopic ratios to infer the fuel parameters. The calculated database of xenon and krypton isotopic ratios was created using a series of reactor analysis codes to model essentially all types of reactors. The isotopic ratios were calculated as a function of burnup for pressurized water reactors, boiling water reactors, CANDUs, graphite moderated reactors, and FBRs. To verify the accuracy of the calculated database, the reactor models were benchmarked to a dozen measurements from literature. Computer modeling showed the ability of the isotopic ratios to distinguish between light-water reactors, heavy-water reactors, and breeder reactors. This system utilizes the fact that noble gases are not chemically bound to the fuel and are thus released during reprocessing.¹³ However, the isotopic ratios and

database developed by this technique are not useful when it comes to analyzing post-processed materials; which is the focus of this thesis study.

A thesis study was performed by M.R. Scott¹⁴ to create a forensics methodology for attributing spent fuel used in a radiological dispersal device (RDD) to a source reactor. The attributes which were determined included the spent fuel burnup, age from discharge, reactor type, and initial fuel enrichment. The attribution process was theorized to begin by accurately measuring the isotopic composition of the RDD debris with mass spectrometry. The method for finding the correct reactor type is a forward model, in which several reactor types are modeled with coupled MCNP (core physics analysis code) and ORIGEN (fuel burnup simulation code). The isotope measurements will be compared to the library of computationally calculated isotopes for different reactor types. An error will exist between the measured isotopes and computationally calculated isotopes. The reactor types will be ranked based on the lowest error as to which type is most likely to be the source reactor. To find a distinction between reactor types, isotopes with cross sections and yields that change significantly between reactor types were needed. Ratios of the chosen isotopes were plotted versus burnup for each reactor. The results showed that these isotopes could easily differentiate between a fast reactor and thermal reactor.¹⁴ The work, however, was focused on higher burnup levels which are more commonly achieved in power reactors, rather than low burnup weapons-grade plutonium. Isotopes coming from more complex production chains do not have sufficient time to be produced in low burnup material. As a result, the monitors

identified in Scott's study were not suitable for the nuclear forensics analysis of low-burnup fuel reported in the current study.

A previous study was completed by Wallenius et al. in which plutonium isotopes were analyzed and used for the purposes of origin determination of plutonium seized in the illicit trafficking of nuclear material.¹⁵ This study used the reactor fuel burnup simulation code, ORIGEN2, to calculate the plutonium composition for multiple reactors, as well as a thermal ionization mass spectrometer to measure plutonium isotope ratios of five plutonium samples. The sources of the plutonium samples included two from the National Bureau of Standards, two from the former Soviet Union, and one sample from an (ITWG) International Technical Working Group round-robin test. Following the measurements, a source reactor for each sample was then inferred by comparing the measured and computationally calculated isotopic compositions. Their study raises a number of concerns, though, with regards to the isotopics and computer modeling. The isotopic analyses consisted of plutonium and actinides only, with no investigation of contaminant fission products. The isotope generation and depletion code, ORIGEN2, uses a zero-dimensional fuel model giving the composition averaged over the whole reactor core. This could be problematic with an FBR, whose core consists of very different fuel and blanket regions. Additionally, irradiation times were considered to be continuous and no cooling time corrections were made to the material. Burnup levels obtained for the computational models were equivalent to typical burnup for each reactor with the exception of the FBR. Here, the FBR blanket material has a

burnup of 20 MWd/kg, which is relatively high. This burnup level is likely due to the averaging of core and blanket fuel, as a result of ORIGEN2 modeling.

A paper was published by A. Glaser¹⁶ over the isotopic signatures of weapons-grade plutonium produced in reactor types which have been historically used for plutonium production. The computer code system, MCODE, which links MCNP and ORIGEN2 was used to model three types of reactors: the Hanford-type, NRX-type, and Calder Hall-type reactors. For each of these reactor types, the plutonium composition was obtained and ratios of plutonium isotopes were analyzed. In addition, sample data, from previous publications, of the same plutonium isotopic ratios from more reactor types were included in Glaser's analysis. It was determined that predictive signatures derived from the plutonium isotopic ratios can distinguish weapons-grade plutonium from basic reactor types including fast breeder reactors, light water reactors using low-enriched fuel, and reactors fueled with natural uranium.¹⁶ Analyses of fission products are absent in the study done by Glaser. While a forensics methodology consisting purely of plutonium isotope ratios could be beneficial due to independence from the separation technique, the plutonium isotopes may be assisted by the inclusion of fission product compositions.

From the literature review, it was evident that previous studies have developed plutonium or fission product isotope analysis techniques which attribute nuclear material to a source reactor. Most of the research completed however, has been focused on reactor spent fuel, irradiated typically to an average burnup level, where the composition of plutonium produced is not of weapons-grade. Lacking are the investigations into

isotopic characteristics of fission product contaminants in weapons-grade material as a result of low burnup fuel from reactor misuse or a breeder reactor blanket. Using ratios of fission products to plutonium in separated weapons-grade plutonium, attempted in the present study, for nuclear forensics isotope analysis, is a novel approach for source reactor attribution.

II. THEORY

Technical nuclear forensics capabilities for the attribution of nuclear material to a source reactor are vital to establishing a credible nuclear deterrent. Previous nuclear forensics research has demonstrated the ability for isotopes within the material to retain information on the system which produced the material. With fast breeder reactors becoming more prevalent there is a need to continue research with a focus on FBR systems and weapons-grade plutonium from low burnup fuel. The theory involved with the creation and remaining presence of isotope characteristics in a material which possess information from the source system is discussed throughout this chapter.

A detailed characterization of weapons-grade plutonium will include fission product contaminants in addition to plutonium isotopes. The potential for trace amounts of fission products in separated plutonium is due to the non-ideal chemical process used to separate plutonium from irradiated nuclear fuel. The degree of purification achieved by a separation process can be quantified by decontamination factors (DF), which are the ratios of a stated impurity to desired component in the feed divided by the equivalent ratio in the product.¹⁷

$$DF = \frac{\left[\frac{Impurity}{Desired Component} \right]_{Feed}}{\left[\frac{Impurity}{Desired Component} \right]_{Product}} \quad (2)$$

The most commonly employed technique for plutonium separation is the Plutonium Uranium Recovery by EXtraction (PUREX) process.¹⁸ Using the PUREX process, decontamination factors of 10^6 - 10^8 have been achieved¹⁷ for the reduction of

fission products in separated plutonium, however a measurable contaminant concentration will remain. The DF is largely insensitive to mass-dependent fractionation. Thus, different mass isotopes of the same element will be separated together and equally during chemical separation.

II.A. Neutron Energy Spectrum

Similar to the plutonium composition, fission product inventories have a large dependency on parameters such as fuel burnup, the type of fuel, enrichment, the neutron energy spectrum of the reactor, and the fuel cooling time after irradiation. The objective of this research is to analyze plutonium and fission products from an FBR blanket fuel with a burnup of 1 MWd/kg, followed by a comparison to PHWR fuel also with 1 MWd/kg burnup. The fast neutron energy spectrum of a sodium cooled FBR is drastically different from the heavy-water moderated PHWR thermal neutron energy spectrum. For these reactor fuels exposed to equal levels of burnup, the neutron energy spectrum will be the source of variations in fission product inventories. The fuel used in PHWR is natural uranium (0.72 atom percent ^{235}U), whereas that used in an FBR blanket fuel is depleted uranium (0.25 atom percent ^{235}U). The energy production in the FBR core is from the seed fuel subassemblies containing mixed oxides (MOX) of PuO_2 and UO_2 . A plot of fast and thermal neutron energy spectra can be seen in Figure 2 which was obtained from reference 19.¹⁹

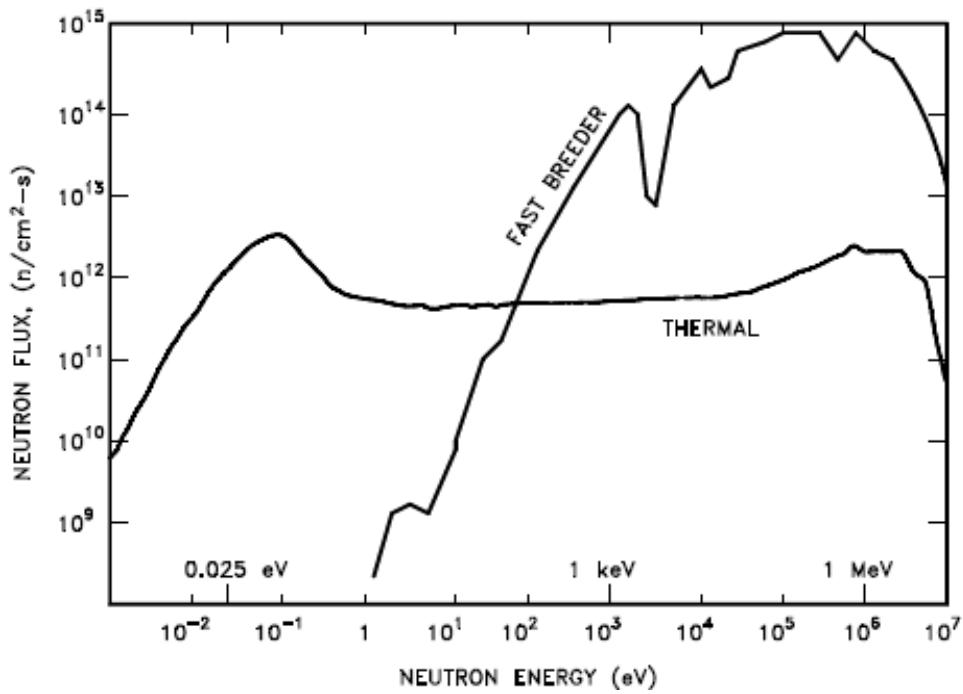


Figure 2. Neutron energy spectra for that of a fast reactor and thermal reactor.

The dissimilarity of the neutron spectra can lead to changes in characteristics of plutonium and fission product concentrations in the fuel. Variations in the fission yield for several fission products, variations in neutron interaction probabilities (cross-sections) for fission product isotopes, and variations in the neutron interaction probabilities for the plutonium isotopes can all lead to characteristic differences in fission product to plutonium ratios, which is the focus of this study. Figure 3 displays the fission product yield curves for the thermal fission of ^{235}U to represent a PHWR, and for the fast fission of ^{238}U and ^{239}Pu which represents an FBR.²⁰ It can be seen from Figure 3 that the fission product yields are, in fact, different between the neutron energy and uranium and plutonium isotopes and thus analysis of fission product inventories, a small

portion of which will be contaminant in the separated plutonium, can assist in attributing plutonium to a source reactor. It should be noted that energy for the fast fission yield in Figure 3 is at a neutron energy of 500 keV. This is a slightly higher energy than the energy of the dominant neutrons in a MOX-fueled FBR, which varies between 100 keV to 400 keV.

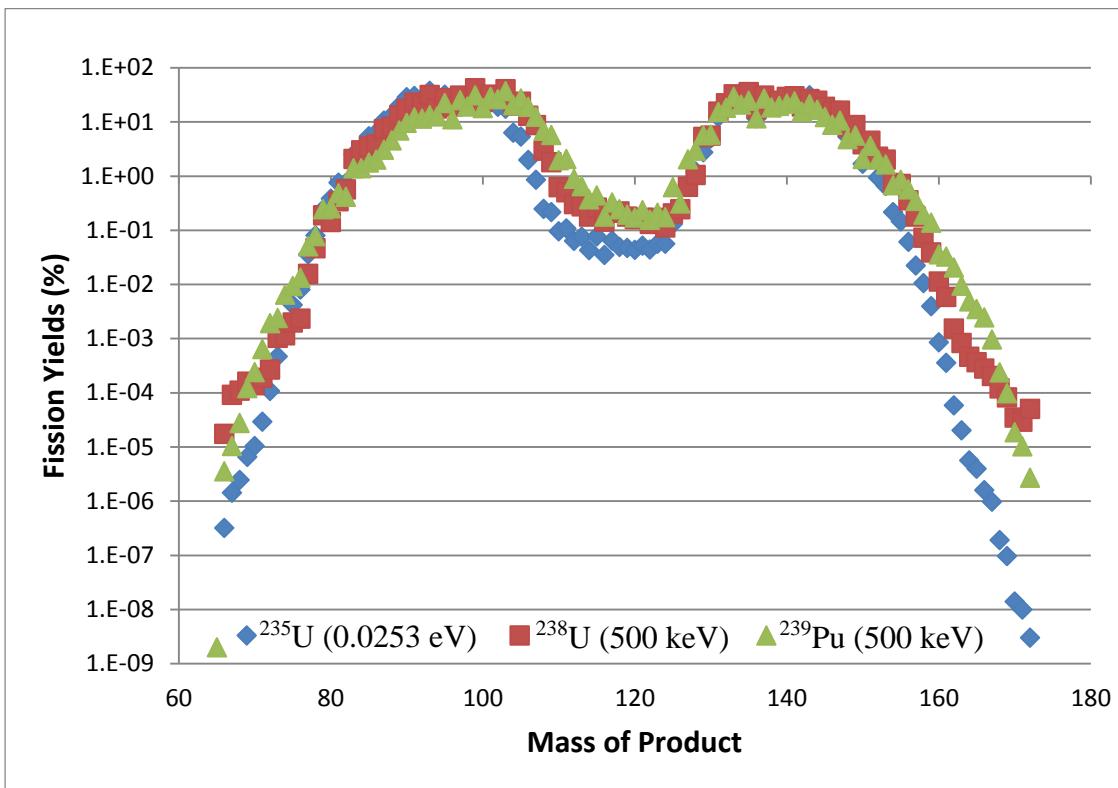


Figure 3. Cumulative fission yield curves for various isotopes and neutron energies.

II.B. India's Three Stage Nuclear Program

Dr. Homi J. Bhabha is regarded as the “father” of India’s nuclear program. Dr. Bhabha conceived the idea of a three stage nuclear program as a way for India to work

around the country's limited amount of uranium while utilizing its vast thorium reserves.

The strategy is to develop a Th – U fuel cycle, breeding fissile ^{233}U from the fertile ^{232}Th . Stage one consists of natural-uranium fueled thermal reactors (PHWRs) used to generate power. The second stage of the program involves a fleet of fast breeder reactors. The stage two fast breeder reactors, beginning with the PFBR, will be fueled with reactor-grade plutonium and depleted uranium from the reprocessed spent fuel of stage one reactors and will breed more plutonium than it consumes. With enough plutonium fuel stockpiled, the FBRs can then switch to a plutonium – thorium breeder cycle capable of producing ^{233}U . The third stage will then be a $^{233}\text{U} – \text{Th}$ breeder cycle, which best utilizes India's resources.²¹ Figure 4 gives a visual representation of India's three stage nuclear power program.

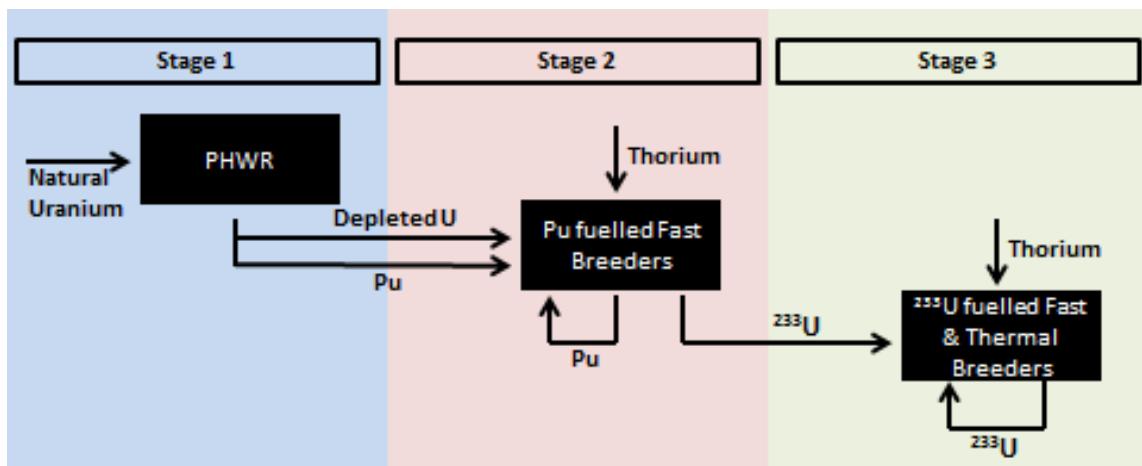


Figure 4. India's three stage nuclear power program.

II.C. The Indian PFBR

The Indian PFBR is the specific fast breeder reactor modeled in this study due to core characteristics being available in open literature.²² However, the objective of this thesis study is to analyze a fast breeder reactor system and the results obtained are applicable to general FBR systems. Design information for the 500-MWe Indian PFBR was obtained from Chirayath et al.²² and essential parameters are listed in Table II. There is an active core, one meter in height, which consists of an inner core and outer core of MOX “driver” fuel. The MOX fuel of the active core has PuO₂ enrichments of 20.7% for the inner core and 27.7% for the outer core. This increases the amount of fissile material around the radial periphery of the active core, thus desirably flattening the neutron flux profile across the reactor. Axial blankets of length 0.3 m each sit above and below the active core, all of which are surrounded by 1.6 m tall radial blankets. Both core blankets are comprised of depleted UO₂ “target” fuel, capturing the neutrons leaking from the core. This large amount of ²³⁸U is where the weapons-grade plutonium breeding will take place, due to the low fuel burnup experienced in the blanket regions. The plutonium bred in the axial and radial blankets are likely to have similar characteristics; however, this project is focused on the plutonium produced in the radial blankets only. A core map of the PFBR, in Figure 5, highlights the inner core, control safety rods (CSR), diverse safety rods (DSR), outer core, and radial blanket.

Table II
PFBR core design parameters.

Core parameter	Value
Reactor power (MWe)	500
Thermal Efficiency (%)	40
Maximum linear heat rating (W/cm)	450
Fuel pin clad O.D./I.D. (cm)	0.66/0.57
Fuel pellet diameter (cm)	0.555
Fuel pins per sub-assembly	217
Fuel pin triangular pitch (cm)	0.825
Assembly pitch (cm)	13.5
Radial blanket pin clad O.D./I.D. (cm)	1.433/1.323
Radial blanket pellet diameter (cm)	1.29
Pins per radial blanket sub-assembly	61
Radial blanket pin triangular pitch (cm)	1.553
Fuel assembly sheath thickness and sub-assembly size (cm)	0.32/13.13
Active core height	100
Axial blanket height top + bottom (cm)	30 + 30
Radial blanket height (cm)	160
Fuel – Density of fuel (g/cm ³)	PuO ₂ – UO ₂ ^a (11.0)
Axial/Radial blanket fuel	Depleted UO ₂ ^a
Fuel clad material	20% CW D9 steel
Core Pu enrichments in MOX, inner core and outer core (%)	20.7/27.7
Plutonium isotope ratios in fuel: ²³⁹ Pu/ ²⁴⁰ Pu/ ²⁴¹ Pu/ ²⁴² Pu (%)	[PuO ₂ / (PuO ₂ + UO ₂)]
Plutonium inventory (t)	68.8/24.6/5.3/1.3
Primary coolant	1.99
Primary inlet/outlet temperature (°C)	Liquid sodium
Fuel average temperature (°C)	397/547
Fuel cycle (effective full power days)	1289
	180

^a Depleted uranium: ²³⁵U (0.25 at.%) and ²³⁸U (99.75 at.%)

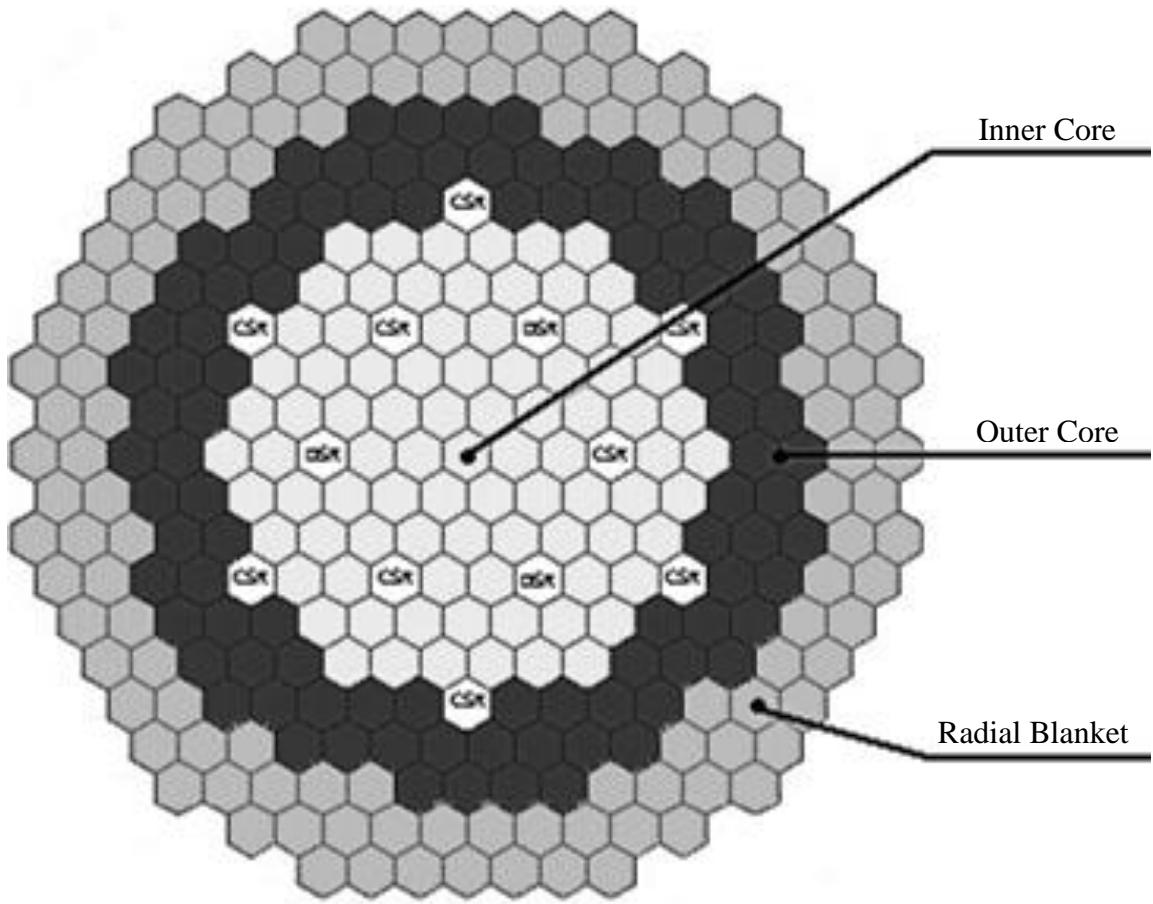


Figure 5. Core map of the PFBR.

