

Combination of CDF and D0 W -Boson Mass Measurements

T. Aaltonen^{†,21} V.M. Abazov^{‡,13} B. Abbott^{‡,116} B.S. Acharya^{‡,81} M. Adams^{‡,98} T. Adams^{‡,97} J.P. Agnew^{‡,94}
 G.D. Alexeev^{‡,13} G. Alkhazov^{‡,89} A. Alton^{‡,jj,32} S. Amerio^{†,vv,40} D. Amidei^{†,32} A. Anastassov^{†,w,15} A. Annovi^{†,17}
 J. Antos^{†,12} G. Apollinari^{†,15} J.A. Appel^{†,15} T. Arisawa^{†,53} A. Artikov^{†,13} J. Asaadi^{†,48} W. Ashmanskas^{†,15}
 A. Askew^{‡,97} S. Atkins^{‡,106} B. Auerbach^{†,2} K. Augsten^{‡,63} A. Aurisano^{†,48} C. Avila^{‡,61} F. Azfar^{†,39} F. Badaud^{‡,66}
 W. Badgett^{†,15} T. Bae^{†,25} L. Bagby^{‡,15} B. Baldin^{‡,15} D.V. Bandurin^{‡,97} S. Banerjee^{‡,81} A. Barbaro-Galtieri^{†,26}
 E. Barberis^{‡,107} P. Baringer^{‡,105} V.E. Barnes^{†,44} B.A. Barnett^{†,23} J. Guimaraes da Costa^{†,20} P. Barria^{†,xx,42}
 J.F. Bartlett^{‡,15} P. Bartos^{†,12} U. Bessler^{†,71} M. Bauce^{†,vv,40} V. Bazterra^{‡,98} A. Bean^{‡,105} F. Bedeschi^{†,42}
 D. Beecher^{†,28} M. Begalli^{‡,58} S. Behari^{†,15} L. Bellantoni^{†,15} G. Bellettini^{†,ww,42} J. Bellinger^{†,55}
 D. Benjamin^{†,14} A. Beretvas^{†,15} S.B. Beri^{†,79} G. Bernardi^{†,70} R. Bernhard^{†,75} I. Bertram^{†,92} M. Besanon^{‡,71}
 R. Beuselinck^{‡,93} P.C. Bhat^{‡,15} S. Bhatia^{‡,108} V. Bhatnagar^{†,79} A. Bhatti^{†,46} I. Bizjak^{†,28} K.R. Bland^{†,5}
 G. Blazey^{†,99} S. Blessing^{†,97} K. Bloom^{‡,109} B. Blumenfeld^{†,23} A. Bocci^{†,14} A. Bodek^{†,45} A. Boehnlein^{‡,15}
 D. Boline^{‡,113} E.E. Boos^{‡,87} G. Borissov^{†,92} D. Bortoletto^{†,44} J. Boudreau^{†,43} A. Boveia^{†,11} A. Brandt^{‡,119}
 O. Brandt^{‡,76} L. Brigliadori^{†,uu,6} R. Brock^{‡,33} C. Bromberg^{†,33} A. Bross^{†,15} D. Brown^{‡,70} E. Brucken^{†,21}
 X.B. Bu^{†,15} J. Budagov^{†,13} H.S. Budd^{†,45} M. Buehler^{‡,15} V. Buescher^{‡,77} V. Bunichev^{‡,87} S. Burdin^{†,kk,92}
 K. Burkett^{†,15} G. Busetto^{†,vv,40} P. Bussey^{†,19} C.P. Buszello^{‡,91} P. Butti^{†,ww,42} A. Buzatu^{†,19} A. Calamba^{†,10}
 E. Camacho-Pérez^{‡,84} S. Camarda^{†,4} M. Campanelli^{†,28} F. Canelli^{†,dd,11} B. Carls^{†,22} D. Carlsmith^{†,55} R. Carosi^{†,42}
 S. Carrillo^{†,16} B. Casal^{†,j,9} M. Casarsa^{†,49} B.C.K. Casey^{†,15} H. Castilla-Valdez^{†,84} A. Castro^{†,uu,6} P. Catastini^{†,20}