II.D. Monte Carlo N-Particle (MCNP) Transport Code

MCNPX Version 2.7, a general-purpose Monte Carlo radiation transport code designed to track particle interactions, was utilized for modeling the PFBR core.²³ A user created input deck contains information regarding geometry specifications, material descriptions, selection of neutron interaction cross-section evaluations, location and characteristics of the radiation particle source, type of output information desired, and any variance reduction techniques if applicable.²⁴ The geometry of MCNPX treats an

arbitrary three-dimensional configuration in geometric cells bounded by first-, second-, and fourth-degree surfaces. The cells are defined by the intersections, unions, and complements of the regions bounded by the surface. The variety of flux estimators available, including flux averaged over a volume, flux averaged over a surface, and flux at a point, make MCNPX a versatile and powerful code for radiation transport calculations.²³ MCNP will read variables of the user input file and perform appropriate particle simulations.

The Monte Carlo method is a numerical analysis technique, which uses random sampling procedures to construct the solution of a physical or mathematical problem. A stochastic model is set up, and by sampling from appropriate probability distributions, estimates the required numerical answers to the problem by statistical means.²⁵ Monte Carlo thus duplicates the process of nuclear particle interactions with matter by sampling, via random numbers, probability distributions calculated from transport data.²⁴

A Monte Carlo code must have a supply of truly random numbers, which are uniformly distributed between 0 and 1. Each particle is followed from its birth to its death or escape from the system, with random sampling of probability distributions contained in radiation transport equations to determine the outcome at each step of its life.²⁴ In the life of a particle, events may include distance traveled to the next collision, collision nuclide selection, and nuclear reaction selection for a nuclide. Probabilities representing the physical process, at each event, are calculated based on physics, transport data and the materials involved. A random number is selected at an event and applied to the probability distribution to determine the events outcome. This process is

repeated along the particle's history, or path from birth until death, with a particle's death coming from absorption or leakage from the system. As a large number of histories are tracked, the average particle behavior more accurately simulates the physical process.²⁴

The Monte Carlo method indirectly solves the Boltzmann Transport Equation by simulating individual particles and averaging the behavior of a large number of particles. Monte Carlo is well suited for solving complicated three-dimensional, time-dependent problems²⁴ and thus is ideal for use in this thesis work.

II.E. CINDER90 Depletion

During the operation of a nuclear reactor the material compositions undergo changes, mainly, due to nuclear reactions caused by neutron interactions. The neutron interaction probabilities of the nuclides, along with the neutron flux, dictate the rate of change of the isotopic composition. As the isotopic compositions of the materials change, the nuclear reaction rates of the system are further altered. It is therefore essential to calculate the isotopic changes of the materials in the reactor in order to accurately simulate the reactor's operation.²⁶

MCNPX comes equipped with the depletion/burnup code, CINDER90, built-in to the code package.²³ An advantage of this integration is that only a single input is needed to run both the transport calculations as well as the isotope generation and depletion calculations. MCNPX runs a steady-state calculation to determine 63-group neutron fluxes, which are then energy-integrated with nuclide transport cross-sections resulting in reaction rates. CINDER90 takes the MCNPX-generated data, including reaction rates,

collapses the 63-group cross-sections into a one-group cross-section and performs the depletion calculation to obtain new isotopic compositions for the next time step. This process is repeated for each burnup time step, specified by the user, until the entire core burnup simulation is completed.²⁷

Solving for the time-dependent change in an isotope composition requires accounting for nuclear reactions which causes a production or loss of the nuclide, and may be described by the Bateman equations.²⁶ A simplified form of the Bateman equations for a specified isotope is^{26,27}

$$\frac{dN_i}{dt} = -N_i(t)\beta_i + \bar{Y}_i + \sum N_k(t)\gamma_{k \rightarrow i} \quad (3)$$

where:

$\frac{dN_i}{dt}$ = time-dependent change in isotope i,

$N_i(t)$ = the time-dependent atom density of isotope i,

β_i = the total transmutation probability of isotope i,

\bar{Y}_i = production of isotope i via an external source, and

$\gamma_{k \rightarrow i}$ = the probability of an isotope k transmuting, by decay or absorption, into isotope i.

Equation 3 is nonlinear because the transmutation probabilities rely on the time-integrated flux, which is also reliant upon the time-dependent isotope compositions.²⁷ To make the equation linear, the assumption must be made that the transmutation probabilities remain constant over the time step. This assumption thus requires attention when selecting the number and duration of time steps. A larger number of steps means more computational run time, however a time step of a long duration may not be

sufficiently accurate in accounting for composition transmutations. In CINDER90, the set of coupled differential equations is reduced to a set of linear differential equations using the Markov Linear Chain method.²⁶ Linear chains are created for each isotope transmutation path, starting from the initial concentrations:

$$\frac{dN_i}{dt} = -N_i(t)\beta_i + \bar{Y}_i + N_{i-1}(t)\gamma_{i-1} \quad (4)$$

where γ_{i-1} is the transmutation probability of forming isotope N_i .²⁷ The solution to each linear chain determines a partial isotope composition, which is then summed to obtain the total isotope inventory. Because of the use of these linear chains, the isotopic inventory is only coupled to preceding elements in the sequence, where the parameters are assumed known.²⁶ The general solution to such a linear sequence is as follows:^{26,27}

$$N_n(t) = \prod_{k=1}^{n-1} \gamma_k \left\{ \bar{Y}_m \left[\frac{1}{\prod_{l=1}^n \beta_l} - \sum_{j=1}^n \frac{e^{-\beta_j t}}{\prod_{i=1, \neq j}^n (\beta_i - \beta_j)} \right] + N_1^0 \sum_{j=1}^n \frac{e^{-\beta_j t}}{\prod_{i=1, \neq j}^n (\beta_i - \beta_j)} \right\} \quad (5)$$

II.F. ORIGEN & Matrix Exponential Method

Multiple depletion codes exist with a commonly used system being ORIGEN2.2.²⁸ The point-depletion and radioactive-decay code can simulate nuclear fuel cycles and calculate nuclide compositions and characteristics of materials contained in the nuclear fuel.²⁸ Both ORIGEN2.2 and CINDER90 are zero dimensional depletion codes that can be linked to MCNP flux generation, but the two codes solve the depletion equation differently. CINDER90 utilizes the Markov Linear Chain Method whereas ORIGEN2.2 is based on the Matrix Exponential Method. The basic equation ORIGEN solves is the first order differential equation for nuclide decay. The solution method will be a matrix of first order differential equations of size $N \times N$, with N being the number

of nuclides. The large sparse matrix requires a large amount of memory in order to store all the necessary computations.²⁶ Thus, to accelerate calculations the number of nuclides followed in the set of equations is limited. ORIGEN2.2 has the ability to track 1700 isotopes, while the CINDER90 library contains 3456 isotopes. Due to the current research work's interest in predicting and analyzing quantities for a vast number of fission product isotopes, CINDER90 was chosen as the ideal isotope generation and depletion code for use.

III. PFBR CORE MODELING AND FUEL BURNUP SIMULATIONS

III.A. PFBR Model & Simulation

In simulating the operation of the PFBR, a detailed pin-by-pin 3D model of the PFBR core was created with MCNPX-2.7. An input deck is attached in Appendix A. As stated previously, the PFBR core is comprised of inner core sub-assemblies, outer core sub-assemblies, and radial blanket (RB) sub-assemblies. The inner and outer core sub-assemblies are geometrically identical with varying fuel enrichments. Cross-sectional views of the fuel pin arrangements for a core sub-assembly and a radial blanket sub-assembly, as generated by MCNP, can be seen in Figure 6.

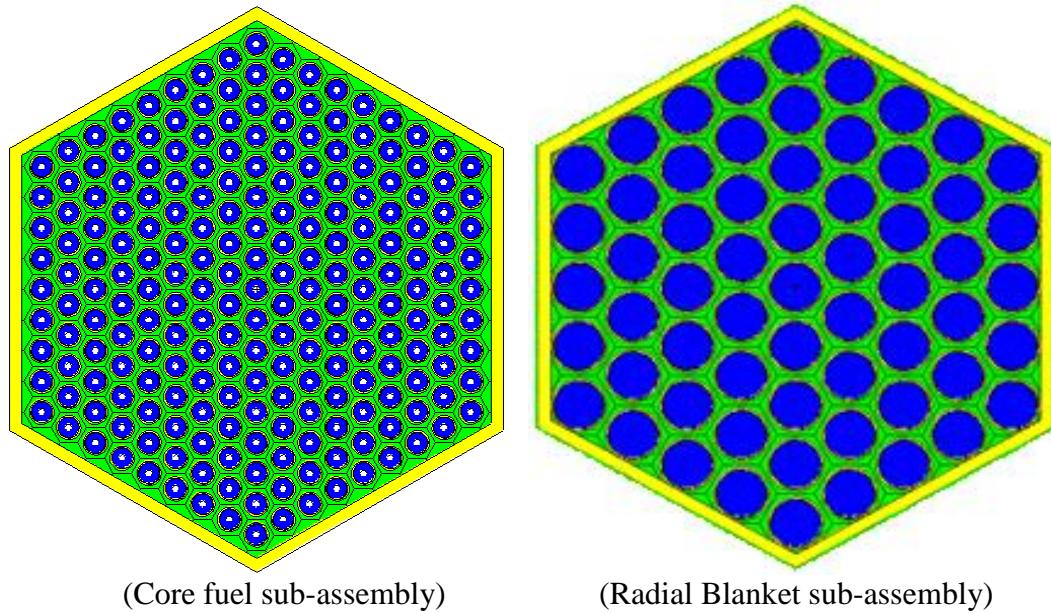


Figure 6. Fuel pin arrangement for core and radial blanket sub-assemblies.

To realistically simulate reactor operations a separate input deck had to be developed for every refueling cycle, which consisted of 180 days full-power operation followed by 60 days of refueling shutdown and decay. The first input deck was built for the fresh PFBR core and two input decks were built, each having material comprised of the predecessor's output, to accurately simulate the refueling and altering of the core as it reached an equilibrium core configuration in the third fuel cycle. When the equilibrium core configuration is reached, the subsequent cycles will continue to have the same core configuration as the equilibrium core. The MCNPX simulations are computer intensive. The run time for each simulation was just above 30 days while running in parallel on a 32 core, 2.7 GHz, 64 GB RAM machine. These computer simulations provided an estimate of the plutonium composition and fission product composition within the discharged PFBR blanket fuel.

MCNPX output data was analyzed to perform checks on the accuracy of the model simulations. The global parameter, neutron multiplication factor (k_{eff}), was obtained for each step and plotted versus time. The values of k_{eff} as a function of time for the first three cycles of operation, shown in Figure 7, were compared with expected values²² and were seen to behave as anticipated. The value for k_{eff} is larger than 1.0 because the lack of control rods in the model. As can be seen in Figure 5 the reactor control rods are not in close proximity to the radial blanket, and therefore the lack of control rods in the model simulations will not have an impact on the radial blanket results. The neutron energy spectrum calculated by MCNPX was used as another check to the model. Figure 8 displays the neutron spectra in the inner core, outer core, and

radial blanket regions. The shape of the neutron spectra are as expected and matched the values for the average neutron flux throughout the PFBR found in Chirayath et al.²² The dominant neutron energy in the inner and outer core can be seen from Figure 8 as between 100 keV to 400 keV. The radial blanket spectrum is slightly softer than the active core, indicative of the small thermalization from the oxygen of the MOX fuel as the neutrons' travel length increases.

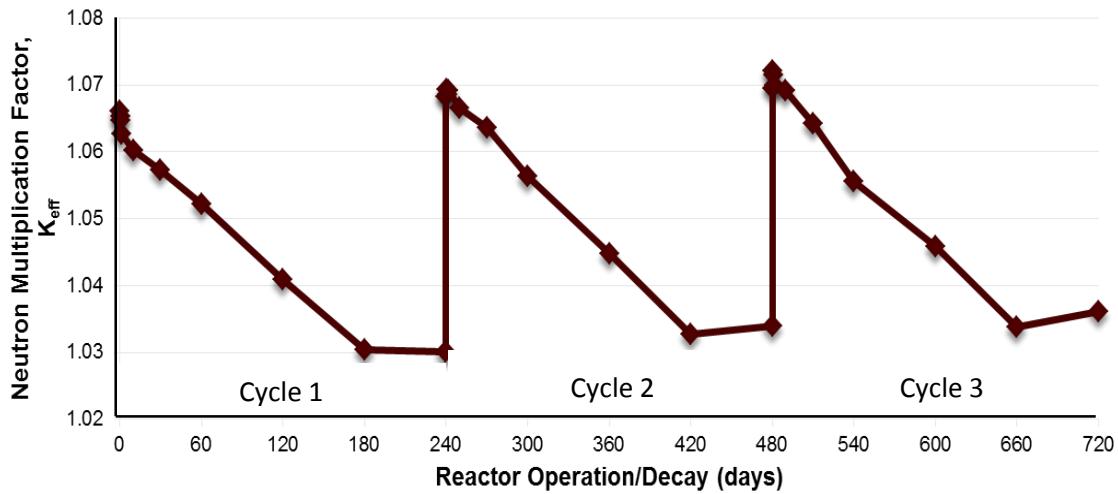


Figure 7. Variation of k_{eff} as a function of reactor operation.

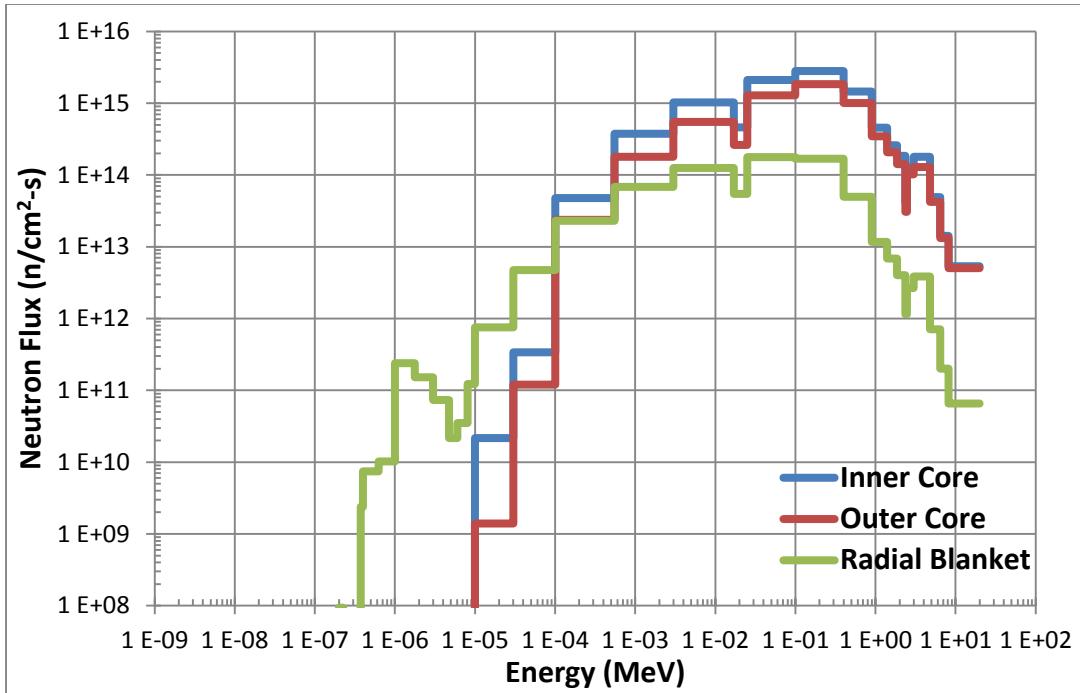


Figure 8. MCNPX-generated neutron energy spectra for the Indian PFBR core.

The standard operating scheme of the PFBR will be cycles of 180 days full-power operation followed by 60 days shutdown and refueling, with one-third of the active core being refueled at the end of each cycle.²² Conversely, the radial blanket sub-assemblies are refueled with a slightly different pattern, yet still on the basis of fuel burnup. Typical operation of the PFBR will be to discharge the radial blanket sub-assemblies during the refueling stage in which the burnup is nearest to 1 MWd/kg. Thus the radial blanket is split into three sections. Forty-two blanket sub-assemblies, which are in close proximity to the core, and therefore exposed to a large neutron flux, are replaced after every cycle of 180 days. Six radial blanket sub-assemblies are refueled after every two cycles, and seventy-two blanket sub-assemblies that are located farther

from the core are irradiated for three cycles before being refueled. Figure 9 displays a core map of the PFBR, highlighting the three radial blanket groups. The radial blanket sub-assemblies which are replaced after a single cycle are shown in red, two cycles in yellow, and the sub-assemblies which are irradiated for three cycles before being replaced are shown in blue. Table III gives the fuel burnup estimated through MCNPX simulations for each radial blanket group at the end of life. It can be seen that the group which is exposed to three irradiation cycles reaches a burnup closest to the target of 1 MWd/kg. Therefore the “three cycles” group of radial blanket sub-assemblies is the material of reference for the analyses and isotope selection described later.

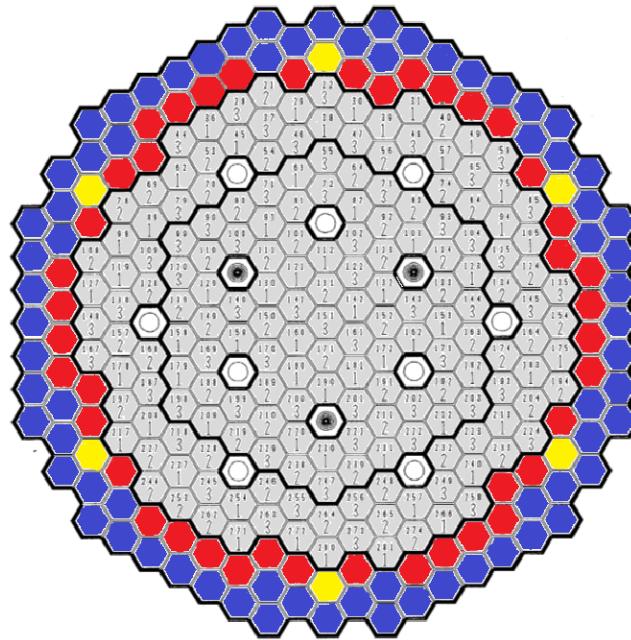


Figure 9. Core map highlighting the three radial blanket refueling groups.

Table III
End of life burnup level for radial blanket groups.

Irradiation Time & No. of RB Assemblies Discharged	Burnup (MWd/kg)
One Cycle (42 RBs)	0.710
Two Cycles (6 RBs)	1.302
Three Cycles (72 RBs)	1.016

III.B. Radial Splitting

As mentioned in the previous chapter, CINDER90 is a zero-dimensional depletion code and therefore has no knowledge of the spatial dependence of the transmutation rates.²⁶ CINDER90 assumes the flux in each burned material is not spatially dependent throughout that material. This averaging of the flux over each material region may affect the isotope concentrations if the spatial dependence of the flux would result in overall transmutation rates significantly different from the average. If necessary, splitting the material region can reduce the magnitude and the effects of the average flux.

The cylindrical fuel pins of the PFBR radial blanket can be split both radially and axially. Radially, transmutation rates may change with the flux and neutron properties as the neutron travels through the fuel. However, the fast neutrons of the PFBR are at high energies. Analytic calculations showed that the mean free path of neutrons at energies common in the PFBR are more than five times larger than the diameter of the radial

blanket fuel pins. Thus, there is no need to radially split the fuel pins due to the fast neutron spectrum.

III.C. Axial Splitting & Super Cell

Figure 10 shows the neutron flux profile as a function of core height for the outer core and the radial blanket groups obtained by running 20E+6 histories (or 20 Megahistories). It is clear that the neutron flux is not constant along the height of the fuel pin, which may affect the final isotopic concentrations from inaccurate burn and production as a result of the difference between the local flux and average flux. To test if, and how many, axial regions are needed, a case should be modeled with increasing axial regions until the resultant data converges. A full-core cycle simulation is too computer intensive, though, to run multiple times. Consequently, a super cell was modeled to accurately represent the full core, while at the same time being small and simple. An MCNPX-generated image of the super cell can be seen in Figure 11.

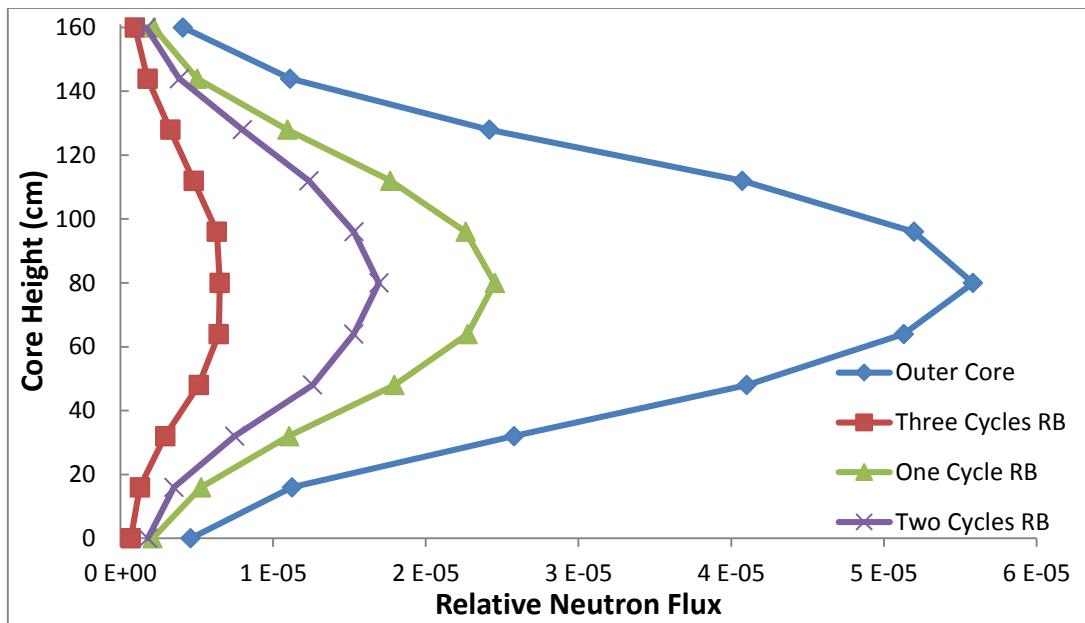


Figure 10. Flux profile as a function of pin height.

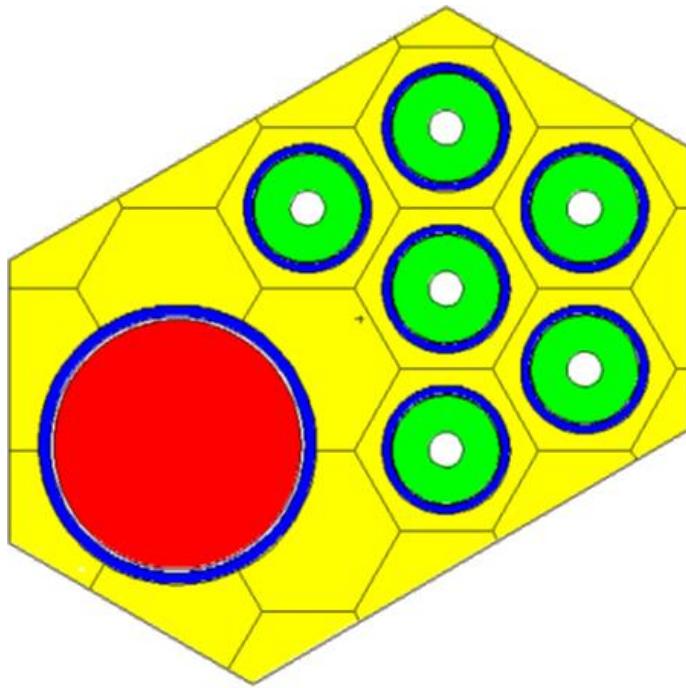


Figure 11. MCNP-generated image of the super cell model.

A super cell was used to determine the appropriate number of axial regions to split the radial blanket pins into. The small and simple super cell model has a short run time and is therefore conducive to several replications. The super cell consists of six outer core fuel pins at the same pitch found in the full core, and a radial blanket fuel pin appropriately placed from the outer core fuel pins and the boundary. The six faces comprising the elongated hexagon have a reflecting boundary condition. This design resulted in a neutron multiplication factor, k_{eff} , for the system of about one. The power of the super cell was set to achieve the same neutron flux as in the full core, and the radial blanket material was burned to the target, 1 MWd/kg.

Simulation of the super cell was repeated four times with 1, 3, 5, and 7 axial regions. For each simulation, the amounts of ^{137}Cs and total plutonium were summed up for the entire radial blanket pin. The results from the super cell simulations are shown in Table IV. For the mass of plutonium the relative error was less than 0.1%, and for the masses of ^{134}Cs , ^{137}Cs , ^{148}Nd , and ^{150}Sm the relative errors were all around 2%. This error was deemed insignificant, and a single axial region was used in the subsequent full core simulations. Thus, for simplicity of the model and shorter computational run time, the blanket fuel pins were not split axially.

Table IV
Super cell simulation results for axial splitting.

Number of Axial Regions	^{134}Cs mass (g)	^{137}Cs mass (g)	^{148}Nd mass (g)	^{150}Sm mass (g)	Plutonium mass (g)
1	1.294 E-4	3.827 E-2	1.374 E-2	2.033 E-4	6.726
3	1.336 E-4	3.843 E-2	1.367 E-2	1.906 E-4	6.716
5	1.356 E-4	3.842 E-2	1.366 E-2	1.921 E-4	6.719
7	1.372 E-4	3.848 E-2	1.368 E-2	1.943 E-4	6.716

III.D. MODEL VERIFICATION

After verifying the neutron energy spectrum and time-dependent multiplication factor (k_{eff}) discussed in section III.A, further model verifications were made. To verify the reactor simulation, the production behaviors of a few select isotopes of interest were analyzed. The isotope mass within each radial blanket (RB) region was plotted as a function of fuel burnup and compared to expected behavior. Mass buildup of ^{137}Cs , ^{148}Nd , ^{134}Cs , and ^{239}Pu are displayed in Figures 12, 13, 14, and 15, respectively.

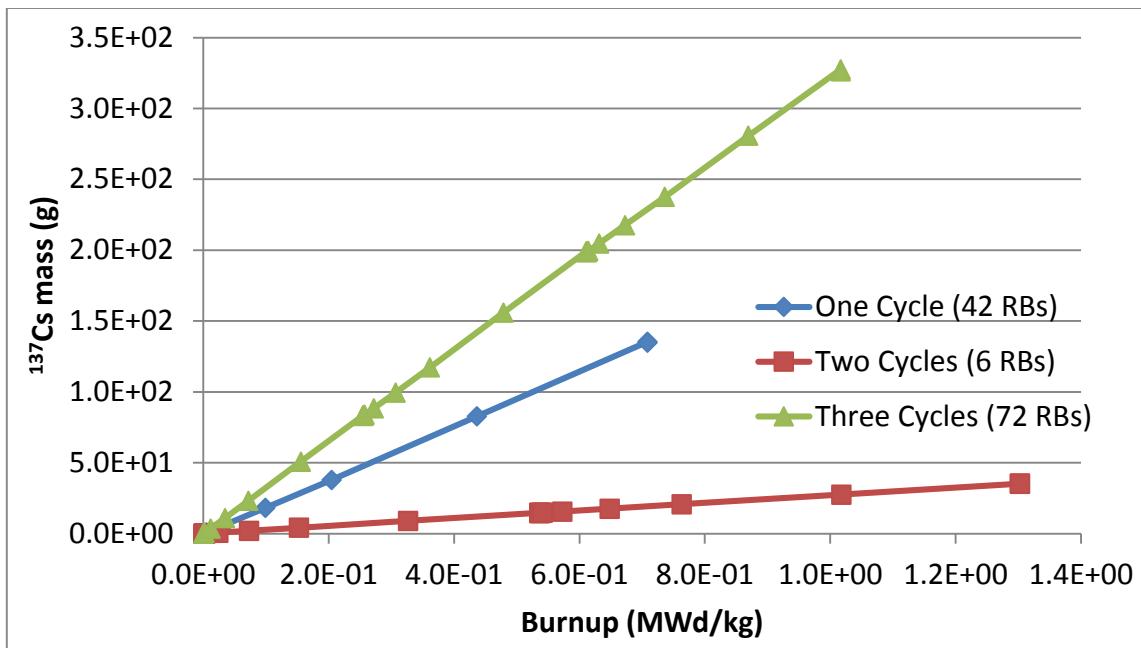


Figure 12. Production of ^{137}Cs as a function of burnup in radial blanket regions.

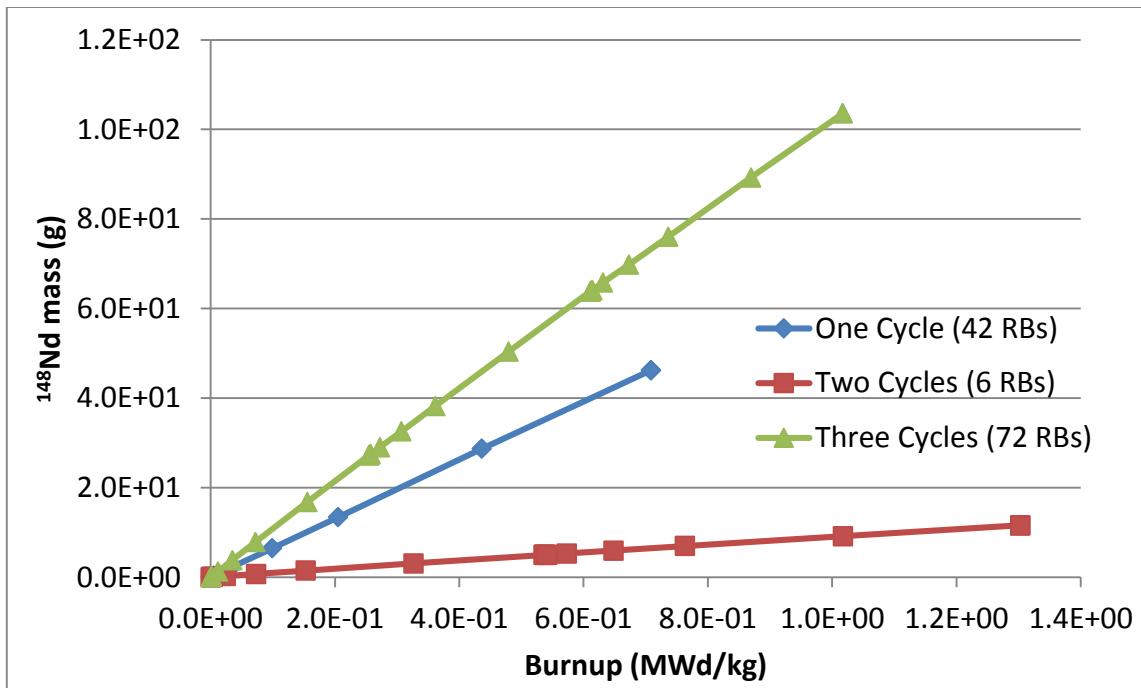


Figure 13. Production of ^{148}Nd as a function of burnup in radial blanket regions.

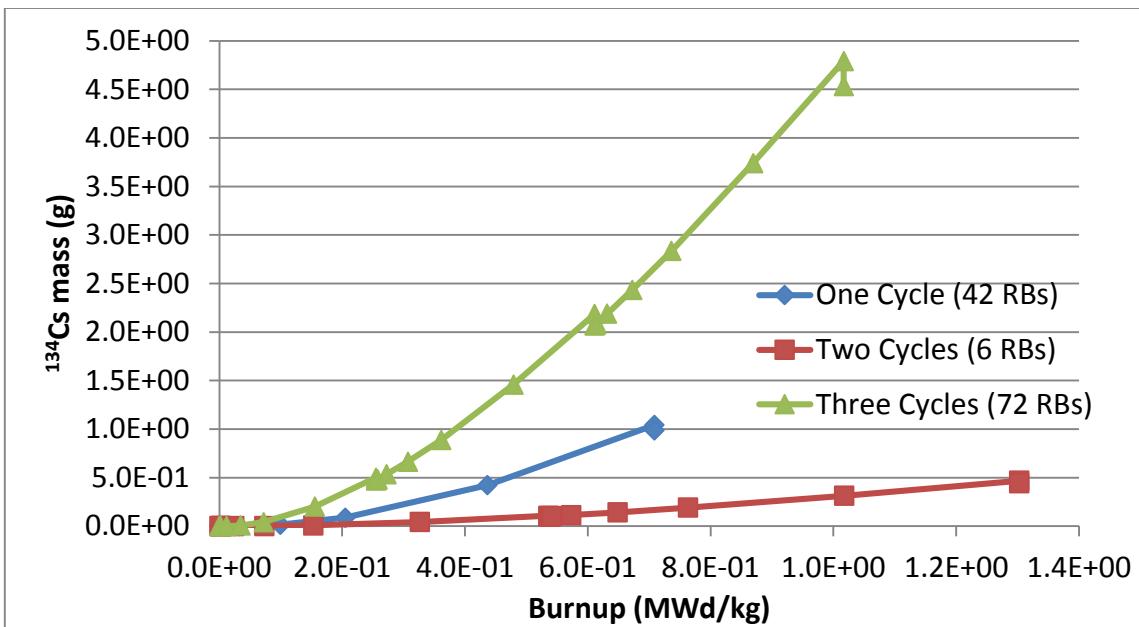


Figure 14. Production of ^{134}Cs as a function of burnup in radial blanket regions.

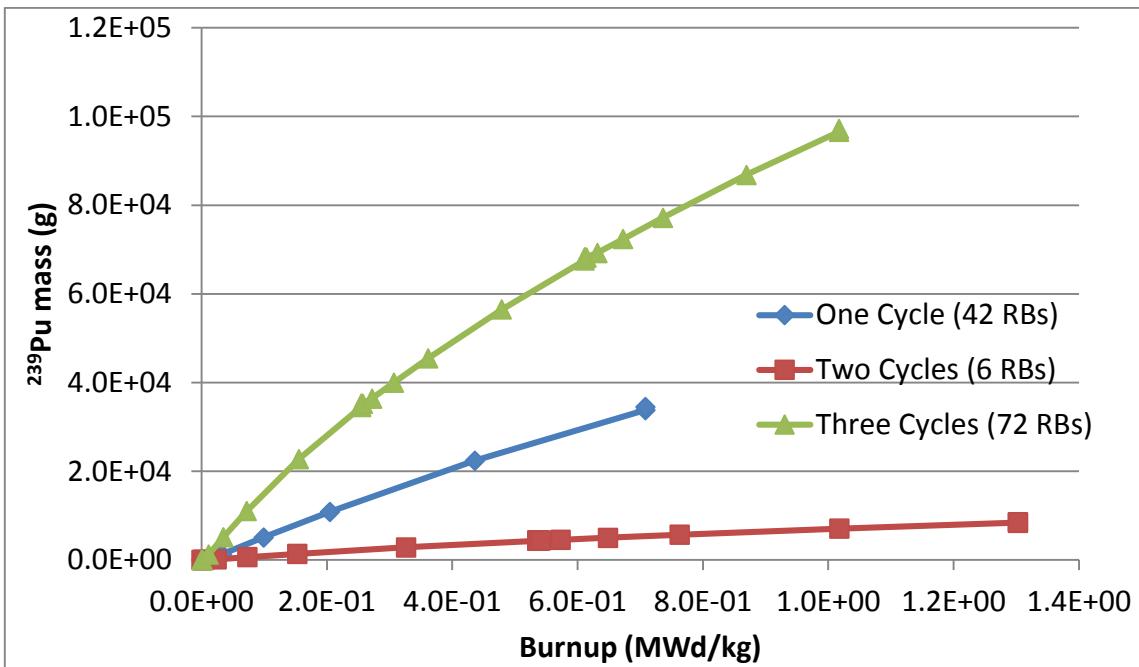


Figure 15. Production of ^{239}Pu as a function of burnup in radial blanket regions.

For each isotope, the isotope production versus material burnup is plotted for the three radial blanket regions. The difference in the rates of production for the three radial blanket regions is due to the mass of fuel, or number of radial blanket sub-assemblies. Forty-two, six and seventy-two radial blanket sub-assemblies are refueled in one cycle, two cycles and three cycles respectively. For all four isotopes, when normalized by fuel mass in each region, the lines for isotope mass as a function of burnup all fall on top of one another, irrespective of the radial blanket region refueled. Both ^{137}Cs and ^{148}Nd are expected to increase linearly with burnup and are common burnup monitors.²⁹ Both isotopes are direct fission products, and have negligible neutron cross-sections for loss mechanisms.³⁰ During the 60-day shutdown periods no significant decay will occur as the half-life of ^{137}Cs is about 30.1 years and ^{148}Nd is stable. Two production methods exist for ^{134}Cs , as a direct fission product and from the neutron capture of the fission product ^{133}Cs . The neutron cross-sections for loss mechanisms of ^{134}Cs are insignificant.³⁰ The production of ^{134}Cs is, therefore, proportional to the square of the neutron flux.²⁹ With a short half-life of around 2.06 years, an observable amount of ^{134}Cs decays away during the 60-day shutdowns. The production of ^{239}Pu , as seen in equation 1, is nearly linear at low levels of fuel burnup. However, due to the cross-sections for radiative capture and fission the amount of ^{239}Pu approaches an asymptotic composition.³⁰ Figures 12 through 15 match the expected production behavior of the four selected isotopes, indicating the model is correctly simulating reality.

IV. RESULTS AND DISCUSSIONS

IV.A. Plutonium Production

The significant quantity (SQ), defined as the approximate amount of nuclear material for which the possibility of manufacturing a nuclear explosive device cannot be excluded, is 8 kg for plutonium.³¹ Estimations for the amount of plutonium produced, as well as isotopic composition, in the PFBR radial blanket were obtained using MCNPX-2.7 burnup simulations. Results from the simulations for the three radial blanket regions are shown in Table V. Model simulations show that the PFBR does, in fact, produce high quality plutonium on a large scale. Based on the concentrations of ^{239}Pu , ^{240}Pu , and the grade definitions found in Table I, the resulting plutonium is categorized as super-grade. With three radial blanket regions and their various refueling frequencies, a refueling pattern is repeated every six cycles. Table VI shows the pattern and the mass of plutonium discharged at the end of each refueling cycle. Six full refueling cycles equals 1440 days, or almost four years. In this period, 432.9 kg of plutonium will be discharged from the radial blanket regions. On average, standard operation of PFBR will yield over 100 kg per year of super-grade, weapons-usable plutonium, from the radial blankets alone.

Table V

Plutonium produced in the PFBR radial blanket regions at the end of the regions' life.

Irradiation Time (Number of Cycles)	No. of Assemblies	Burnup (MWd/kg)	Total Mass of Plutonium (kg)	Plutonium Isotopes (%)				
				^{238}Pu	^{239}Pu	^{240}Pu	^{241}Pu	^{242}Pu
3	72	1.0157	99.06	0.0052	98.05	1.907	0.0382	0.00028
2	6	1.3023	8.664	0.0074	98.24	1.729	0.0236	0.00016
1	42	0.7096	34.85	0.0047	99.01	0.975	0.0073	0.00003

Table VI

Plutonium yield from the PFBR radial blanket regions.

Cycle Number	Days	Mass of Plutonium Discharged From One Cycle RB Group	Mass of Plutonium Discharged From Two Cycles RB Group	Mass of Plutonium Discharged From Three Cycles RB Group
1	240	34.8 kg	-	-
2	480	34.8 kg	8.7 kg	-
3	720	34.8 kg	-	99.0 kg
4	960	34.8 kg	8.7 kg	-
5	1200	34.8 kg	-	-
6	1440	34.8 kg	8.7 kg	99.0 kg

IV.B. Development & Selection of Isotopic Ratios

With output data obtained from reactor core physics and fuel burnup simulations, analyses of various isotopes were necessary for determining if isotopic characteristics

exist within the material, which are unique to the PFBR. The ultimate goal being the development of a suite of isotopic ratios capable of attributing the source reactor from which separated weapons-grade plutonium would have been extracted. As mentioned previously, the “three cycles” radial blanket group comes closest to the target fuel burnup of 1 MWd/kg, and is the material of reference for isotopic analysis. The MCNPX output prints mass and radioactivity data for all the isotopes present in the irradiated fuel, which has mass above a threshold limit. The default threshold limit for isotope mass is 1E-10 grams.²³ Table XII, attached in Appendix B, contains mass and activity information for all isotopes existing at the end of the last 60-day decay period, for the three cycles material.

It would be extremely time consuming to measure every isotope present within the material. Therefore, certain characteristics were evaluated in determining which isotopes would be selected for the fission product to plutonium ratio analysis. Selected isotopes are reported in Table VII, as the expected mass and activity of each isotope which would be present in 1 kg of PUREX separated plutonium. One benefit to reporting the selected isotopes as ratios per mass of plutonium is the ability to scale the data to the mass of the interdicted weapons-grade plutonium. Selection was based on (a) the amount of isotope production, at least a few pico-grams per kg of plutonium, (b) the probability of detection (high gamma energy > 100 keV, long half-life > 100 days, high

radioactivity > 1 micro-curie), (c) reactor type dependency in isotope production, and (d) the PUREX plutonium reprocessing decontamination factor (DF) of the isotope.

Specific decontamination factors on an elemental basis could not be found in open literature; hence, a DF of 10^6 was applied universally. Although, the same DF is applied to all isotopes, with the exception of plutonium, it is still an important characteristic.

Given that the material is separated, fission products will be reduced to trace contaminants in the nearly pure plutonium. For fission products with a small amount of production, the inclusion of a DF from separation results in levels which are undetectable. In Table VII, the isotopes are classified into four groups namely, Prompt Gamma, Delayed Alpha, Other Gamma, and Mass Spectrometry based on the type of detection and how fast results can be obtained. Results can be acquired in a few hours with gamma spectroscopy, whereas it can require on the order of a days and weeks for alpha spectroscopy and mass spectrometry processes, respectively.

Table VII

Selected isotopes per kg of PUREX processed plutonium from PFBR radial blanket fuel.

Candidate Isotope	Expected mass (g) per 1 kg Pu with DF of 10^6	Expected activity (Ci) per 1 kg Pu with DF of 10^6
<i>Prompt Gamma</i>		
^{137}Cs	3.29E-06	2.87E-04
^{144}Ce	1.12E-06	3.57E-03
<i>Delayed Alpha</i>		
^{239}Pu	9.81E+02	6.08E+01
^{242}Pu	2.83E-03	1.12E-05
<i>Other Gamma</i>		
^{134}Cs	4.12E-08	5.33E-05
^{125}Sb	3.99E-08	4.19E-05
^{154}Eu	1.41E-08	3.80E-06
<i>Mass Spectrometry</i>		
^{85}Rb	1.97E-07	Stable Isotope
^{90}Sr	1.09E-06	1.54E-04
^{148}Nd	1.05E-06	Stable Isotope
^{147}Pm	9.56E-07	8.87E-04
^{150}Sm	6.17E-08	Stable Isotope

Table VII contains data on the expected mass and expected radioactivity of selected isotopes within a kilogram of separated weapons-grade plutonium produced in the radial blanket of the PFBR. The same data was collected from the fuel of an Indian PHWR discharged at 1 MWD/kg.¹² The ratio of isotope mass per unit plutonium from the PHWR divided by the isotope mass per unit plutonium from the PFBR depicts the

isotope ratio's reactor dependency. These values are given in Table VIII. Several important observations can be drawn from Table VII and Table VIII and are discussed in the following paragraphs.