 S. Caughron^{‡,33} D. Cauz^{†,ccddd,49} V. Cavaliere^{†,22} M. Cavalli-Sforza^{†,4} A. Cerri^{†,e,26} L. Cerrito^{†,r,28}
 S. Chakrabarti^{‡,113} K.M. Chan^{†,103} A. Chandra^{†,121} E. Chapon^{†,71} G. Chen^{‡,105} Y.C. Chen^{†,1} M. Chertok^{†,7}
 G. Chiarelli^{†,42} G. Chlachidze^{†,15} K. Cho^{†,25} S.W. Cho^{‡,83} S. Choi^{‡,83} D. Chokheli^{†,13} B. Choudhary^{‡,80}
 S. Cihangir^{†,15} D. Claes^{†,109} A. Clark^{†,18} C. Clarke^{†,54} J. Clutter^{†,105} M.E. Convery^{†,15} J. Conway^{†,7} M. Cooke^{†,15}
 W.E. Cooper^{†,15} M. Corbo^{†,z,15} M. Corcoran^{†,121} M. Cordelli^{†,17} F. Couderc^{†,71} M.-C. Cousinou^{†,68} C.A. Cox^{†,7}
 D.J. Cox^{†,7} M. Cremonesi^{†,42} D. Cruz^{†,48} J. Cuevas^{†,y,9} R. Culbertson^{†,15} D. Cutts^{†,118} A. Das^{†,95}
 N. d'Ascenzo^{†,v,15} M. Datta^{†,gg,15} G. Davies^{†,93} P. de Barbaro^{†,45} S.J. de Jong^{†,85,86} E. De La Cruz-Burelo^{†,84}
 F. Déliot^{†,71} R. Demina^{†,45} L. Demortier^{†,46} M. Deninno^{†,6} D. Denisov^{†,15} S.P. Denisov^{†,88} M. D'Errico^{†,vv,40}
 S. Desai^{†,15} C. Deterre^{†,mm,76} K. DeV Vaughan^{†,109} F. Devoto^{†,21} A. Di Canto^{†,ww,42} B. Di Ruzza^{†,p,15} H.T. Diehl^{†,15}
 M. Diesburg^{†,15} P.F. Ding^{†,94} J.R. Dittmann^{†,5} A. Dominguez^{†,109} S. Donati^{†,ww,42} M. D'Onofrio^{†,27}
 M. Dorigo^{†,eee,49} A. Driutti^{†,ccddd,49} A. Dubey^{†,80} L.V. Dudko^{†,87} A. Duperrin^{†,68} S. Dutt^{†,79} M. Eads^{†,99}
 K. Ebina^{†,53} R. Edgar^{†,32} D. Edmunds^{†,33} A. Elagin^{†,48} J. Ellison^{†,96} V.D. Elvira^{†,15} Y. Enari^{†,70} R. Erbacher^{†,7}
 S. Errede^{†,22} B. Esham^{†,22} R. Eusebi^{†,48} H. Evans^{†,101} V.N. Evdokimov^{†,88} S. Farrington^{†,39} L. Feng^{†,99}
 T. Ferbel^{†,45} J.P. Fernández Ramos^{†,29} F. Fiedler^{†,77} R. Field^{†,16} F. Filthaut^{†,85,86} W. Fisher^{†,33} H.E. Fisk^{†,15}
 G. Flanagan^{†,t,15} R. Forrest^{†,7} M. Fortner^{†,99} H. Fox^{†,92} M. Franklin^{†,20} J.C. Freeman^{†,15} H. Frisch^{†,11} S. Fuess^{†,15}
 Y. Funakoshi^{†,53} C. Galloni^{†,ww,42} A. Garcia-Bellido^{†,45} J.A. García-González^{†,84} A.F. Garfinkel^{†,44} P. Garosi^{†,xx,42}
 V. Gavrilov^{†,34} W. Geng^{†,68,33} C.E. Gerber^{†,98} H. Gerberich^{†,22} E. Gerchtein^{†,15} Y. Gershtein^{†,110} S. Giagu^{†,47}
 V. Giakoumopoulou^{†,3} K. Gibson^{†,43} C.M. Ginsburg^{†,15} G. Ginter^{†,15,45} N. Giokaris^{†,3} P. Giromini^{†,17}
 G. Giurgiu^{†,23} V. Glagolev^{†,13} D. Glenzinski^{†,15} M. Gold^{†,35} D. Goldin^{†,48} A. Golossanov^{†,15} G. Golovanov^{†,13}
 G. Gomez^{†,9} G. Gomez-Ceballos^{†,30} M. Goncharov^{†,30} O. González López^{†,29} I. Gorelov^{†,35} A.T. Goshaw^{†,14}
 K. Goulianos^{†,46} E. Gramellini^{†,6} P.D. Grannis^{†,113} S. Greder^{†,72} H. Greenlee^{†,15} G. Grenier^{†,73} S. Grinstein^{†,4}
 Ph. Gris^{†,66} J.-F. Grivaz^{†,69} A. Grohsjean^{†,ll,71} C. Grosso-Pilcher^{†,11} R.C. Group^{†,52,15} S. Grünendahl^{†,15}
 M.W. Grünewald^{†,82} T. Guillemín^{†,69} G. Gutierrez^{†,15} P. Gutierrez^{†,116} S.R. Hahn^{†,15} J. Haley^{†,107}
 J.Y. Han^{†,45} L. Han^{†,60} F. Happacher^{†,17} K. Hara^{†,50} K. Harder^{†,94} M. Hare^{†,51} A. Harel^{†,45} R.F. Harr^{†,54}
 T. Harrington-Taber^{†,m,15} K. Hatakeyama^{†,5} J.M. Hauptman^{†,104} C. Hays^{†,39} J. Hays^{†,93} T. Head^{†,94}
 T. Hebbeker^{†,74} D. Hedin^{†,99} H. Hegab^{†,117} J. Heinrich^{†,41} A.P. Heinson^{†,96} U. Heintz^{†,118} C. Hensel^{†,76}
 I. Heredia-De La Cruz^{†,mm,84} M. Herndon^{†,55} K. Herner^{†,15} G. Hesketh^{†,oo,94} M.D. Hildreth^{†,103} R. Hirosky^{†,122}
 T. Hoang^{†,97} J.D. Hobbs^{†,113} A. Hocker^{†,15} B. Hoeneisen^{†,65} J. Hogan^{†,121} M. Hohlfield^{†,77} J.L. Holzbauer^{†,108}
 Z. Hong^{†,48} W. Hopkins^{†,f,15} S. Hou^{†,1} I. Howley^{†,119} Z. Hubacek^{†,63,71} R.E. Hughes^{†,36} U. Husemann^{†,56}
 M. Hussein^{†,bb,33} J. Huston^{†,33} V. Hynek^{†,63} I. Iashvili^{†,112} Y. Ilchenko^{†,120} R. Illingworth^{†,15} G. Introzzi^{†,zzaaa,42}
 M. Iori^{†,bbb,47} A.S. Ito^{†,15} A. Ivanov^{†,o,7} S. Jabeen^{†,118} M. Jaffré^{†,69} E. James^{†,15} D. Jang^{†,10} A. Jayasinghe^{†,116}
 B. Jayatilaka^{†,15} E.J. Jeon^{†,25} M.S. Jeong^{†,83} R. Jesik^{†,93} P. Jiang^{†,60} S. Jindariani^{†,15} K. Johns^{†,95} E. Johnson^{†,33}
 M. Johnson^{†,15} A. Jonckheere^{†,15} M. Jones^{†,44} P. Jonsson^{†,93} K.K. Joo^{†,25} J. Joshi^{†,96} S.Y. Jun^{†,10} A.W. Jung^{†,15}
 T.R. Junk^{†,15} A. Juste^{†,90} E. Kajfasz^{†,68} M. Kambeitz^{†,24} T. Kamon^{†,25,48} P.E. Karchin^{†,54} D. Karmanov^{†,87}
 A. Kasmi^{†,5} Y. Kato^{†,n,38} I. Katsanos^{†,109} R. Kehoe^{†,120} S. Kermiche^{†,68} W. Ketchum^{†,hh,11} J. Keung^{†,41}

N. Khalatyan^{‡,15} A. Khanov^{‡,117} A. Kharchilava^{‡,112} Y.N. Kharzheev^{‡,13} B. Kilminster^{†dd,15} D.H. Kim^{†,25}