Table VIII
Reactor dependency of selected isotope ratios

Candidate Isotope	Ratio of expected mass PHWR/PFBR
<i>Prompt Gamma</i>	
^{137}Cs	12.86
^{144}Ce	28.80
<i>Delayed Alpha</i>	
^{239}Pu	0.98
^{242}Pu	19.77
<i>Other Gamma</i>	
^{134}Cs	3.16
^{125}Sb	5.93
^{154}Eu	3.72
<i>Mass Spectrometry</i>	
^{85}Rb	19.00
^{90}Sr	22.29
^{148}Nd	12.51
^{147}Pm	15.48
^{150}Sm	107.99

The radioactivity concentration of ^{137}Cs and ^{144}Ce isotopes are sufficiently high in 1 kg of plutonium and gamma spectroscopy measurements can be made quickly

(prompt measurements) once such material has been interdicted. Both ^{137}Cs and ^{134}Ce undergo beta radiation decay followed by gamma emissions of 662 keV and 134 keV, respectively. The commonly used burnup monitor, ^{137}Cs , is an interesting isotope to note when being used to display a reactor dependency.

We found that selected fission product to plutonium ratios provide more information and result in larger differences between reactors, than just the isotope mass. The radio-isotope ^{137}Cs , for example, is an attractive isotope for selection. The individual fission yield is high at around 6%, it has a long half-life of over 30 years, and the gamma radiation is easily measurable. However, ^{137}Cs is a direct fission product with a fission yield that is constant regardless of fissile isotope or neutron energy. The amount of ^{137}Cs can provide information on the burnup of a material but no information regarding the source reactor. The ratio of ^{137}Cs to plutonium, though, is found in this study to result in a significant difference between the PFBR and PHWR. The ratios of fission products to plutonium have the ability to decipher between fast and thermal reactors. This is due, in large part, to the amount of plutonium the PFBR breeds. The PFBR has a larger percentage of ^{238}U in the depleted uranium fuel in addition to the effect of a fast neutron spectrum. Thus the PFBR radial blanket produces much more plutonium per initial loading of uranium ($\sim 1\%$ of ^{238}U is converted to ^{239}Pu) than the PHWR.

The radioactivities of ^{239}Pu and ^{242}Pu isotopes are sufficiently high in 1 kg of plutonium and alpha spectroscopy measurements can be made. However, sample preparations are needed for performing alpha spectrometry, which makes this method

slower than prompt gamma radiation measurements. Both ^{239}Pu and ^{242}Pu undergo alpha decay with energies of 5156 keV and 4901 keV, respectively. These alpha energies are distinct enough to be separately seen in the alpha spectra. A fast or thermal neutron spectrum irradiation of the fuel can likely be determined from an alpha or mass spectrometry measurement of ^{239}Pu and ^{242}Pu , alone. The PHWR to PFBR ratio of ^{239}Pu concentration is 0.98, while the ratio of ^{242}Pu concentration is 19.15. This indicates that much less ^{242}Pu is present in the plutonium produced in the fast spectrum. This is a result of the relative differences in neutron interaction cross-sections between absorption and fission at varying neutron energies. Lower concentrations of heavier plutonium isotopes, specifically ^{241}Pu and ^{242}Pu , are present in the PFBR blanket fuel due to fission being more likely than radiative capture at fast neutron energies.

The next set of isotopes ^{134}Cs , ^{125}Sb , and ^{154}Eu are again proposed to be measured via gamma spectroscopy. These three isotopes all decay via beta radiation and a subsequent gamma emission, with the gamma energies high enough to be detected easily. The radioactivity concentrations for these gamma emitting isotopes, however, are orders-of-magnitude less than the prompt gamma measurement isotopes, ^{137}Cs and ^{144}Ce . Mass spectrometry is anticipated as the measurement technique for ^{85}Rb , ^{90}Sr , ^{148}Nd , ^{147}Pm , and ^{150}Sm . The isotopes ^{85}Rb , ^{148}Nd , and ^{150}Sm are stable and are thus undetectable using radiation measurements. The isotopes ^{90}Sr and ^{147}Pm are pure beta radiation emitters without any gamma energy emissions. Additionally, isotopes having a range in values for both reactor dependency and isotope mass number were considered desirable during selection. Of the isotopes proposed to be measured using mass

spectroscopy, ^{150}Sm is particularly significant. The PHWR to PFBR ratio of ^{150}Sm concentration is ~ 107 , meaning plutonium produced in a thermal neutron spectrum will have two orders-of-magnitude more ^{150}Sm contamination than plutonium produced in a fast spectrum. The source of this large difference is a result of the radiative capture cross-section of the well-known fission product neutron poison, ^{149}Sm . The second most important neutron absorber in nuclear reactor physics, ^{149}Sm , has a very large cross-section for absorption of thermal neutrons. A plot obtained from reference 32 of the ^{149}Sm radiative capture cross-section per incident neutron energy is shown in Figure 16.³² The dominant neutron energy of the PFBR is 100 keV to 400 keV, whereas the dominant neutron energy of the PHWR is 0.01 eV to 0.1 eV. When applying these dominant neutron energies to the cross-section plot in Figure 16, it can be seen that the ^{149}Sm neutron absorption cross-section in the PFBR is less than 1 barn, while the ^{149}Sm neutron absorption cross-section in the PHWR is around 1E+5 barns (1 barn = 1E-24 cm²).

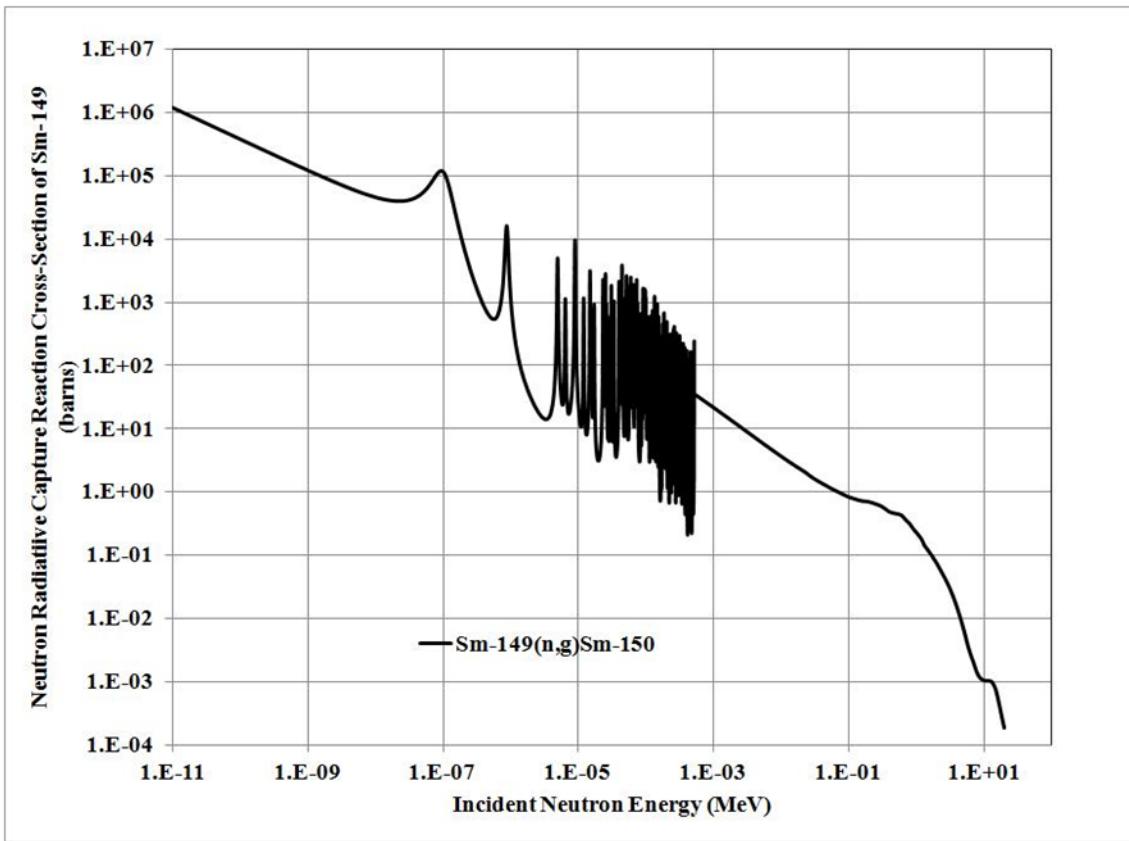


Figure 16. Plot of the neutron radiative capture cross-section for ^{149}Sm .

IV.C. Stochastic Uncertainty

The MCNPX code used to simulate core operation is based on the principles of stochastic methods for solving the Boltzmann radiation transport equation. Because of the stochastic nature of the solution method, the burnup simulations were repeated by altering the stochastic procedures to estimate the stochastic uncertainty associated with the predicted values of fission product and plutonium isotope concentrations. The cycle 1 simulation was repeated by changing the random number seed, changing the sampling procedures, and resulting in nine independent simulations. For every burnup time step in

fuel cycle 1 (180-days), the average isotope concentration in radial blanket fuel (μ) and one sigma standard deviation (σ) values for each isotope were calculated from the results of the nine independent simulations. The issue came in propagating the random error calculated in cycle 1 to the final isotope concentrations at the end of cycle 3. However, it was found, by plotting the data for the selected isotope mass from the nine cycle-1 simulations (8 burnup time steps data points) that the variance (σ^2) increases roughly linearly with isotope mass. An example plot obtained for ^{137}Cs is shown in Figure 17. Similar plots for the other eleven selected isotopes showed the correlation between mass and variance. The regression line equation and fit for the twelve isotope plots are listed in Table IX. It was determined that, using the regression line equation, an estimate of the variance can be extrapolated out to the final isotope mass at the end of cycle 3. The variance was then adjusted for a larger sample size. The nine independent simulations were performed at an earlier time in the research with 9E+5 particle histories, while the final production simulations were run with 2E+7 histories. The larger number of particles increases the precision of the simulations; however, limits on computational time and power made running the nine independent simulations again with 2E+7 histories unviable. The relative random error (σ/μ) was thus obtained for the selected isotopes and is listed in Table X. From Table X it can be observed that the relative random error in predicted isotope concentrations are insignificant and are in the range of 0.019% to 0.054%. This small error indicates that the differences in isotopic compositions seen in the simulations are not due to the Monte Carlo method's random behavior.

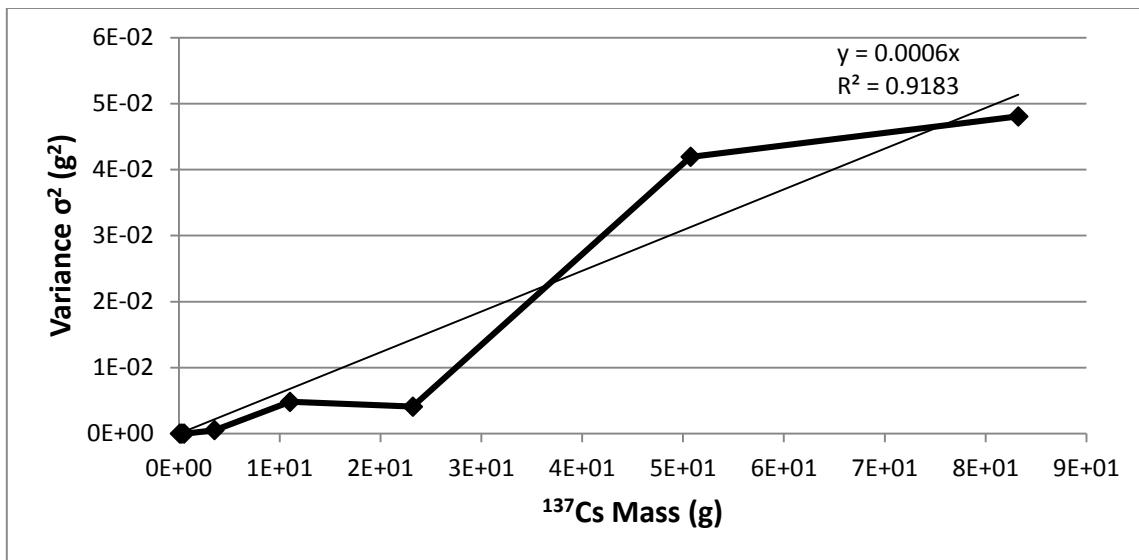


Figure 17. Plot of variance per average mass of ^{137}Cs from nine cycle-1 simulations.

Table IX

Linear regression equation and fit between variance and mass for the selected isotopes.

Isotope	Regression Line Equation*	Regression Line Fit (R^2)
^{137}Cs	$y = 0.0006x$	0.9183
^{144}Ce	$y = 0.0004x$	0.8896
^{239}Pu	$y = 0.0819x$	0.9243
^{242}Pu	$y = (5.0\text{E-}6)x$	1
^{134}Cs	$y = (2.0\text{E-}5)x$	0.9831
^{125}Sb	$y = (5.0\text{E-}6)x$	0.9468
^{154}Eu	$y = (3.0\text{E-}6)x$	0.9585
^{85}Rb	$y = (3.0\text{E-}5)x$	0.9178
^{90}Sr	$y = (2.0\text{E-}4)x$	0.8971
^{148}Nd	$y = (3.0\text{E-}4)x$	0.9450
^{147}Pm	$y = (2.0\text{E-}4)x$	0.9565
^{150}Sm	$y = (4.0\text{E-}5)x$	0.9920

*x being the respective mass of the isotope in grams

Table X
Stochastic uncertainty associated with MCNPX simulations.

Isotope	Relative Random Error (%)
^{137}Cs	0.029
^{144}Ce	0.040
^{239}Pu	0.019
^{242}Pu	0.090
^{134}Cs	0.047
^{125}Sb	0.024
^{154}Eu	0.031
^{85}Rb	0.026
^{90}Sr	0.029
^{148}Nd	0.036
^{147}Pm	0.031
^{150}Sm	0.054

The possibility exists for other sources of error in the model simulations that can affect the isotopic results. Monte Carlo methods have two types of uncertainties; random and systematic. Table X gives the random uncertainty associated with the Monte Carlo method for the twelve selected isotopes. The systematic uncertainty is associated with how close to reality the model is. We assume the systematic uncertainty is small. Uncertainty in the isotopes' neutron interaction cross-sections used in the simulation has the possibility to lead to larger error, however this analysis is beyond the scope of this thesis.

IV.D. Isotopic Ratios of the Same Element

Although the isotopes of an element behave very differently in nuclear reactions, they have very similar chemical properties. The fact that the isotopes of an element have similar chemical properties means they will behave similarly during PUREX chemical separation.¹⁸ Therefore, select isotope ratios of the same element may infer details of the reactor system while being independent from the chemical separation process used.

Table XI gives ratios of the mass of isotopes of the same element present in weapons-grade plutonium produced in a PFBR and PHWR. The values for the isotope ratios found in PHWR plutonium are then divided by the PFBR values to represent the reactor dependency of the isotope of the same element ratios.

Table XI
Mass ratios for isotopes of the same element.

Isotope Ratio	Ratio of Isotope Mass from PFBR	Ratio of Isotope Mass from PHWR	Reactor Comparison PHWR / PFBR
$^{137}\text{Cs} / ^{134}\text{Cs}$	8.00E+01	3.25E+02	4.07
$^{144}\text{Ce} / ^{142}\text{Ce}$	4.18E-01	8.04E-01	1.92
$^{150}\text{Sm} / ^{154}\text{Sm}$	5.04E-01	8.88E+00	17.6
$^{242}\text{Pu} / ^{239}\text{Pu}$	2.89E-06	5.83E-05	20.2

The $^{137}\text{Cs} / ^{134}\text{Cs}$ has an observable difference while the ratios of $^{150}\text{Sm} / ^{154}\text{Sm}$ and $^{242}\text{Pu} / ^{239}\text{Pu}$ both contain more than an order-of-magnitude difference between the

PFBR and PHWR. These three ratios could be measured by gamma spectroscopy, mass spectrometry, and alpha spectroscopy, respectively. It is therefore possible to deduce information of the producing reactor system from isotope ratios of the same element. The benefit of such ratios being, that there are no longer assumptions involved in determining which chemical separation process was used and applying the appropriate decontamination factors.

IV.E. Theoretical Procedures

In the event of an interdicted sample of weapons-grade plutonium multiple measurements of several sub-samples will be performed in order to obtain information on the material and its source as quickly as possible. Immediately gamma spectroscopy measurements will be started and within a few hours results can be drawn on the gamma emitting isotopes present in the material. Due to sample preparation times it is expected that results will be obtained from alpha spectroscopy and mass spectrometry on the order of a few days and a week, respectively. Measured isotopic results will be scaled based on the size of the sub-sample and compared with the expected mass of the selected isotopes in Table VII. As time progresses, results for more of the isotopes contained in Table VII will become available and the certainty in attributing the sample material to a fast or thermal neutron spectrum reactor type will increase.

V. CONCLUSIONS AND FUTURE WORK

V.A. Conclusions

The goal of this thesis was to develop a nuclear forensics capability for targeting plutonium produced in foreign fuel cycles. Monte Carlo computational radiation transport methods coupled with fuel burnup calculations were used to simulate the isotopic composition of plutonium produced in the radial blankets of the 500-MWe Indian Prototype Fast Breeder Reactor. Detailed investigation of fission product contaminants were analyzed to determine the feasibility of predicting intrinsic characteristics in separated weapons-grade plutonium produced by certain reactors, specifically a fast breeder reactor and thermal heavy water reactor. These two reactor types possess the capability to produce weapons-grade plutonium and will likely be operating in a non-safeguarded manner in some countries.

The radiation transport code, MCNPX-2.7, along with the burnup/depletion code, CINDER90, was used to computationally simulate the operation of the PFBR. The PFBR radial blanket fuel will be discharged with a low burnup of around 1 MWD/kg during normal operation; thus, resulting in the average production of around 100 kg of weapons-grade plutonium per year. The data output by MCNPX-2.7 for the fission product and actinide isotopes of the PFBR radial blankets were analyzed to develop isotopic ratios expected to be present in separated weapons-grade plutonium which was produced in an FBR.

Isotope selection was based on four criteria: the amount of isotope production, the probability of isotope detection, the magnitude of reactor type dependency in the isotope's production, and the PUREX decontamination factor. By using this selection criteria, down-selection of the possible isotopes resulted in a list of twelve isotopic ratios. The twelve selected isotopes include: ^{85}Rb , ^{90}Sr , ^{125}Sb , ^{134}Cs , ^{137}Cs , ^{144}Ce , ^{147}Pm , ^{148}Nd , ^{150}Sm , ^{154}Eu , ^{239}Pu , and ^{242}Pu . Table VII reports the expected mass and activity of the selected isotopes present in 1 kg of separated plutonium, which was produced in the PFBR radial blanket. In the event of an interdicted sample Table VII gives a scalable database for the amount of an isotope expected to be present in material produced in a fast reactor system to compare versus measurement data.

Evaluation of isotopic ratios between the PFBR and PHWR helped to select isotope ratios which contain information dependent upon the reactor system. Table VIII shows the magnitude of the reactor dependency for the selected isotopes. It is clearly indicated from Table VIII that the suite of selected isotopes can attribute the source of separated weapons-grade plutonium to a fast or thermal neutron reactor system. The significance of some isotopes stand out from Table VIII.

When normalizing to total plutonium the concentration of ^{137}Cs is an order-of-magnitude larger in the thermal neutron system. This is due to the efficiency of the FBR system at breeding plutonium out of uranium. The rate of ^{137}Cs production remains constant, while they fast breeder system produces ten times more plutonium per initial uranium loading than the thermal neutron system.

A fast or thermal neutron spectrum irradiation of the fuel can likely be determined from an alpha or mass spectrometry measurement of ^{239}Pu and ^{242}Pu , alone. The PHWR to PFBR ratio of ^{239}Pu concentration is 0.98, while the ratio of ^{242}Pu concentration is 19.15. Much less ^{242}Pu is present in the plutonium produced in the fast spectrum as a result of fission being more likely than radiative capture at fast neutron energies.

Of the isotopes proposed to be measured using mass spectroscopy, ^{150}Sm is the most significant. The PHWR to PFBR ratio of ^{150}Sm concentration is ~ 107 , meaning plutonium produced in a thermal neutron spectrum will have two orders-of-magnitude more ^{150}Sm contamination than plutonium produced in a fast spectrum. The source of this drastic ^{150}Sm difference is the very large thermal neutron absorption cross-section of the fission product poison, ^{149}Sm .

Lastly, isotope ratios of the same element were explored. The ratios of $^{137}\text{Cs} / ^{134}\text{Cs}$, $^{150}\text{Sm} / ^{154}\text{Sm}$, and $^{242}\text{Pu} / ^{239}\text{Pu}$ show that such ratios may lead to attribution of a source reactor system, while being independent of the chemical separation process used for plutonium separation.

In conclusion, the computational results indicate a suite of selected isotopic ratios can attribute separated weapons-grade plutonium to a fast or thermal neutron source reactor system.

V.B. Future Work

Experimental data is vital to the verification of the computational results. To accomplish this, samples of depleted UO_2 were irradiated at the Oak Ridge National

Laboratory – High Flux Isotope Reactor (ORNL-HFIR) facility. An image of the HFIR core can be seen in Figure 18.³³ Depleted UO₂ fuel samples were irradiated in a simulated fast neutron spectrum to replicate the PFBR irradiation. The samples were burnt to achieve a burnup between 0.85 – 1.25 MWd/kg. To achieve the fast spectrum, a gadolinium sheath was placed around the fuel samples during irradiation. The high thermal neutron absorption of gadolinium was expected to create a fast neutron spectrum environment with a fast to thermal flux ratio of ≥ 100 .

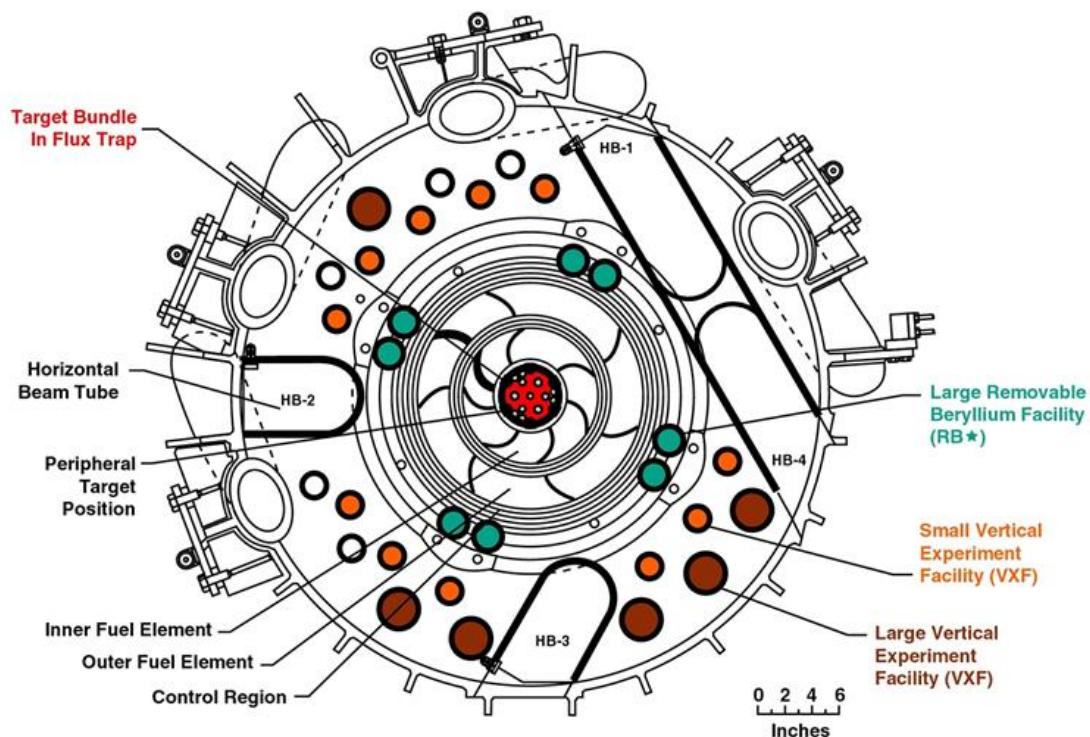


Figure 18. Image of the HFIR core.

The irradiated samples were received by Texas A&M University for chemical reprocessing. A lab scale PUREX process will be performed to separate the plutonium from fission products and uranium. The PUREX process was chosen for this research because it is the most commonly employed reprocessing technique currently in use at reprocessing facilities worldwide, including India. A challenge involved with using the PUREX process is attempting to accurately simulate this process in a laboratory setting. It is believed though, that reprocessing of weapons-grade plutonium will likely occur in a small batch process.

An MCNPX model of the ORNL-HFIR core will be completed with the fuel samples, sample holder, and gadolinium sheath in the irradiation location. Although the PFBR model is simulating the Indian PFBR, the fuel samples are being irradiated in a replicated fast neutron spectrum at HFIR. Thus, accuracy in the HFIR simulation when compared to the experimental samples will provide confidence in the MCNPX models, indirectly verifying the selected isotopes from the PFBR simulation.

Similar to the depleted UO₂ samples, samples of natural UO₂ will be irradiated at the ORNL-HFIR facility in the future. The natural UO₂ fuel samples will be irradiated in a thermal neutron spectrum to replicate the PHWR irradiation. Again, the burnup of the fuel will be to a level near 1 MWd/kg. The lab scale PUREX process will be repeated, as will an MCNPX model of the ORNL-HFIR core with the natural UO₂ fuel samples in place. The measured isotopic data will then be compared with the computational results in order to verify the PHWR model and simulations.

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APPENDIX A

FBR Core FIRST CRITICAL With Depleted UO₂ Axial and Radial Blanket
c Changing to a single burn/decay cycle and removing material concentrations manually *
BOC-3 * **** 2e7 particles ****

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1 0 -1 17 -21 fill=1           imp:n=1 $ Core inner
2 0 -101 102 -103 104 -105 106 lat=2 u=1 imp:n=1
fill=-13:13 -13:13 0:0
17 26R
17 12R 17 17 6 6 6 6 6 6 6 6 17 17 17
17 11R 17 6 6 6 6 6 6 6 6 6 17 17
17 10R 6 6 6 6 900 900 900 900 900 900 900 900 6 6 6 6 17
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17 8R 6 6 900 900 902 901 901 803 802 801 802 901 902 900 900 6 6 17
17 7R 6 6 900 900 901 803 801 802 803 802 801 802 803 901 900 900 6 6 17
17 6R 6 6 900 901 803 801 801 803 804 23 804 803 801 802 901 901 900 6 &
6 17
17 5R 6 6 900 901 801 803 802 703 701 703 702 701 703 803 801 803 901 &
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17 4R 6 6 900 901 802 801 804 703 702 702 701 703 702 701 804 803 801 &
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17 3R 6 6 900 901 801 803 23 702 701 27 703 701 23 703 703 23 801 802 &
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17 17 6 6 6 6 6 6 6 6 6 17 17 11R

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17 17 17 6 6 6 6 6 6 6 6 17 17 17 12R
 17 26R
 C Universe 701 is FUEL SA CORE INNER
 3 0 -401 402 -403 404 -405 406 15 -16 fill=2 u=701 imp:n=1 \$ SA hex can inner
 4 0 -201 202 -203 204 -205 206 lat=2 u=2 imp:n=1
 fill=-9:9 -9:9 0:0
 12 18R
 12 8R 11 11 11 11 11 11 11 11 11 11 11 11 12
 12 7R 11 11 11 11 11 11 11 11 11 11 11 11 12
 12 6R 11 11 11 11 11 11 11 11 11 11 11 11 12
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 12 4R 11 11 11 11 11 11 11 11 11 11 11 11 12
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 12 18R
 5 2 -8.00 (-501 502 -503 504 -505 506) &
 (401:-402:403:-404:405:-406) 15 -16 u=701 imp:n=1 \$ SA hex can
 6 3 -0.90304 (501:-502:503:-504:505:-506) &
 15 -16 u=701 imp:n=1 \$ SA hex can outer
 7 11 0.05969518 -10 -15 u=701 imp:n=1 \$ SA bottom
 8 12 0.03004312 -10 16 -18 u=701 imp:n=1 \$ SA top homo plenum
 9 15 0.05386495 -10 18 -19 u=701 imp:n=1 \$ Core top SS
 10 16 0.05761362 -10 19 u=701 imp:n=1 \$ Core top B4C
 11 0 -4 -2 u=11 imp:n=1 \$ plenum bot
 12 0 -26 2 -8 u=11 imp:n=1 \$ ax blank hole
 13 501 -11.65154927 26 -4 2 -8 u=11 imp:n=1 vol=6.49425 \$ ax blanket bot
 14 0 -26 8 -9 u=11 imp:n=1 \$ fuel hole
 15 101 -10.64258118 26 -4 8 -9 u=11 imp:n=1 vol=21.6475 \$ fuel 10.7737803
 16 0 -26 9 -3 u=11 imp:n=1 \$ ax blank hole
 17 501 -11.65154927 26 -4 9 -3 u=11 imp:n=1 vol=6.49425 \$ ax blanket top
 18 0 -4 3 u=11 imp:n=1 \$ plenum top
 19 0 4 -5 u=11 imp:n=1 \$ fuel clad gap
 20 2 -8.00 5 -6 u=11 imp:n=1 \$ fuel clad
 21 3 -0.90304 6 -7 u=11 imp:n=1 \$ Na out pin
 22 3 -0.90304 -7 u=12 imp:n=1 \$ Na filling tube
 C Universe 702 is FUEL SA CORE INNER
 23 0 -401 402 -403 404 -405 406 15 -16 fill=3 u=702 imp:n=1 \$ SA hex can inner
 24 0 -201 202 -203 204 -205 206 lat=2 u=3 imp:n=1

127 11 0.05969518 -10 -15 u=804 imp:n=1 \$ SA bottom
 128 12 0.03004312 -10 16 -18 u=804 imp:n=1 \$ SA top homo plenum
 129 15 0.05386495 -10 18 -19 u=804 imp:n=1 \$ Core top SS
 130 16 0.05761362 -10 19 u=804 imp:n=1 \$ Core top B4C
 131 0 -4 -2 u=46 imp:n=1 \$ plenum bot
 132 0 -26 2 -8 u=46 imp:n=1 \$ ax blank hole
 133 514 -11.65154927 26 -4 2 -8 u=46 imp:n=1 vol=6.49425 \$ ax blanket bot
 134 0 -26 8 -9 u=46 imp:n=1 \$ fuel hole
 135 404 -10.67464722 26 -4 8 -9 u=46 imp:n=1 vol=21.6475 \$ fuel 10.7737803
 136 0 -26 9 -3 u=46 imp:n=1 \$ ax blank hole
 137 514 -11.65154927 26 -4 9 -3 u=46 imp:n=1 vol=6.49425 \$ ax blanket top
 138 0 -4 3 u=46 imp:n=1 \$ plenum top
 139 0 4 -5 u=46 imp:n=1 \$ fuel clad gap
 140 2 -8.00 5 -6 u=46 imp:n=1 \$ fuel clad
 141 3 -0.90304 6 -7 u=46 imp:n=1 \$ Na out pin
 142 3 -0.90304 -7 u=47 imp:n=1 \$ Na filling tube
 C Universe 4 is Na tube of FA size
 143 3 -0.90304 -10 u=4 imp:n=1 \$ NA filling
 C Universe 900 is Radial Blanket SA
 144 0 -401 402 -403 404 -405 406 22 -20 fill=31 u=900 imp:n=1 \$ SA hex can inner
 145 0 -301 302 -303 304 -305 306 lat=2 u=31 imp:n=1
 fill=-5:5 -5:5 0:0
 16 10R
 16 4R 15 15 15 15 15 16
 16 3R 15 15 15 15 15 16
 16 2R 15 15 15 15 15 15 16
 16 1R 15 15 15 15 15 15 15 16
 16 15 15 15 15 15 15 15 15 16
 16 15 15 15 15 15 15 15 15 16 1R
 16 15 15 15 15 15 15 15 16 2R
 16 15 15 15 15 15 15 16 3R
 16 15 15 15 15 15 16 4R
 16 10R
 146 2 -8.00 (-501 502 -503 504 -505 506) &
 (401:-402:403:-404:405:-406) 22 -20 u=900 imp:n=1 \$ SA hex can
 147 3 -0.90304 (501:-502:503:-504:505:-506) &
 22 -20 u=900 imp:n=1 \$ SA hex can out
 148 11 0.05969518 -10 -15 u=900 imp:n=1 \$ SA bottom
 149 13 0.06846700 -10 15 -22 u=900 imp:n=1 \$ RBPSS
 150 14 0.02912191 -10 20 -18 u=900 imp:n=1 \$ RBPT
 151 15 0.05386495 -10 18 -19 u=900 imp:n=1 \$ RBSS top
 152 16 0.05761362 -10 19 u=900 imp:n=1 \$ RBB4C top
 153 0 -11 -2 u=15 imp:n=1 \$ rad blank ple bot
 154 600 -10.59230329 -11 2 -3 u=15 imp:n=1 vol=204.603 \$ rad blanket
 155 0 -11 3 u=15 imp:n=1 \$ rad blank ple top
 156 0 11 -12 u=15 imp:n=1 \$ blank clad gap
 157 2 -8.00 12 -13 u=15 imp:n=1 \$ blanket clad
 158 3 -0.90304 13 -10 u=15 imp:n=1 \$ NA out blanket

159 3 -0.90304 -10 u=16 imp:n=1 \$ NA filling tube
 C Universe 901 is Radial Blanket SA
 160 0 -401 402 -403 404 -405 406 22 -20 fill=32 u=901 imp:n=1 \$ SA hex can inner
 161 0 -301 302 -303 304 -305 306 lat=2 u=32 imp:n=1
 fill=-5:5 -5:5 0:0
 49 10R
 49 4R 48 48 48 48 48 49
 49 3R 48 48 48 48 48 49
 49 2R 48 48 48 48 48 49
 49 1R 48 48 48 48 48 48 49
 49 48 48 48 48 48 48 49
 49 48 48 48 48 48 48 49 1R
 49 48 48 48 48 48 49 2R
 49 48 48 48 48 48 49 3R
 49 48 48 48 48 49 4R
 49 10R
 162 2 -8.00 (-501 502 -503 504 -505 506) &
 (401:-402:403:-404:405:-406) 22 -20 u=901 imp:n=1 \$ SA hex can
 163 3 -0.90304 (501:-502:503:-504:505:-506) &
 22 -20 u=901 imp:n=1 \$ SA hex can out
 164 11 0.05969518 -10 -15 u=901 imp:n=1 \$ SA bottom
 165 13 0.06846700 -10 15 -22 u=901 imp:n=1 \$ RBPSS
 166 14 0.02912191 -10 20 -18 u=901 imp:n=1 \$ RBPT
 167 15 0.05386495 -10 18 -19 u=901 imp:n=1 \$ RBSS top
 168 16 0.05761362 -10 19 u=901 imp:n=1 \$ RBB4C top
 169 0 -11 -2 u=48 imp:n=1 \$ rad blank ple bot
 170 601 -10.59230329 -11 2 -3 u=48 imp:n=1 vol=204.603 \$ rad blanket
 171 0 -11 3 u=48 imp:n=1 \$ rad blank ple top
 172 0 11 -12 u=48 imp:n=1 \$ blank clad gap
 173 2 -8.00 12 -13 u=48 imp:n=1 \$ blanket clad
 174 3 -0.90304 13 -10 u=48 imp:n=1 \$ NA out blanket
 175 3 -0.90304 -10 u=49 imp:n=1 \$ NA filling tube
 C Universe 902 is Radial Blanket SA
 176 0 -401 402 -403 404 -405 406 22 -20 fill=33 u=902 imp:n=1 \$ SA hex can inner
 177 0 -301 302 -303 304 -305 306 lat=2 u=33 imp:n=1
 fill=-5:5 -5:5 0:0
 51 10R
 51 4R 50 50 50 50 50 51
 51 3R 50 50 50 50 50 51
 51 2R 50 50 50 50 50 50 51
 51 1R 50 50 50 50 50 50 50 51
 51 50 50 50 50 50 50 50 51
 51 50 50 50 50 50 50 50 51 1R
 51 50 50 50 50 50 50 51 2R
 51 50 50 50 50 50 51 3R
 51 50 50 50 50 51 4R
 51 10R
 178 2 -8.00 (-501 502 -503 504 -505 506) &

(401:-402:403:-404:405:-406) 22 -20 u=902 imp:n=1 \$ SA hex can
 179 3 -0.90304 (501:-502:503:-504:505:-506) &
 22 -20 u=902 imp:n=1 \$ SA hex can out
 180 11 0.05969518 -10 -15 u=902 imp:n=1 \$ SA bottom
 181 13 0.06846700 -10 15 -22 u=902 imp:n=1 \$ RBPSS
 182 14 0.02912191 -10 20 -18 u=902 imp:n=1 \$ RBPT
 183 15 0.05386495 -10 18 -19 u=902 imp:n=1 \$ RBSS top
 184 16 0.05761362 -10 19 u=902 imp:n=1 \$ RBB4C top
 185 0 -11 -2 u=50 imp:n=1 \$ rad blank ple bot
 186 602 -10.59230329 -11 2 -3 u=50 imp:n=1 vol=204.603 \$ rad blanket
 187 0 -11 3 u=50 imp:n=1 \$ rad blank ple top
 188 0 11 -12 u=50 imp:n=1 \$ blank clad gap
 189 2 -8.00 12 -13 u=50 imp:n=1 \$ blanket clad
 190 3 -0.90304 13 -10 u=50 imp:n=1 \$ NA out blanket
 191 3 -0.90304 -10 u=51 imp:n=1 \$ NA filling tube
 C Universe 904 is Fuel SA CORE OUTER *All universe 904 will be replaced with 901 at second refueling
 192 0 -401 402 -403 404 -405 406 15 -16 fill=34 u=904 imp:n=1 \$ SA hex can inner
 193 0 -201 202 -203 204 -205 206 lat=2 u=34 imp:n=1
 fill=-9.9 -9.9 0:0
 53 18R
 53 8R 52 52 52 52 52 52 52 52 52 52 53
 53 7R 52 52 52 52 52 52 52 52 52 52 53
 53 6R 52 52 52 52 52 52 52 52 52 52 53
 53 5R 52 52 52 52 52 52 52 52 52 52 53
 53 4R 52 52 52 52 52 52 52 52 52 52 53
 53 3R 52 52 52 52 52 52 52 52 52 52 53
 53 2R 52 52 52 52 52 52 52 52 52 52 53
 53 1R 52 52 52 52 52 52 52 52 52 52 53
 53 52 52 52 52 52 52 52 52 52 52 53
 53 52 52 52 52 52 52 52 52 52 52 53 1R
 53 52 52 52 52 52 52 52 52 52 52 53 2R
 53 52 52 52 52 52 52 52 52 52 52 53 3R
 53 52 52 52 52 52 52 52 52 52 52 53 4R
 53 52 52 52 52 52 52 52 52 52 52 53 5R
 53 52 52 52 52 52 52 52 52 52 52 53 6R
 53 52 52 52 52 52 52 52 52 52 52 53 7R
 53 52 52 52 52 52 52 52 52 52 52 53 8R
 53 18R
 194 2 -8.00 (-501 502 -503 504 -505 506) &
 (401:-402:403:-404:405:-406) 15 -16 u=904 imp:n=1 \$ SA hex can
 195 3 -0.90304 (501:-502:503:-504:505:-506) &
 15 -16 u=904 imp:n=1 \$ SA hex can outer
 196 11 0.05969518 -10 -15 u=904 imp:n=1 \$ SA bottom
 197 12 0.03004312 -10 16 -18 u=904 imp:n=1 \$ SA top homo plenum
 198 15 0.05386495 -10 18 -19 u=904 imp:n=1 \$ Core top SS
 199 16 0.05761362 -10 19 u=904 imp:n=1 \$ Core top B4C
 200 0 -4 -2 u=52 imp:n=1 \$ plenum bot

201 0 -26 2 -8 u=52 imp:n=1 \$ ax blank hole
 202 524 -11.65154927 26 -4 2 -8 u=52 imp:n=1 vol=6.49425 \$ ax blanket bot
 203 0 -26 8 -9 u=52 imp:n=1 \$ fuel hole
 204 604 -10.67464722 26 -4 8 -9 u=52 imp:n=1 vol=21.6475 \$ fuel 10.80996965
 205 0 -26 9 -3 u=52 imp:n=1 \$ ax blank hole
 206 524 -11.65154927 26 -4 9 -3 u=52 imp:n=1 vol=6.49425 \$ ax blanket top
 207 0 -4 3 u=52 imp:n=1 \$ plenum top
 208 0 4 -5 u=52 imp:n=1 \$ fuel clad gap
 209 2 -8.00 5 -6 u=52 imp:n=1 \$ fuel clad
 210 3 -0.90304 6 -7 u=52 imp:n=1 \$ Na out pin
 211 3 -0.90304 -7 u=53 imp:n=1 \$ Na filling tube
 C Universe 6 is SS reflector SA
 212 11 0.05969518 -10 -15 u=6 imp:n=1 \$ SS Refl Ass bot
 213 7 0.09365394 -10 15 -8 u=6 imp:n=1 \$ SS Reflector B4C
 214 8 0.06154800 -10 8 -20 u=6 imp:n=1 \$ SS Reflector
 215 14 0.02912191 -10 20 -18 u=6 imp:n=1 \$ SS reflector top
 216 7 0.09365394 -10 18 u=6 imp:n=1 \$ SS refle B4C top
 C Universe 17 is B4C Shield SA
 217 11 0.05969518 -10 -15 u=17 imp:n=1 \$ B4C SHLD bottom
 218 17 0.01835245 -10 15 -23 u=17 imp:n=1 \$ SHLD Plenum bot
 219 7 0.09365394 -10 23 -24 u=17 imp:n=1 \$ B4C Shld I layer
 220 17 0.01835245 -10 24 -25 u=17 imp:n=1 \$ SHLD Plenum top
 221 18 0.06221962 -10 25 u=17 imp:n=1 \$ SHLD SS top
 C Universe 18 is CSR/DSR
 222 9 0.03393119 -10 -8 u=18 imp:n=1 \$ CSR/DSR bottom
 223 10 0.06340921 -10 8 -14 u=18 imp:n=1 \$ CSR/DSR
 224 9 0.03393119 -10 14 u=18 imp:n=1 \$ CSR/DSR top
 C Universe 19 is Diluent SA *All universe 19 will be replaced with 703 at first refueling
 225 11 0.05969518 -10 -15 u=19 imp:n=1 \$ SA bottom
 226 13 0.06846700 -10 15 -22 u=19 imp:n=1 \$ RBPSS
 227 20 0.02324489 -10 22 -2 u=19 imp:n=1 \$ RB Plenum bot
 228 19 0.05133189 -10 2 -3 u=19 imp:n=1 vol= 43740 \$ Diluent with RBP
 229 20 0.02324489 -10 3 -20 u=19 imp:n=1 \$ RB Plenum top
 230 14 0.02912191 -10 20 -18 u=19 imp:n=1 \$ RBPT
 231 15 0.05386495 -10 18 -19 u=19 imp:n=1 \$ RBSS top
 232 16 0.05761362 -10 19 u=19 imp:n=1 \$ RBB4C top
 C Universe 23 is pinwise CSR
 233 0 -401 402 -403 404 -405 406 36 -37 fill=36 u=23 imp:n=1 \$ SA hex can inner
 234 0 -601 602 -603 604 -605 606 lat=2 u=36 imp:n=1
 fill=-3:3 -3:3 0:0
 21 6R
 21 2R 20 20 20 21
 21 1R 20 20 20 20 21
 21 20 20 20 20 20 21
 21 20 20 20 20 21 1R
 21 20 20 20 21 2R
 21 6R
 235 2 -8.00 (-501 502 -503 504 -505 506) &