 H.S. Kim^{†,25} J.E. Kim^{†,25} M.J. Kim^{†,17} S.H. Kim^{†,50} S.B. Kim^{†,25} Y.J. Kim^{†,25} Y.K. Kim^{†,11} N. Kimura^{†,53}
 M. Kirby^{†,15} I. Kiselevich^{‡,34} K. Knoepfel^{†,15} J.M. Kohli^{‡,79} K. Kondo^{†,53,*} D.J. Kong^{†,25} J. Konigsberg^{†,16}
 A.V. Kotwal^{†,14} A.V. Kozelov^{‡,88} J. Kraus^{‡,108} M. Kreps^{†,24} J. Kroll^{†,41} M. Kruse^{†,14} T. Kuhr^{†,24}
 A. Kumar^{‡,112} A. Kupco^{‡,64} M. Kurata^{†,50} T. Kura^{‡,73} V.A. Kuzmin^{‡,87} A.T. Laasanen^{†,44} S. Lammel^{†,15}
 S. Lammers^{‡,101} M. Lancaster^{†,28} K. Lannon^{†x,36} G. Latino^{†xx,42} P. Lebrun^{‡,73} H.S. Lee^{‡,83} H.S. Lee^{†,25}
 J.S. Lee^{†,25} S.W. Lee^{‡,104} W.M. Lee^{‡,97} X. Lei^{‡,95} J. Lellouch^{‡,70} S. Leo^{†,42} S. Leone^{†,42} J.D. Lewis^{†,15}
 D. Li^{‡,70} H. Li^{‡,122} L. Li^{‡,96} Q.Z. Li^{†,15} J.K. Lim^{†,83} A. Limosani^{†s,14} D. Lincoln^{†,15} J. Linnemann^{‡,33}
 V.V. Lipaev^{†,88} E. Lipeles^{†,41} R. Lipton^{†,15} A. Lister^{†a,18} H. Liu^{†,52} H. Liu^{†,120} Q. Liu^{†,44} T. Liu^{†,15} Y. Liu^{†,60}
 A. Lobodenko^{‡,89} S. Lockwitz^{†,56} A. Loginov^{†,56} M. Lokajicek^{†,64} R. Lopes de Sa^{†,113} D. Lucchesi^{†vv,40}
 A. Lucà^{†,17} J. Lueck^{†,24} P. Lujan^{†,26} P. Lukens^{†,15} R. Luna-Garcia^{†pp,84} G. Lungu^{†,46} A.L. Lyon^{†,15} J. Lys^{†,26}
 R. Lysak^{†d,12} A.K.A. Maciel^{‡,57} R. Madar^{†,75} R. Madrak^{†,15} P. Maestro^{†xx,42} R. Magaa-Villalba^{‡,84} S. Malik^{†,46}
 S. Malik^{†,109} V.L. Malyshev^{‡,13} G. Manca^{†b,27} A. Manousakis-Katsikakis^{†,3} J. Mansour^{‡,76} L. Marchese^{†ii,6}
 F. Margaroli^{†,47} P. Marino^{†yy,42} J. Martínez-Ortega^{‡,84} M. Martínez^{†,4} K. Matera^{†,22} M.E. Mattson^{†,54}
 A. Mazzacane^{†,15} P. Mazzanti^{†,6} R. McCarthy^{‡,113} C.L. McGivern^{†,94} R. McNulty^{†i,27} A. Mehta^{†,27} P. Mehtala^{†,21}
 M.M. Meijer^{‡,85,86} A. Melnitchouk^{‡,15} D. Menezes^{‡,99} P.G. Mercadante^{‡,59} M. Merkin^{‡,87} C. Mesropian^{†,46}
 A. Meyer^{‡,74} J. Meyer^{†rr,76} T. Miao^{†,15} F. Miconi^{‡,72} D. Mietlicki^{†,32} A. Mitra^{†,1} H. Miyake^{†,50} S. Moed^{†,15}
 N. Moggi^{†,6} N.K. Mondal^{†,81} H.E. Montgomery^{†tt,15} C.S. Moon^{†z,15} R. Moore^{†eff,15} M.J. Morello^{†yy,42}
 A. Mukherjee^{†,15} M. Mulhearn^{†,122} Th. Muller^{†,24} P. Murat^{†,15} M. Mussini^{†uu,6} J. Nachtman^{†m,15} Y. Nagai^{†,50}