(401:-402:403:-404:405:-406) 36 -37 u=23 imp:n=1 \$ SA hex can
 236 3 -0.90304 (501:-502:503:-504:505:-506) &
 36 -37 u=23 imp:n=1 \$ SA hex can out
 237 9 0.03393119 -10 -36 u=23 imp:n=1 \$ CSR Follower bot
 238 9 0.03393119 -10 37 u=23 imp:n=1 \$ CSR Follower top
 239 22 -2.4 -27 -40 u=20 imp:n=1 \$ CSR pin bot
 240 21 -2.4 -27 40 -41 u=20 imp:n=1 \$ CSR pin mid
 241 22 -2.4 -27 41 u=20 imp:n=1 \$ CSR pin top
 242 0 27 -28 u=20 imp:n=1 \$ CSR clad gap
 243 2 -8.00 28 -29 u=20 imp:n=1 \$ CSR clad
 244 3 -0.90304 29 -10 u=20 imp:n=1 \$ NA out CSR
 245 3 -0.90304 -10 u=21 imp:n=1 \$ NA filling tube
 C Universe 27 is pinwise DSR
 246 0 -401 402 -403 404 -405 406 38 -39 fill=37 u=27 imp:n=1 \$ SA hex can inner
 247 0 -601 602 -603 604 -605 606 lat=2 u=37 imp:n=1
 fill=-3:3 -3:3 0:0
 25 6R
 25 2R 24 24 24 25
 25 1R 24 24 24 24 25
 25 24 24 24 24 25
 25 24 24 24 25 1R
 25 24 24 25 2R
 25 6R
 248 2 -8.00 (-501 502 -503 504 -505 506) &
 (401:-402:403:-404:405:-406) 38 -39 u=27 imp:n=1 \$ SA hex can
 249 3 -0.90304 (501:-502:503:-504:505:-506) &
 38 -39 u=27 imp:n=1 \$ SA hex can out
 250 9 0.03393119 -10 -38 u=27 imp:n=1 \$ DSR Follower bot
 251 9 0.03393119 -10 39 u=27 imp:n=1 \$ DSR Follower top
 252 21 -2.4 -33 u=24 imp:n=1 \$ DSR pin mid
 253 0 33 -34 u=24 imp:n=1 \$ DSR clad gap
 254 2 -8.00 34 -35 u=24 imp:n=1 \$ DSR clad
 255 3 -0.90304 35 -10 u=24 imp:n=1 \$ NA out DSR
 256 3 -0.90304 -10 u=25 imp:n=1 \$ NA filling tube
 C Universe 28 is ALSO Diluent SA *Both universe 28 will be replaced with 702 at second refueling
 257 11 0.05969518 -10 -15 u=28 imp:n=1 \$ SA bottom
 258 13 0.06846700 -10 15 -22 u=28 imp:n=1 \$ RBPSS
 259 20 0.02324489 -10 22 -2 u=28 imp:n=1 \$ RB Plenum bot
 260 23 0.05133189 -10 2 -3 u=28 imp:n=1 vol= 43740 \$ Diluent with RBP
 261 20 0.02324489 -10 3 -20 u=28 imp:n=1 \$ RB Plenum top
 262 14 0.02912191 -10 20 -18 u=28 imp:n=1 \$ RBPT
 263 15 0.05386495 -10 18 -19 u=28 imp:n=1 \$ RBSS top
 264 16 0.05761362 -10 19 u=28 imp:n=1 \$ RBB4C top
 265 0 1:-17:21 imp:n=0
 1 cz 155 \$ core vessel rad
 2 pz 0 \$ blanket bottom

3	pz 160	\$ blanket top
4	cz 0.2775	\$ fuel pellet rad
5	cz 0.285	\$ fuel clad ID
6	cz 0.33	\$ fuel clad OD
7	cz 5.0	\$ outer pin Na
8	pz 30	\$ bot blank end
9	pz 130	\$ top blank start
10	cz 20	\$ dummy NA
11	cz 0.638	\$ radi blank rad
12	cz 0.6565	\$ blank clad ID
13	cz 0.7165	\$ blank clad OD
14	pz 141	\$ CRF top Start
15	pz -75	\$ plenum bottom
16	pz 183	\$ plenum top
17	pz -101	\$ SA bottom
18	pz 191.5	\$ remaining plenum
19	pz 257.0	\$ SA SS top
20	pz 170	\$ Rad blnk ple top
21	pz 267	\$ Core B4C top
22	pz -60	\$ RB ple bot SS
23	pz 3.9	\$ SHLD plenum bot
24	pz 238.2	\$ B4C shld top
25	pz 248.4	\$ SHPL top
26	cz 0.09	\$ fuel annular rad
27	cz 0.87	\$ CSRB4C pellet OR
28	cz 1.02	\$ CSRB4C clad IR
29	cz 1.12	\$ CSRB4C clad OR

C The following three cards are not required any more

C 30	pz 50	\$ CSRnatB4C bot
C 31	pz 121	\$ CSRnatB4C top
C 32	pz 131	\$ DSRB4C top
33	cz 0.89	\$ DSRB4C pellet OR
34	cz 1.00	\$ DSRB4C clad IR
35	cz 1.07	\$ DSRB4C clad OR

C ****

C The following PZ's are for pin wise CSR and DSR inserion and withdrawal

C Change the comment card accordingly

C ****

C 36	pz 29.0	\$ CSR DOWN (bottom edge)
C 37	pz 140.0	\$ CSR DOWN (top edge)
36	pz 137.5	\$ CSR UP (bottom edge)
37	pz 248.5	\$ CSR UP (top edge)
C 38	pz 29.5	\$ DSR DOWN (bottom edge)
C 39	pz 130.5	\$ DSR DOWN (top edge)
38	pz 131.5	\$ DSR UP (bottom edge)
39	pz 232.5	\$ DSR UP (top edge)

```

C ****
*
C following pairs are CSR down & up >>> how the pin axial profile change
C ****
*
C 40 pz 49.0 $ CSR DOWN (bottom nat B4C pin top)
C 41 pz 120.0 $ CSR DOWN (mid enrich B4C pin top)
40 pz 157.5 $ CSR UP (bottom nat B4C pin top)
41 pz 228.5 $ CSR UP (mid enrich B4C pin top)
C ****
*
101 px 6.75 $ hexside FA
102 px -6.75 $ hexside FA
103 p 1 1.7320508076 0 13.5 $ hexside FA
104 p 1 1.7320508076 0 -13.5 $ hexside FA
105 p -1 1.7320508076 0 13.5 $ hexside FA
106 p -1 1.7320508076 0 -13.5 $ hexside FA
201 py 0.4125 $ hexside pin
202 py -0.4125 $ hexside pin
203 p 1.7320508076 1 0 0.825 $ hexside pin
204 p 1.7320508076 1 0 -0.825 $ hexside pin
205 p 1.7320508076 -1 0 0.825 $ hexside pin
206 p 1.7320508076 -1 0 -0.825 $ hexside pin
301 py 0.8 $ hexside pin
302 py -0.8 $ hexside pin
303 p 1.7320508076 1 0 1.6 $ hexside pin
304 p 1.7320508076 1 0 -1.6 $ hexside pin
305 p 1.7320508076 -1 0 1.6 $ hexside pin
306 p 1.7320508076 -1 0 -1.6 $ hexside pin
401 px 6.26 $ hexside FA
402 px -6.26 $ hexside FA
403 p 1 1.7320508076 0 12.52 $ hexside FA
404 p 1 1.7320508076 0 -12.52 $ hexside FA
405 p -1 1.7320508076 0 12.52 $ hexside FA
406 p -1 1.7320508076 0 -12.52 $ hexside FA
501 px 6.58 $ hexside FA
502 px -6.58 $ hexside FA
503 p 1 1.7320508076 0 13.16 $ hexside FA
504 p 1 1.7320508076 0 -13.16 $ hexside FA
505 p -1 1.7320508076 0 13.16 $ hexside FA
506 p -1 1.7320508076 0 -13.16 $ hexside FA
601 py 1.2 $ hexside pin
602 py -1.2 $ hexside pin
603 p 1.7320508076 1 0 2.4 $ hexside pin
604 p 1.7320508076 1 0 -2.4 $ hexside pin

```

```

605 p 1.7320508076 -1 0 2.4          $ hexside pin
606 p 1.7320508076 -1 0 -2.4

kcode 4000 1 100 5100
ksrc 0.15 0 80
burn time= 0.3, 0.6, 0.6, 8.5, 20, 30, 60, 60, 60
pfrac= 1, 1, 1, 1, 1, 1, 1, 1, 0
power= 1250
mat= 101, 102, 103, 401, &
      402, 403, 404, 501, 502, &
      503, 511, 512, 513, 514, &
      600, 601, 602
matvol= 126832.919, 126832.919, 145622.981, 150320.496,
       103345.341, 140925.465, 56370.186, 76099.751, 76099.751,
       87373.788, 90192.298, 62007.205, 84555.279, 33822.112,
       898614.753, 524191.939, 74884.563
bopt= 1.0 -24 1.0
c Material 101 is from the output of BOC-2
m101 4009 -0.00003159
      6012 -0.8175
      6013 -24.96
      6014 -0.0006309
      7015 -0.04534
      8016 -159900
      8017 -0.004517
      30066 -0.0001049
      30067 -0.0004304
      30068 -0.0009347
      30070 -0.007126
      31069 -0.005577
      31071 -0.03342
      32072 -0.08508
      32073 -0.1241
      32074 -0.3187
      32076 -1.16
      33075 -0.482
      34076 -0.02092
      34077 -2.485
      34078 -4.935
      34079 -12.21
      34080 -19.9
      34082 -50.01
      35079 -0.0001455
      35081 -29.36
      36080 -0.0008921
      36082 -1.096
      36083 -73.37
      36084 -133.5

```

36085	-30.18
36086	-207.8
37085	-111.3
37086	-0.03891
37087	-275
38086	-2.642
38087	-0.01824
38088	-353
38089	-38.52
38090	-554.9
39088	-0.0001797
39089	-432.8
39090	-0.1442
39091	-72.38
40090	-10.56
40091	-613.5
40092	-833.4
40093	-1034
40094	-1168
40095	-153.9
40096	-1382
41094	-0.009213
41095	-122.6
42094	-0.01517
42095	-985.3
42096	-25.9
42097	-1430
42098	-1610
42100	-1888
43098	-0.0122
43099	-1623
44099	-0.0576
44100	-86.91
44101	-1822
44102	-2080
44103	-103.6
44104	-1945
44106	-820.4
45103	-1816
45106	-0.0007615
46102	-0.0004506
46104	-79.29
46105	-1477
46106	-626.1
46107	-922.9
46108	-709.2
46110	-212.7
47109	-334.7

47111	-0.01269
48108	-0.002961
48110	-17.26
48111	-111.2
48112	-64.66
48113	-39.43
48114	-30.94
48116	-20.59
49113	-0.02731
49115	-22.53
50114	-0.001446
50115	-1.208
50116	-1.534
50117	-22.41
50118	-20.71
50119	-19.02
50120	-20.15
50122	-24.46
50123	-1.835
50124	-40.21
50125	-0.01417
50126	-87.59
51121	-19.44
51123	-24.71
51124	-0.1165
51125	-44.91
51126	-0.005904
52122	-0.7428
52123	-0.009007
52124	-0.5654
52125	-8.103
52126	-3.692
52127	-0.02
52128	-292.6
52129	-0.003141
52130	-885
53127	-142
53129	-478.2
53131	-0.2622
54126	-0.0004783
54128	-7.092
54129	-0.09537
54130	-15.9
54131	-1400
54132	-2053
54133	-0.02409
54134	-2921
54136	-2752

55133 -2547
55134 -88.85
55135 -2841
55136 -0.3224
55137 -2569
56132 -0.0003224
56134 -17.02
56135 -0.1233
56136 -94.86
56137 -46.79
56138 -2482
56140 -4.156
57138 -0.04268
57139 -2310
57140 -0.6298
58139 -0.01831
58140 -2220
58141 -72.8
58142 -1979
58144 -868.2
59141 -2002
59143 -4.936
60142 -21.63
60143 -1750
60144 -788.2
60145 -1262
60146 -1160
60147 -0.8588
60148 -748.8
60150 -452.3
61146 -0.01936
61147 -645.4
61148 -0.02346
62146 -0.003767
62147 -126.1
62148 -72.93
62149 -483.4
62150 -98
62151 -271.4
62152 -347
62154 -120.3
63151 -1.3
63152 -0.1168
63153 -165.5
63154 -32.03
63155 -85.22
63156 -0.3436
64152 -0.07784

64153 -0.002549
64154 -1.436
64155 -8.758
64156 -76.57
64157 -45.86
64158 -35.09
64160 -7.71
65159 -15.49
65160 -0.5375
66160 -1.832
66161 -3.999
66162 -3.368
66163 -1.4
66164 -0.9046
67165 -0.3717
68166 -0.2748
68167 -0.1152
68168 -0.0503
68170 -0.004627
69169 -0.008066
69171 -0.000838
70172 -0.0005614
92234 -0.6786
92235.17c -1481
92236 -190.1
92237 -0.01447
92238.17c -886100
93236 -0.005099
93237 -210.7
94237 -0.002623
94238 -81.26
94239.17c -149800
94240 -65880
94241 -11640
94242 -3882
94244 -0.02081
95241 -644.8
95242 -6.17
95243 -250.3
96242 -37.27
96243 -1.001
96244 -33.72
96245 -1.376
96246 -0.03487
96247 -0.0006421
m102 92235.17c -0.0017256 \$ -0.00173
92238.17c -0.6973080 \$ -0.69759
94239.17c -0.1254001 \$ -0.12521

94240.60c -0.0450321 \$ -0.04496

94241.60c -0.0096691 \$ -0.00965

94242.60c -0.0024919 \$ -0.00249

8016.60c -0.1183732 \$ -0.11837 changed wt as per Pu buildup table core I

c Material 103 is from the output of BOC-2

m103 6012 -0.4705

6013 -17.15

6014 -0.0001622

7015 -0.0261

8016 -183400

8017 -0.002514

31069 -0.003248

31071 -0.01946

32072 -0.04913

32073 -0.0726

32074 -0.1835

32076 -0.6714

33075 -0.2841

34076 -0.006115

34077 -1.462

34078 -2.843

34079 -7.172

34080 -11.42

34082 -28.93

35081 -17.42

36080 -0.0004195

36082 -0.3703

36083 -42.96

36084 -76.95

36085 -17.91

36086 -120.2

37085 -64.74

37086 -0.02051

37087 -159.3

38086 -0.7509

38087 -0.003562

38088 -204

38089 -44.12

38090 -323.9

39089 -228.6

39090 -0.08415

39091 -81.19

40090 -3.302

40091 -317

40092 -481.1

40093 -599.8

40094 -672

40095 -169.5

40096	-799.3
41094	-0.003949
41095	-130.4
42094	-0.001929
42095	-438.7
42096	-5.093
42097	-840.1
42098	-922
42100	-1094
43099	-961.4
44099	-0.03272
44100	-20.97
44101	-1084
44102	-1172
44103	-119.3
44104	-1129
44106	-582.3
45103	-1015
46104	-15.35
46105	-880.9
46106	-220.1
46107	-551.8
46108	-386.4
46110	-123.1
47109	-194.8
47111	-0.01463
48108	-0.0008603
48110	-4.526
48111	-65.1
48112	-35.79
48113	-23.09
48114	-17.5
48116	-11.93
49113	-0.01128
49115	-13.48
50114	-0.00025
50115	-0.7017
50116	-0.3819
50117	-13
50118	-11.88
50119	-11.06
50120	-11.54
50122	-14.12
50123	-1.7
50124	-23.21
50125	-0.01682
50126	-50.49
51121	-11.41

51123	-13.8
51124	-0.05256
51125	-28.95
51126	-0.003924
52122	-0.2144
52123	-0.001395
52124	-0.1516
52125	-2.703
52126	-1.318
52128	-168.9
52130	-511
53127	-84.31
53129	-283.7
53131	-0.3105
54128	-1.867
54129	-0.01193
54130	-4.599
54131	-818.4
54132	-1177
54133	-0.02844
54134	-1687
54136	-1589
55133	-1505
55134	-26.23
55135	-1653
55136	-0.2242
55137	-1495
56134	-2.984
56135	-0.008789
56136	-37.38
56137	-15.33
56138	-1432
56140	-4.927
57138	-0.01712
57139	-1335
57140	-0.7465
58139	-0.005921
58140	-1277
58141	-85.85
58142	-1143
58144	-649.5
59141	-1123
59143	-5.844
60142	-4.642
60143	-1033
60144	-294.1
60145	-743.7
60146	-655.9

60147 -1.007
60148 -435
60150 -262.9
61147 -436.5
61148 -0.01241
62147 -44.06
62148 -19.29
62149 -302.7
62150 -28.84
62151 -176.8
62152 -181.4
62154 -70.16
63151 -0.5214
63152 -0.01735
63153 -104.1
63154 -10.3
63155 -50.88
63156 -0.3583
64152 -0.01049
64154 -0.2757
64155 -3.074
64156 -41.68
64157 -26.9
64158 -18.71
64160 -4.47
65159 -9.396
65160 -0.2386
66160 -0.4596
66161 -2.377
66162 -1.788
66163 -0.8033
66164 -0.4914
67165 -0.2218
68166 -0.1533
68167 -0.06641
68168 -0.02498
68170 -0.002684
92234 -0.311
92235.17c -2146
92236 -120.5
92237 -0.01673
92238.17c -1049000
93236 -0.001618
93237 -131.8
94237 -0.003003
94238 -25.98
94239.17c -182500
94240 -72850

94241	-14000
94242	-4138
94244	-0.01015
95241	-422.8
95242	-1.977
95243	-146.1
96242	-12.44
96243	-0.1717
96244	-9.73
96245	-0.2011
96246	-0.002377
m2	26000.55c -0.66598
	6000.66c -0.00052
	24000.50c -0.13800
	28000.50c -0.15200
	42000.66c -0.01460
	14000.60c -0.00920
	25055.60c -0.01740
	22000.62c -0.00230
	\$ SS
m3	11023.62c 1.0
	\$ Na
c	Material 401 is from the output of BOC-2
m401	4009 -0.00002569
	6012 -0.7768
	6013 -23.77
	6014 -0.000397
	7015 -0.04309
	8016 -190000
	8017 -0.003203
	30067 -0.0004042
	30068 -0.0008995
	30070 -0.007165
	31069 -0.005611
	31071 -0.03433
	32072 -0.08688
	32073 -0.1284
	32074 -0.3244
	32076 -1.182
	33075 -0.499
	34076 -0.01275
	34077 -2.564
	34078 -4.98
	34079 -12.55
	34080 -20.11
	34082 -50.36
	35079 -0.0001509
	35081 -30.44
	36080 -0.0008473
	36082 -0.7548

36083	-73.97
36084	-131.2
36085	-30.3
36086	-206.6
37085	-112.4
37086	-0.02327
37087	-274
38086	-1.605
38087	-0.008727
38088	-350
38089	-37.83
38090	-551.7
39089	-429.1
39090	-0.1434
39091	-71.7
40090	-10.41
40091	-613.7
40092	-831.1
40093	-1039
40094	-1162
40095	-153.8
40096	-1390
41094	-0.007922
41095	-122.7
42094	-0.01226
42095	-1004
42096	-16.09
42097	-1459
42098	-1616
42100	-1914
43098	-0.01003
43099	-1670
44099	-0.05891
44100	-53.45
44101	-1888
44102	-2067
44103	-103.7
44104	-1980
44106	-838.3
45103	-1876
45106	-0.0007781
46102	-0.0003736
46104	-49.16
46105	-1548
46106	-596.4
46107	-971.7
46108	-698.9
46110	-219.8

47109	-348.6
47111	-0.01282
48108	-0.002469
48110	-10.44
48111	-116
48112	-64.54
48113	-40.9
48114	-31.2
48116	-20.99
49113	-0.02333
49115	-23.63
50114	-0.000796
50115	-1.239
50116	-0.9763
50117	-22.9
50118	-20.96
50119	-19.34
50120	-20.41
50122	-24.93
50123	-1.882
50124	-40.97
50125	-0.01426
50126	-89.47
51121	-20.07
51123	-25.4
51124	-0.07513
51125	-46.66
51126	-0.004034
52122	-0.4713
52123	-0.003715
52124	-0.3797
52125	-8.476
52126	-2.705
52127	-0.02036
52128	-298.1
52129	-0.003121
52130	-890.3
53127	-147.9
53129	-490.9
53131	-0.2607
54126	-0.0003969
54128	-4.334
54129	-0.03612
54130	-10.14
54131	-1424
54132	-2050
54133	-0.02385
54134	-2933

54136 -2766
55133 -2606
55134 -53.8
55135 -2880
55136 -0.2228
55137 -2588
56132 -0.0002632
56134 -10.27
56135 -0.05015
56136 -72.16
56137 -47.22
56138 -2495
56140 -4.111
57138 -0.03814
57139 -2330
57140 -0.623
58139 -0.01458
58140 -2221
58141 -72.82
58142 -1989
58144 -872
59141 -2033
59143 -4.943
60142 -12.29
60143 -1795
60144 -774.8
60145 -1289
60146 -1141
60147 -0.8553
60148 -753.5
60150 -457.8
61146 -0.01658
61147 -689.2
61148 -0.01525
62146 -0.003195
62147 -136.8
62148 -45.3
62149 -516.7
62150 -59.6
62151 -301.8
62152 -323
62154 -122.7
63151 -1.565
63152 -0.08651
63153 -178.2
63154 -21.03
63155 -85.8
63156 -0.3152

64152 -0.05399
64153 -0.001033
64154 -0.9366
64155 -9.341
64156 -74.86
64157 -47.32
64158 -33.65
64160 -7.955
65159 -16.38
65160 -0.3432
66160 -1.177
66161 -4.191
66162 -3.239
66163 -1.42
66164 -0.8815
67165 -0.3921
68166 -0.274
68167 -0.1177
68168 -0.04584
68170 -0.004722
69169 -0.007995
69171 -0.0008108
70172 -0.0005267
92234 -0.6528
92235.17c -1913
92236 -144
92237 -0.01247
92238.17c -981600
93236 -0.003954
93237 -194
94237 -0.003088
94238 -74.82
94239.17c -234600
94240 -100300
94241 -18450
94242 -5821
94244 -0.01142
95241 -1100
95242 -7.212
95243 -245.9
96242 -37.83
96243 -0.6003
96244 -19.93
96245 -0.506
96246 -0.007651
m402 92235.17c -0.0015732 \$ -0.00157
92238.17c -0.6357256 \$ -0.63568
94239.17c -0.1678273 \$ -0.16786

94240.60c -0.0602681 \$ -0.06028

94241.60c -0.0129404 \$ -0.01294

94242.60c -0.0033350 \$ -0.00334

8016.60c -0.1183303 \$ -0.11833 changed wt as per Pu buildup table core II

c Material 403 is from the output of BOC-2

m403 6012 -0.3893

6013 -14.14

6014 -0.00009539

7015 -0.0216

8016 -178000

8017 -0.001603

31069 -0.002885

31071 -0.01769

32072 -0.04461

32073 -0.06625

32074 -0.1662

32076 -0.6066

33075 -0.2586

34076 -0.003682

34077 -1.326

34078 -2.547

34079 -6.484

34080 -10.26

34082 -25.79

35081 -15.81

36080 -0.0003684

36082 -0.2562

36083 -38.11

36084 -66.97

36085 -15.91

36086 -105.6

37085 -57.56

37086 -0.01209

37087 -140.2

38086 -0.449

38087 -0.002313

38088 -179

38089 -38.65

38090 -284.6

39089 -200.3

39090 -0.07395

39091 -71.65

40090 -2.883

40091 -279.4

40092 -424.8

40093 -532.5

40094 -593.2

40095 -151.2

40096 -711
41094 -0.003199
41095 -116.4
42094 -0.001487
42095 -392.2
42096 -3.128
42097 -754.3
42098 -822.4
42100 -981
43099 -866.1
44099 -0.02943
44100 -12.41
44101 -981.8
44102 -1043
44103 -107.4
44104 -1017
44106 -528.5
45103 -918.7
46104 -9.111
46105 -811.4
46106 -189.4
46107 -507.4
46108 -346.1
46110 -112.7
47109 -179.3
47111 -0.01338
48108 -0.0006807
48110 -2.615
48111 -59.91
48112 -32.25
48113 -21.13
48114 -15.8
48116 -10.77
49113 -0.009386
49115 -12.35
50115 -0.6372
50116 -0.2333
50117 -11.77
50118 -10.69
50119 -9.956
50120 -10.4
50122 -12.77
50123 -1.553
50124 -21
50125 -0.01519
50126 -45.84
51121 -10.37
51123 -12.5

51124	-0.03393
51125	-26.41
51126	-0.002688
52122	-0.1324
52123	-0.0006923
52124	-0.102
52125	-2.467
52126	-0.9772
52128	-152.8
52130	-455.9
53127	-76.95
53129	-255.2
53131	-0.2765
54128	-1.101
54129	-0.004769
54130	-2.888
54131	-733.4
54132	-1045
54133	-0.02523
54134	-1501
54136	-1415
55133	-1350
55134	-15.2
55135	-1480
55136	-0.1605
55137	-1336
56134	-1.726
56135	-0.004795
56136	-28.79
56137	-13.69
56138	-1276
56140	-4.358
57138	-0.01452
57139	-1192
57140	-0.6604
58139	-0.004516
58140	-1133
58141	-76.78
58142	-1018
58144	-578.8
59141	-1005
59143	-5.234
60142	-2.576
60143	-926.8
60144	-257.2
60145	-665.8
60146	-576.9
60147	-0.8997

60148 -386.7
60150 -234.9
61147 -398.9
61148 -0.007476
62147 -40.38
62148 -11.32
62149 -276.1
62150 -16.69
62151 -165
62152 -155.9
62154 -63.16
63151 -0.4976
63152 -0.0111
63153 -95.68
63154 -6.171
63155 -46.17
63156 -0.3119
64152 -0.006463
64154 -0.1648
64155 -2.822
64156 -37.07
64157 -24.46
64158 -16.5
64160 -4.085
65159 -8.656
65160 -0.1454
66160 -0.2801
66161 -2.197
66162 -1.582
66163 -0.7273
66164 -0.4377
67165 -0.2049
68166 -0.1379
68167 -0.0603
68168 -0.02152
68170 -0.002418
92234 -0.2526
92235.17c -2041
92236 -76.7
92237 -0.01288
92238.17c -936900
93236 -0.001106
93237 -102.5
94237 -0.003352
94238 -22.26
94239.17c -234200
94240 -92660
94241 -18230

94242	-5240
94244	-0.005569
95241	-564.1
95242	-1.86
95243	-126.6
96242	-10.97
96243	-0.09684
96244	-5.545
96245	-0.0766
96246	-0.0005938
m404	92235.17c -0.0015732 \$ -0.00157
	92238.17c -0.6357256 \$ -0.63568
	94239.17c -0.1678273 \$ -0.16786
	94240.60c -0.0602681 \$ -0.06028
	94241.60c -0.0129404 \$ -0.01294
	94242.60c -0.0033350 \$ -0.00334
	8016.60c -0.1183303 \$ -0.11833 core II fuel
c	Material 501 is from the output of BOC-2
m501	6012 -0.05483
	6013 -1.764
	7015 -0.003023
	8016 -105100
	8017 -0.001511
	30070 -0.0002267
	31069 -0.000106
	31071 -0.000658
	32072 -0.001626
	32073 -0.00262
	32074 -0.006223
	32076 -0.03216
	33075 -0.01076
	34076 -0.0001749
	34077 -0.06928
	34078 -0.145
	34079 -0.3302
	34080 -0.5305
	34082 -1.54
	35081 -0.8165
	36082 -0.01566
	36083 -2.546
	36084 -5.218
	36085 -1.024
	36086 -7.979
	37085 -3.869
	37086 -0.0006188
	37087 -10.41
	38086 -0.03627
	38087 -0.0001231

38088	-13.71
38089	-1.914
38090	-21.54
39089	-16.55
39090	-0.005596
39091	-3.41
40090	-0.3549
40091	-22.35
40092	-29.2
40093	-34.3
40094	-35.41
40095	-6.114
40096	-42.4
41095	-4.729
42095	-28.01
42096	-0.3558
42097	-41.77
42098	-45.13
42100	-52.25
43099	-46.65
44099	-0.001612
44100	-0.9355
44101	-49.29
44102	-51.98
44103	-3.68
44104	-44.19
44106	-17.64
45103	-45.19
46104	-0.8033
46105	-34.33
46106	-9.223
46107	-17.45
46108	-10.37
46110	-3.085
47109	-4.886
47111	-0.0003279
48110	-0.09612
48111	-1.667
48112	-0.9956
48113	-0.6928
48114	-0.5739
48116	-0.445
49113	-0.0002025
49115	-0.4638
50115	-0.02431
50116	-0.01452
50117	-0.4741
50118	-0.4616

50119	-0.4452
50120	-0.4378
50122	-0.5247
50123	-0.04483
50124	-0.8403
50125	-0.0004897
50126	-1.657
51121	-0.4252
51123	-0.5385
51124	-0.00126
51125	-0.9753
52122	-0.00666
52124	-0.005407
52125	-0.1323
52126	-0.0368
52127	-0.0006337
52128	-6.143
52130	-23.2
53127	-3.102
53129	-12.41
53131	-0.0109
54128	-0.06582
54129	-0.0002805
54130	-0.1581
54131	-38
54132	-57.8
54133	-0.001039
54134	-84.37
54136	-76.06
55133	-73.48
55134	-1.364
55135	-78.27
55136	-0.006852
55137	-70.73
56134	-0.2376
56135	-0.0006542
56136	-1.481
56137	-1.068
56138	-69.74
56140	-0.1726
57138	-0.0002782
57139	-64.84
57140	-0.02615
58140	-63.09
58141	-2.957
58142	-56.4
58144	-28.72
59141	-57.29

59143 -0.2081
60142 -0.3297
60143 -51.57
60144 -21.07
60145 -38.99
60146 -35.76
60147 -0.03647
60148 -22.02
60150 -12.84
61147 -20.91
61148 -0.0004962
62147 -3.43
62148 -1.298
62149 -14.7
62150 -1.812
62151 -7.858
62152 -7.409
62154 -2.719
63151 -0.0342
63152 -0.001616
63153 -4.468
63154 -0.4465
63155 -1.74
63156 -0.009662
64152 -0.0008996
64154 -0.01781
64155 -0.1523
64156 -1.412
64157 -0.8458
64158 -0.5598
64160 -0.1117
65159 -0.2795
65160 -0.004323
66160 -0.01169
66161 -0.05567
66162 -0.04374
66163 -0.0196
66164 -0.01198
67165 -0.005454
68166 -0.004128
68167 -0.001853
68168 -0.000712
69169 -0.000174
92234 -0.0448
92235.17c -1589
92236 -86.08
92237 -0.001586
92238.17c -762600

93237	-21.93
94238	-1.569
94239.17c	-14900
94240	-460.6
94241	-12.41
94242	-0.1468
95241	-0.2471
95242	-0.0003901
95243	-0.001527
96242	-0.002198 \$ axial blanket
m502	92235.17c -0.00218
	92238.17c -0.87932
	8016.60c -0.11850 \$ axial blanket
c	Material 503 is from the output of BOC-2
m503	6012 -0.03146
	6013 -1.203
	7015 -0.001738
	8016 -120600
	8017 -0.0008497
	31071 -0.0002149
	32072 -0.0005094
	32073 -0.0009352
	32074 -0.002091
	32076 -0.01362
	33075 -0.004096
	34077 -0.02926
	34078 -0.06328
	34079 -0.137
	34080 -0.2201
	34082 -0.6819
	35081 -0.3474
	36082 -0.00381
	36083 -1.179
	36084 -2.506
	36085 -0.4749
	36086 -3.819
	37085 -1.781
	37086 -0.0002544
	37087 -4.964
	38086 -0.008543
	38088 -6.561
	38089 -1.562
	38090 -10.4
	39089 -7.314
	39090 -0.002703
	39091 -2.685
	40090 -0.09952
	40091 -9.589

40092	-13.61
40093	-15.77
40094	-15.75
40095	-4.346
40096	-19.05
41095	-3.25
42095	-9.807
42096	-0.05691
42097	-18.37
42098	-19.55
42100	-22.56
43099	-20.37
44099	-0.0006873
44100	-0.1803
44101	-20.99
44102	-21.57
44103	-2.47
44104	-17.5
44106	-7.399
45103	-17.85
46104	-0.1201
46105	-13.85
46106	-2.317
46107	-6.221
46108	-3.271
46110	-0.9511
47109	-1.521
47111	-0.0001755
48110	-0.01372
48111	-0.508
48112	-0.323
48113	-0.2384
48114	-0.2028
48116	-0.1752
49115	-0.1723
50115	-0.008889
50116	-0.002444
50117	-0.18
50118	-0.1834
50119	-0.1812
50120	-0.1738
50122	-0.2034
50123	-0.01888
50124	-0.3103
50125	-0.0002929
50126	-0.5673
51121	-0.1672
51123	-0.2095

51124 -0.0004295
51125 -0.3636
52122 -0.001424
52124 -0.001205
52125 -0.02931
52126 -0.008259
52128 -2.235
52130 -9.651
53127 -1.143
53129 -5.108
53131 -0.00809
54128 -0.0108
54130 -0.03369
54131 -16.23
54132 -24.91
54133 -0.0007959
54134 -37.09
54136 -32.82
55133 -32.25
55134 -0.3188
55135 -33.7
55136 -0.003037
55137 -30.57
56134 -0.03503
56136 -0.3588
56137 -0.278
56138 -30.32
56140 -0.131
57139 -28.34
57140 -0.01985
58140 -27.61
58141 -2.194
58142 -24.64
58144 -15.7
59141 -24.35
59143 -0.1597
60142 -0.05801
60143 -23.21
60144 -6.557
60145 -17.73
60146 -16.13
60147 -0.02822
60148 -9.855
60150 -5.679
61147 -10.37
61148 -0.0001972
62147 -0.9791
62148 -0.2864

62149 -6.845
62150 -0.4442
62151 -3.537
62152 -2.858
62154 -1.099
63151 -0.009924
63152 -0.0001782
63153 -1.91
63154 -0.1052
63155 -0.6934
63156 -0.005789
64154 -0.002712
64155 -0.03872
64156 -0.4945
64157 -0.3008
64158 -0.1836
64160 -0.03465
65159 -0.09427
65160 -0.001038
66160 -0.001765
66161 -0.01674
66162 -0.01248
66163 -0.005221
66164 -0.00321
67165 -0.001501
68166 -0.001135
68167 -0.000562
68168 -0.0002193
92234 -0.02438
92235.17c -2016
92236 -51.32
92237 -0.001617
92238.17c -885500
93237 -12.33
94238 -0.4195
94239.17c -8822
94240 -138.7
94241 -2.072
94242 -0.01142 \$ axial blanket

c Material 511 is from the output of BOC-2

m511 6012 -0.04317

6013 -1.413
7015 -0.002378
8016 -124500
8017 -0.001112
30070 -0.0001265
31071 -0.0003623
32072 -0.0008583

32073	-0.001531
32074	-0.003467
32076	-0.02021
33075	-0.006206
34077	-0.04346
34078	-0.09231
34079	-0.2064
34080	-0.3334
34082	-1.009
35081	-0.5231
36082	-0.00682
36083	-1.712
36084	-3.583
36085	-0.679
36086	-5.519
37085	-2.598
37086	-0.0002605
37087	-7.163
38086	-0.01581
38088	-9.385
38089	-1.237
38090	-14.84
39089	-11.4
39090	-0.003856
39091	-2.222
40090	-0.252
40091	-15.54
40092	-19.77
40093	-23.04
40094	-23.35
40095	-3.796
40096	-28.24
41095	-2.954
42095	-18.91
42096	-0.1592
42097	-27.46
42098	-29.46
42100	-34.07
43099	-30.42
44099	-0.001054
44100	-0.4009
44101	-31.85
44102	-33.11
44103	-2.18
44104	-27.32
44106	-10.28
45103	-29.08
46104	-0.3461

46105 -21.4
46106 -5.46
46107 -10.1
46108 -5.633
46110 -1.662
47109 -2.649
47111 -0.0001684
48110 -0.03394
48111 -0.8925
48112 -0.5488
48113 -0.3958
48114 -0.3324
48116 -0.2737
49115 -0.2772
50115 -0.01436
50116 -0.005709
50117 -0.2856
50118 -0.2866
50119 -0.2718
50120 -0.2713
50122 -0.319
50123 -0.0226
50124 -0.4861
50125 -0.0002471
50126 -0.8998
51121 -0.2615
51123 -0.3312
51124 -0.0005251
51125 -0.543
52122 -0.002731
52124 -0.002425
52125 -0.07698
52126 -0.01543
52127 -0.0003431
52128 -3.518
52130 -14.6
53127 -1.764
53129 -7.386
53131 -0.006332
54128 -0.02342
54130 -0.06321
54131 -24.03
54132 -36.66
54133 -0.000613
54134 -54.25
54136 -49.02
55133 -47.42
55134 -0.5847

55135 -49.88
55136 -0.002993
55137 -45.02
56134 -0.1041
56135 -0.0001982
56136 -0.6642
56137 -0.7033
56138 -44.57
56140 -0.105
57139 -42.04
57140 -0.0159
58140 -41.34
58141 -1.796
58142 -36.25
58144 -18.79
59141 -37.53
59143 -0.1275
60142 -0.1407
60143 -34.55
60144 -14.27
60145 -26.17
60146 -23.71
60147 -0.02253
60148 -14.55
60150 -8.487
61147 -14.23
61148 -0.0002246
62147 -2.472
62148 -0.5931
62149 -10.25
62150 -0.8325
62151 -5.227
62152 -4.485
62154 -1.699
63151 -0.02463
63152 -0.0007274
63153 -2.899
63154 -0.1953
63155 -1.053
63156 -0.005114
64152 -0.0004248
64154 -0.008001
64155 -0.09867
64156 -0.8022
64157 -0.4776
64158 -0.296
64160 -0.05989
65159 -0.1502

65160	-0.001494
66160	-0.004155
66161	-0.029
66162	-0.02113
66163	-0.009649
66164	-0.005828
67165	-0.002783
68166	-0.002047
68167	-0.0009631
68168	-0.0002622
92234	-0.0353
92235.17c	-2027
92236	-66.29
92237	-0.001143
92238.17c	-911400
93237	-17.08
94238	-0.7946
94239.17c	-11420
94240	-220.9
94241	-4.099
94242	-0.0293
95241	-0.08086
96242	-0.0004389
	\$ axial blanket
m512	92235.17c -0.00218
	92238.17c -0.87932
	8016.60c -0.11850
	\$ axial blanket
c	Material 513 is from the output of BOC-2
m513	6012 -0.02186
	6013 -0.8553
	7015 -0.001204
	8016 -116700
	8017 -0.0005618
	31071 -0.0001197
	32072 -0.0002706
	32073 -0.0005514
	32074 -0.00118
	32076 -0.008378
	33075 -0.00233
	34077 -0.01797
	34078 -0.03919
	34079 -0.08403
	34080 -0.1363
	34082 -0.4328
	35081 -0.2174
	36082 -0.001711
	36083 -0.7581
	36084 -1.629
	36085 -0.3031

36086 -2.501
37085 -1.144
37086 -0.000113
37087 -3.238
38086 -0.003832
38088 -4.255
38089 -0.9934
38090 -6.788
39089 -4.755
39090 -0.001764
39091 -1.729
40090 -0.06547
40091 -6.296
40092 -8.803
40093 -10.14
40094 -10.05
40095 -2.717
40096 -12.25
41095 -2.046
42095 -6.334
42096 -0.02588
42097 -11.73
42098 -12.45
42100 -14.35
43099 -12.89
44099 -0.0004353
44100 -0.07985
44101 -13.26
44102 -13.56
44103 -1.502
44104 -10.77
44106 -4.397
45103 -11.28
46104 -0.05486
46105 -8.5
46106 -1.372
46107 -3.62
46108 -1.82
46110 -0.5206
47109 -0.8384
48110 -0.004928
48111 -0.2757
48112 -0.182
48113 -0.138
48114 -0.1191
48116 -0.107
49115 -0.1031
50115 -0.005289

50116 -0.001061
50117 -0.1087
50118 -0.1133
50119 -0.1108
50120 -0.1072
50122 -0.1235
50123 -0.009475
50124 -0.18
50125 -0.0001455
50126 -0.3092
51121 -0.1023
51123 -0.1283
51124 -0.0001991
51125 -0.2034
52122 -0.0006203
52124 -0.0005866
52125 -0.01689
52126 -0.003641
52128 -1.285
52130 -5.996
53127 -0.6479
53129 -2.957
53131 -0.004705
54128 -0.004641
54130 -0.01428
54131 -10.03
54132 -15.43
54133 -0.0004686
54134 -23.16
54136 -20.64
55133 -20.14
55134 -0.1411
55135 -20.95
55136 -0.001362
55137 -19.01
56134 -0.01561
56136 -0.1639
56137 -0.1747
56138 -18.88
56140 -0.07999
57139 -17.88
57140 -0.01212
58140 -17.59
58141 -1.344
58142 -15.4
58144 -10.08
59141 -15.46
59143 -0.09767

60142 -0.02586
60143 -14.83
60144 -4.226
60145 -11.43
60146 -10.33
60147 -0.01746
60148 -6.289
60150 -3.639
61147 -6.758
62147 -0.6499
62148 -0.1324
62149 -4.564
62150 -0.2084
62151 -2.257
62152 -1.742
62154 -0.6694
63151 -0.006515
63153 -1.208
63154 -0.04745
63155 -0.4218
63156 -0.003146
64154 -0.001232
64155 -0.02419
64156 -0.2853
64157 -0.1703
64158 -0.09889
64160 -0.01882
65159 -0.0504
65160 -0.0003943
66160 -0.0006834
66161 -0.008707
66162 -0.005998
66163 -0.002554
66164 -0.001563
67165 -0.0007649
68166 -0.0005635
68167 -0.0002928
92234 -0.01744
92235.17c -2011
92236 -34.54
92237 -0.001109
92238.17c -859900
93237 -8.731
94238 -0.2061
94239.17c -5924
94240 -63.95
94241 -0.7165
94242 -0.002649 \$ axial blanket

m514	92235.17c	-0.00218
	92238.17c	-0.87932
	8016.60c	-0.11850
		\$ axial blanket
m524	92235.17c	-0.00218
	92238.17c	-0.87932
	8016.60c	-0.11850
		\$ axial blanket
c	Material 600 is from the output of BOC-2	
m600	6012	-0.1795
	6013	-5.883
	7015	-0.009862
	8016	-1128000
	8017	-0.00706
	31071	-0.001016
	32072	-0.003617
	32073	-0.006839
	32074	-0.01542
	32076	-0.1051
	33075	-0.03016
	34077	-0.2186
	34078	-0.452
	34079	-0.9642
	34080	-1.597
	34082	-4.602
	35081	-2.454
	36082	-0.02443
	36083	-7.842
	36084	-16.27
	36085	-3.168
	36086	-25.57
	37085	-12.18
	37087	-33.24
	38086	-0.05052
	38088	-43.79
	38089	-5.617
	38090	-69.22
	39089	-52.99
	39090	-0.01798
	39091	-10.06
	40090	-1.187
	40091	-71.96
	40092	-90.26
	40093	-104.4
	40094	-105.6
	40095	-16.74
	40096	-125.9
	41095	-13.05
	42095	-85.8
	42096	-0.538

42097	-122.8
42098	-130.2
42100	-150.1
43099	-135.3
44099	-0.00469
44100	-1.247
44101	-139.4
44102	-142.3
44103	-9.143
44104	-114.7
44106	-42.1
45103	-124.4
46104	-1.062
46105	-90.07
46106	-22.04
46107	-41.42
46108	-22.57
46110	-6.759
47109	-10.74
48110	-0.1014
48111	-3.645
48112	-2.262
48113	-1.657
48114	-1.401
48116	-1.179
49115	-1.187
50115	-0.06122
50116	-0.01947
50117	-1.231
50118	-1.232
50119	-1.162
50120	-1.175
50122	-1.378
50123	-0.09213
50124	-2.093
50126	-3.805
51121	-1.134
51123	-1.436
51124	-0.001674
51125	-2.304
52122	-0.00853
52124	-0.008023
52125	-0.334
52126	-0.05403
52127	-0.00145
52128	-15.26
52130	-62.69
53127	-7.717

53129 -32.77
53131 -0.02695
54128 -0.0767
54130 -0.1989
54131 -105.2
54132 -158.8
54133 -0.002596
54134 -238.9
54136 -214.6
55133 -210.2
55134 -1.852
55135 -218.6
55136 -0.01031
55137 -197.9
56134 -0.3328
56136 -2.355
56137 -3.125
56138 -197.6
56140 -0.4558
57139 -186.9
57140 -0.06907
58140 -183.4
58141 -7.812
58142 -161.3
58144 -83.94
59141 -167.4
59143 -0.559
60142 -0.4383
60143 -155.7
60144 -64.24
60145 -116.9
60146 -104.2
60147 -0.09751
60148 -63.7
60150 -36.55
61147 -63.47
62147 -11.19
62148 -1.933
62149 -45.22
62150 -2.72
62151 -22.79
62152 -18.55
62154 -7.037
63151 -0.1104
63152 -0.00233
63153 -12.45
63154 -0.605
63155 -4.37

63156	-0.02028
64154	-0.02499
64155	-0.4171
64156	-3.253
64157	-1.951
64158	-1.182
64160	-0.2422
65159	-0.6105
65160	-0.004164
66160	-0.0111
66161	-0.1167
66162	-0.08222
66163	-0.03804
66164	-0.0228
67165	-0.01102
68166	-0.008048
68167	-0.002752
92234	-0.1497
92235.17c	-19110
92236	-426
92237	-0.005009
92238.17c	-8296000
93237	-73.39
94238	-2.342
94239.17c	-68130
94240	-933.8
94241	-14.46
94242	-0.07244
95241	-0.3203 \$ radial blanket
m601	92235.17c -0.00218
	92238.17c -0.87932
	8016.60c -0.11850 \$ radial blanket
m602	92235.17c -0.00218
	92238.17c -0.87932
	8016.60c -0.11850 \$ radial blanket
m604	92235.17c -0.0015732 \$ -0.00157
	92238.17c -0.6357256 \$ -0.63568
	94239.17c -0.1678273 \$ -0.16786
	94240.60c -0.0602681 \$ -0.06028
	94241.60c -0.0129404 \$ -0.01294
	94242.60c -0.0033350 \$ -0.00334
	8016.60c -0.1183303 \$ -0.11833 core II fuel
m7	26054.62c 0.49285e-03
	26056.62c 0.77367e-02
	26057.62c 0.17867e-03
	26058.62c 0.23778e-04
	24050.62c 0.83368e-04
	24052.62c 0.16077e-02

	24053.62c 0.18230e-03	
	24054.62c 0.45377e-04	
	28058.62c 0.12397e-02	
	28060.62c 0.47752e-03	
	28061.62c 0.20759e-04	
	28062.62c 0.66175e-04	
	28064.62c 0.16862e-04	
	42000.66c 0.16722e-03	
	6012.50c 0.15150e-01	
	11023.62c 0.55700e-02	
	1001.62c 0.95384e-20	
	14000.60c 0.15859e-03	
	25055.62c 0.25943e-03	
	5010.66c 0.11915e-01	
	5011.66c 0.48262e-01	\$ B4C Shld & SS bot
m8	26054.62c 0.25312e-02	
	26056.62c 0.39734e-01	
	26057.62c 0.91763e-03	
	26058.62c 0.12212e-03	
	24050.62c 0.54017e-04	
	24052.62c 0.10417e-02	
	24053.62c 0.11812e-03	
	24054.62c 0.29402e-04	
	28058.62c 0.52485e-02	
	28060.62c 0.20217e-02	
	28061.62c 0.87891e-04	
	28062.62c 0.28017e-03	
	28064.62c 0.71392e-04	
	42000.66c 0.96251e-03	
	6012.50c 0.83109e-04	
	11023.62c 0.64104e-02	
	1001.62c 0.95384e-20	
	14000.60c 0.65748e-03	
	25055.62c 0.11766e-02	\$ SS Reflector
m9	26054.62c 0.54798e-03	
	26056.62c 0.86021e-02	
	26057.62c 0.19866e-03	
	26058.62c 0.26438e-04	
	24050.62c 0.92692e-04	
	24052.62c 0.17875e-02	
	24053.62c 0.20268e-03	
	24054.62c 0.50453e-04	
	28058.62c 0.13783e-02	
	28060.62c 0.53091e-03	
	28061.62c 0.23080e-04	
	28062.62c 0.73574e-04	
	28064.62c 0.18748e-04	
	42000.66c 0.18592e-03	

6012.50c 0.28365e-04	
11023.62c 0.19719e-01	
1001.62c 0.95384e-20	
14000.60c 0.17636e-03	
25055.62c 0.28844e-03	\$ CSR/DSR follower
m10 26054.62c 0.62343e-03	
26056.62c 0.97865e-02	
26057.62c 0.22601e-03	
26058.62c 0.30078e-04	
24050.62c 0.10546e-03	
24052.62c 0.20336e-02	
24053.62c 0.23060e-03	
24054.62c 0.57401e-04	
28058.62c 0.15682e-02	
28060.62c 0.60405e-03	
28061.62c 0.26260e-04	
28062.62c 0.83709e-04	
28064.62c 0.21330e-04	
42000.66c 0.21153e-03	
5010.66c 0.15779e-01	
5011.66c 0.12131e-01	
6012.50c 0.69823e-02	
11023.62c 0.12380e-01	
1001.62c 0.95384e-20	
14000.60c 0.20061e-03	
25055.62c 0.32817e-03	\$ CSR/DSR homo
m11 26054.62c 0.19242e-02	
26056.62c 0.30205e-01	
26057.62c 0.69757e-03	
26058.62c 0.92834e-04	
24050.62c 0.32548e-03	
24052.62c 0.62765e-02	
24053.62c 0.71170e-03	
24054.62c 0.17716e-03	
28058.62c 0.48398e-02	
28060.62c 0.18643e-02	
28061.62c 0.81046e-04	
28062.62c 0.25835e-03	
28064.62c 0.65832e-04	
42000.66c 0.65284e-03	
6012.50c 0.99601e-04	
11023.62c 0.97907e-02	
1001.62c 0.95384e-20	
14000.60c 0.61914e-03	
25055.62c 0.10128e-02	\$ SA bottom
m12 26054.62c 0.78387e-03	
26056.62c 0.12305e-01	
26057.62c 0.28418e-03	

	26058.62c 0.37819e-04	
	24050.62c 0.13260e-03	
	24052.62c 0.25571e-02	
	24053.62c 0.28995e-03	
	24054.62c 0.72175e-04	
	28058.62c 0.19717e-02	
	28060.62c 0.75950e-03	
	28061.62c 0.33018e-04	
	28062.62c 0.10525e-03	
	28064.62c 0.26820e-04	
	42000.66c 0.26597e-03	
	6012.50c 0.40578e-04	
	11023.62c 0.97126e-02	
	1001.62c 0.99184e-20	
	14000.60c 0.25224e-03	
	25055.62c 0.41263e-03	\$ Core Plenum homog
m13	26054.62c 0.23927e-02	
	26056.62c 0.37560e-01	
	26057.62c 0.86743e-03	
	26058.62c 0.11544e-03	
	24050.62c 0.40474e-03	
	24052.62c 0.78049e-02	
	24053.62c 0.88502e-03	
	24054.62c 0.22030e-03	
	28058.62c 0.60183e-02	
	28060.62c 0.23182e-02	
	28061.62c 0.10078e-03	
	28062.62c 0.32126e-03	
	28064.62c 0.81863e-04	
	42000.66c 0.81182e-03	
	6012.50c 0.12386e-03	
	11023.62c 0.64104e-02	
	1001.62c 0.95384e-20	
	14000.60c 0.76992e-03	
	25055.62c 0.12595e-02	\$ RPBSS
m14	26054.62c 0.29254e-03	
	26056.62c 0.45923e-02	
	26057.62c 0.10606e-03	
	26058.62c 0.14114e-04	
	24050.62c 0.49438e-04	
	24052.62c 0.95424e-03	
	24053.62c 0.10820e-03	
	24054.62c 0.26934e-04	
	28058.62c 0.73597e-03	
	28060.62c 0.28349e-03	
	28061.62c 0.12324e-04	
	28062.62c 0.39286e-04	
	28064.62c 0.10011e-04	

	42000.66c 0.91980e-04	
	6012.50c 0.10890e-04	
	11023.62c 0.21539e-01	
	14000.60c 0.94113e-04	
	25055.62c 0.15398e-03	\$ RBPT/RFTSS1
m15	6012.50c 0.83867e-04	
	14000.60c 0.52075e-03	
	25055.62c 0.85201e-03	
	26054.62c 0.16187e-02	
	26056.62c 0.25410e-01	
	26057.62c 0.58684e-03	
	26058.62c 0.78097e-04	
	24050.62c 0.27381e-03	
	24052.62c 0.52802e-02	
	24053.62c 0.59873e-03	
	24054.62c 0.14904e-03	
	28058.62c 0.40722e-02	
	28060.62c 0.15686e-02	
	28061.62c 0.68192e-04	
	28062.62c 0.21738e-03	
	28064.62c 0.55391e-04	
	42000.66c 0.54889e-03	
	11023.62c 0.11882e-01	\$ Core-SS
m16	5010.66c 0.47121e-02	
	5011.66c 0.19329e-01	
	6012.50c 0.60415e-02	
	14000.60c 0.19450e-03	
	25055.62c 0.31822e-03	
	42000.66c 0.20501e-03	
	11023.62c 0.11882e-01	
	26054.62c 0.60458e-03	
	26056.62c 0.94907e-02	
	26057.62c 0.21918e-03	
	26058.62c 0.29169e-04	
	24050.62c 0.10227e-03	
	24052.62c 0.19721e-02	
	24053.62c 0.22362e-03	
	24054.62c 0.55664e-04	
	28058.62c 0.15209e-02	
	28060.62c 0.58586e-03	
	28061.62c 0.25469e-04	
	28062.62c 0.81189e-04	
	28064.62c 0.20688e-04	\$ Core B4C
m17	26054.62c 0.49285e-03	
	26056.62c 0.77367e-02	
	26057.62c 0.17867e-03	
	26058.62c 0.23778e-04	
	24050.62c 0.83368e-04	

24052.62c 0.16077e-02	
24053.62c 0.18230e-03	
24054.62c 0.45377e-04	
28058.62c 0.12397e-02	
28060.62c 0.47752e-03	
28061.62c 0.20759e-04	
28062.62c 0.66175e-04	
28064.62c 0.16862e-04	
42000.66c 0.16722e-03	
6012.50c 0.25512e-04	
11023.62c 0.55700e-02	
1001.62c 0.95384e-20	
14000.60c 0.15859e-03	
25055.62c 0.25943e-03	\$ SHPLenum
m18 26054.62c 0.20590e-02	
26056.62c 0.32322e-01	
26057.62c 0.74646e-03	
26058.62c 0.99340e-04	
24050.62c 0.34829e-03	
24052.62c 0.67164e-02	
24053.62c 0.76158e-03	
24054.62c 0.18957e-03	
28058.62c 0.51790e-02	
28060.62c 0.19949e-02	
28061.62c 0.86726e-04	
28062.62c 0.27646e-03	
28064.62c 0.70445e-04	
42000.66c 0.69860e-03	
6012.50c 0.10658e-03	
11023.62c 0.88178e-02	
1001.62c 0.95384e-20	
14000.60c 0.66254e-03	
25055.62c 0.10838e-02	\$ SHLD SS top
m19 92235.17c 0.187309e-04	
92238.17c 0.747364e-02	
26000.55c 0.132470e-01	
24000.50c 0.301425e-02	
28000.50c 0.286071e-02	
8016.60c 0.149847e-01	
6012.50c 0.400788e-04	
11023.62c 0.903609e-02	
14000.60c 0.249138e-03	
25055.62c 0.407556e-03	\$ diluent homog
m20 26054.62c 0.63886e-03	
26056.62c 0.10029e-01	
26057.62c 0.23161e-03	
26058.62c 0.30823e-04	
24050.62c 0.10807e-03	

24052.62c	0.20840e-02	
24053.62c	0.23631e-03	
24054.62c	0.58822e-04	
28058.62c	0.16070e-02	
28060.62c	0.61899e-03	
28061.62c	0.26910e-04	
28062.62c	0.85781e-04	
28064.62c	0.21858e-04	
42000.66c	0.21676e-03	
6012.50c	0.33070e-04	
11023.62c	0.66755e-02	
1001.62c	0.95384e-20	
14000.60c	0.20557e-03	
25055.62c	0.33629e-03	\$ Blanket plenum
m21	5010.66c 0.52	
	5011.66c 0.28	
	6012.50c 0.20	\$ pinwise CSRmid/DSR
m22	5010.66c 0.1592	
	5011.66c 0.6408	
	6012.50c 0.2000	\$ pinwise CSRtop/bot
m23	92235.17c 0.187309e-04	
	92238.17c 0.747364e-02	
	26000.55c 0.132470e-01	
	24000.50c 0.301425e-02	
	28000.50c 0.286071e-02	
	8016.60c 0.149847e-01	
	6012.50c 0.400788e-04	
	11023.62c 0.903609e-02	
	14000.60c 0.249138e-03	
	25055.62c 0.407556e-03	\$ diluent homog

APPENDIX B

Table XII

Mass and activity values for every isotope printed in MCNPX output for “three cycles”

radial blanket material at End-of-Life.

Isotope	ZAID	Mass (g)	Activity (Ci)
C-12	6012	2.677E-01	0.00E+00
C-13	6013	8.193E+00	0.00E+00
N-15	7015	1.471E-02	0.00E+00
O-16	8016	1.128E+06	0.00E+00
O-17	8017	1.018E-02	0.00E+00
Ga-71	31071	2.362E-03	0.00E+00
Ge-72	32072	6.952E-03	0.00E+00
Ge-73	32073	1.224E-02	0.00E+00
Ge-74	32074	2.832E-02	0.00E+00
Ge-76	32076	1.711E-01	0.00E+00
As-75	33075	5.204E-02	0.00E+00
Se-77	34077	3.582E-01	0.00E+00
Se-78	34078	7.336E-01	0.00E+00
Se-79	34079	1.611E+00	2.21E-01
Se-80	34080	2.653E+00	0.00E+00
Se-82	34082	7.495E+00	2.35E-16
Br-81	35081	4.051E+00	0.00E+00
Kr-82	36082	5.439E-02	0.00E+00
Kr-83	36083	1.253E+01	0.00E+00
Kr-84	36084	2.560E+01	0.00E+00
Kr-85	36085	4.984E+00	1.96E+03
Kr-86	36086	4.020E+01	0.00E+00
Rb-85	37085	1.951E+01	0.00E+00
Rb-86	37086	1.250E-03	1.02E+02
Rb-87	37087	5.234E+01	4.49E-06
Sr-86	38086	1.154E-01	0.00E+00
Sr-88	38088	6.880E+01	0.00E+00
Sr-89	38089	5.857E+00	1.70E+05
Sr-90	38090	1.078E+02	1.52E+04

Y-89	39089	8.610E+01	0.00E+00
Y-90	39090	2.802E-02	1.52E+04
Y-91	39091	1.062E+01	2.61E+05
Zr-90	40090	2.693E+00	0.00E+00
Zr-91	40091	1.186E+02	0.00E+00
Zr-92	40092	1.438E+02	0.00E+00
Zr-93	40093	1.676E+02	4.22E-01
Zr-94	40094	1.718E+02	0.00E+00
Zr-95	40095	1.858E+01	3.99E+05
Zr-96	40096	2.044E+02	0.00E+00
Nb-95	41095	1.461E+01	5.75E+05
Mo-95	42095	1.545E+02	0.00E+00
Mo-96	42096	1.403E+00	0.00E+00
Mo-97	42097	2.012E+02	0.00E+00
Mo-98	42098	2.150E+02	0.00E+00
Mo-100	42100	2.484E+02	0.00E+00
Tc-99	43099	2.228E+02	3.82E+00
Ru-99	44099	7.932E-03	0.00E+00
Ru-100	44100	3.072E+00	0.00E+00
Ru-101	44101	2.325E+02	0.00E+00
Ru-102	44102	2.403E+02	0.00E+00
Ru-103	44103	1.056E+01	3.41E+05
Ru-104	44104	1.987E+02	0.00E+00
Ru-106	44106	6.540E+01	2.17E+05
Rh-103	45103	2.154E+02	0.00E+00
Pd-104	46104	2.846E+00	0.00E+00
Pd-105	46105	1.547E+02	0.00E+00
Pd-106	46106	5.074E+01	0.00E+00
Pd-107	46107	7.576E+01	3.90E-02
Pd-108	46108	4.352E+01	0.00E+00
Pd-110	46110	1.315E+01	0.00E+00
Ag-109	47109	2.083E+01	0.00E+00
Cd-110	48110	2.974E-01	0.00E+00
Cd-111	48111	7.119E+00	0.00E+00
Cd-112	48112	4.285E+00	0.00E+00
Cd-113	48113	3.057E+00	1.04E-12
Cd-114	48114	2.547E+00	0.00E+00
Cd-116	48116	2.039E+00	0.00E+00
In-115	49115	2.109E+00	1.49E-11
Sn-115	50115	1.095E-01	0.00E+00

Sn-116	50116	5.039E-02	0.00E+00
Sn-117	50117	2.162E+00	0.00E+00
Sn-118	50118	2.122E+00	0.00E+00
Sn-119	50119	1.990E+00	0.00E+00
Sn-120	50120	2.021E+00	0.00E+00
Sn-122	50122	2.398E+00	0.00E+00
Sn-123	50123	1.351E-01	1.11E+03
Sn-124	50124	3.728E+00	0.00E+00
Sn-126	50126	7.052E+00	2.00E-01
Sb-121	51121	1.961E+00	0.00E+00
Sb-123	51123	2.516E+00	0.00E+00
Sb-124	51124	2.798E-03	4.90E+01
Sb-125	51125	3.957E+00	4.15E+03
Te-122	52122	2.044E-02	0.00E+00
Te-124	52124	1.910E-02	0.00E+00
Te-125	52125	8.395E-01	0.00E+00
Te-126	52126	1.227E-01	0.00E+00
Te-127	52127	1.795E-03	4.74E+03
Te-128	52128	2.732E+01	0.00E+00
Te-130	52130	1.058E+02	0.00E+00
I-127	53127	1.376E+01	0.00E+00
I-129	53129	5.575E+01	9.85E-03
I-131	53131	3.006E-02	3.73E+03
Xe-128	54128	1.996E-01	0.00E+00
Xe-130	54130	4.756E-01	0.00E+00
Xe-131	54131	1.752E+02	0.00E+00
Xe-132	54132	2.643E+02	0.00E+00
Xe-133	54133	2.857E-03	5.35E+02
Xe-134	54134	3.925E+02	0.00E+00
Xe-136	54136	3.556E+02	0.00E+00
Cs-133	55133	3.461E+02	0.00E+00
Cs-134	55134	4.081E+00	5.28E+03
Cs-135	55135	3.628E+02	4.18E-01
Cs-136	55136	1.493E-02	1.09E+03
Cs-137	55137	3.262E+02	2.84E+04
Ba-134	56134	1.058E+00	0.00E+00
Ba-135	56135	1.703E-03	0.00E+00
Ba-136	56136	5.194E+00	0.00E+00
Ba-137	56137	7.382E+00	0.00E+00
Ba-138	56138	3.258E+02	0.00E+00

Ba-140	56140	4.995E-01	3.66E+04
La-139	57139	3.072E+02	0.00E+00
La-140	57140	7.569E-02	4.21E+04
Ce-140	58140	3.013E+02	0.00E+00
Ce-141	58141	8.612E+00	2.46E+05
Ce-142	58142	2.652E+02	1.34E-11
Ce-144	58144	1.109E+02	3.53E+05
Pr-141	59141	2.785E+02	0.00E+00
Pr-143	59143	6.100E-01	4.11E+04
Nd-142	60142	1.112E+00	0.00E+00
Nd-143	60143	2.538E+02	0.00E+00
Nd-144	60144	1.298E+02	1.54E-10
Nd-145	60145	1.889E+02	7.77E-12
Nd-146	60146	1.690E+02	0.00E+00
Nd-147	60147	1.059E-01	8.57E+03
Nd-148	60148	1.039E+02	0.00E+00
Nd-150	60150	6.006E+01	0.00E+00
Pm-147	61147	9.474E+01	8.79E+04
Sm-147	62147	2.487E+01	5.71E-07
Sm-148	62148	4.560E+00	1.39E-12
Sm-149	62149	7.240E+01	8.69E-11
Sm-150	62150	6.114E+00	0.00E+00
Sm-151	62151	3.723E+01	9.80E+02
Sm-152	62152	3.245E+01	0.00E+00
Sm-154	62154	1.213E+01	0.00E+00
Eu-151	63151	2.528E-01	0.00E+00
Eu-152	63152	8.457E-03	1.49E+00
Eu-153	63153	2.085E+01	0.00E+00
Eu-154	63154	1.393E+00	3.77E+02
Eu-155	63155	7.393E+00	3.65E+03
Eu-156	63156	2.506E-02	1.38E+03
Gd-152	64152	3.707E-03	8.08E-14
Gd-154	64154	8.140E-02	0.00E+00
Gd-155	64155	1.002E+00	0.00E+00
Gd-156	64156	5.989E+00	0.00E+00
Gd-157	64157	3.581E+00	0.00E+00
Gd-158	64158	2.253E+00	0.00E+00
Gd-160	64160	4.714E-01	0.00E+00
Tb-159	65159	1.154E+00	0.00E+00
Tb-160	65160	8.485E-03	9.58E+01

Dy-160	66160	3.405E-02	0.00E+00
Dy-161	66161	2.311E-01	0.00E+00
Dy-162	66162	1.687E-01	0.00E+00
Dy-163	66163	7.950E-02	0.00E+00
Dy-164	66164	4.767E-02	0.00E+00
Ho-165	67165	2.264E-02	0.00E+00
Er-166	68166	1.663E-02	0.00E+00
Er-167	68167	6.552E-03	0.00E+00
U-234	92234	2.326E-01	1.45E-03
U-235	92235	1.842E+04	3.98E-02
U-236	92236	6.022E+02	3.90E-02
U-237	92237	4.863E-03	3.97E+02
U-238	92238	8.263E+06	2.78E+00
Np-237	93237	1.096E+02	7.73E-02
Pu-238	94238	5.128E+00	8.78E+01
Pu-239	94239	9.713E+04	6.03E+03
Pu-240	94240	1.889E+03	4.29E+02
Pu-241	94241	3.788E+01	3.91E+03
Pu-242	94242	2.804E-01	1.11E-03
Am-241	95241	1.199E+00	4.11E+00
Cm-242	96242	8.033E-03	2.66E+01