 J. Naganoma^{†,53} E. Nagy^{†,68} I. Nakanō^{†,37} A. Napier^{†,51} M. Narain^{†,118} R. Nayyar^{†,95} H.A. Neal^{‡,32} J.P. Negret^{†,61}
 J. Nett^{†,48} C. Neu^{†,52} P. Neustroev^{†,89} H.T. Nguyen^{†,122} T. Nigmanov^{†,43} L. Nodulman^{†,2} S.Y. Noh^{†,25}
 O. Norniella^{†,22} T. Nunnemann^{†,78} E. Nurse^{†,28} L. Oakes^{†,39} S.H. Oh^{†,14} Y.D. Oh^{†,25} I. Oksuzian^{†,52}
 T. Okusawa^{†,38} R. Orava^{†,21} J. Orduna^{†,121} L. Ortolan^{†,4} N. Osman^{†,68} J. Osta^{†,103} C. Pagliarone^{†,49} A. Pal^{†,119}
 E. Palencia^{†e,9} P. Palni^{†,35} V. Papadimitriou^{†,15} N. Parashar^{†,102} V. Parihar^{†,118} S.K. Park^{†,83} W. Parker^{†,55}
 R. Partridge^{†nn,118} N. Parua^{†,101} A. Patwa^{†ss,114} G. Pauletta^{†ccddd,49} M. Paulini^{†,10} C. Paus^{†,30} B. Penning^{†,15}
 M. Perfilov^{‡,87} Y. Peters^{‡,76} K. Petridis^{†,94} G. Petrillo^{‡,45} P. Pétroff^{‡,69} T.J. Phillips^{†,14} G. Piacentino^{†,42}
 E. Pianori^{†,41} J. Pilot^{†,7} K. Pitts^{†,22} C. Plager^{†,8} M.-A. Pleier^{†,114} V.M. Podstavkov^{‡,15} L. Pondrom^{†,55}
 A.V. Popov^{‡,88} S. Poprocki^{†f,15} K. Potamianos^{†,26} A. Pranko^{†,26} M. Prewitt^{‡,121} D. Price^{†,101} N. Prokopenko^{‡,88}
 F. Prokoshin^{†aa,13} F. Ptohos^{†g,17} G. Punzi^{†ww,42} J. Qian^{†,32} A. Quadri^{†,76} B. Quinn^{†,108} N. Ranjan^{†,44}
 P.N. Ratoff^{†,92} I. Razumov^{†,88} I. Redondo Fernández^{†,29} P. Renton^{†,39} M. Rescigno^{†,47} T. Riddick^{†,28}
 F. Rimondi^{†,6,*} I. Ripp-Baudot^{‡,72} L. Ristori^{†,42,15} F. Rizatdinova^{†,117} A. Robson^{†,19} T. Rodriguez^{†,41} S. Rolli^{†h,51}
 M. Rominsky^{†,15} M. Ronzani^{†ww,42} R. Roser^{†,15} J.L. Rosner^{†,11} A. Ross^{†,92} C. Royon^{†,71} P. Rubinov^{†,15}
 R. Ruchti^{†,103} F. Ruffini^{†xx,42} A. Ruiz^{†,9} J. Russ^{†,10} V. Rusu^{†,15} G. Sajot^{†,67} W.K. Sakumoto^{†,45} Y. Sakurai^{†,53}
 A. Sánchez-Hernández^{†,84} M.P. Sanders^{†,78} L. Santi^{†ccddd,49} A.S. Santos^{†qq,57} K. Sato^{†,50} G. Savage^{†,15}
 V. Saveliev^{†v,15} A. Savoy-Navarro^{†z,15} L. Sawyer^{†,106} T. Scanlon^{†,93} R.D. Schamberger^{†,113} Y. Scheglov^{†,89}
 H. Schellman^{†,100} P. Schlabach^{†,15} E.E. Schmidt^{†,15} C. Schwanenberger^{†,94} T. Schwarz^{†,32} R. Schwienhorst^{†,33}
 L. Scodellaro^{†,9} F. Scuri^{†,42} S. Seidel^{†,35} Y. Seiya^{†,38} J. Sekaric^{†,105} A. Semenov^{†,13} H. Severini^{†,116} F. Sforza^{†ww,42}
 E. Shabalina^{†,76} S.Z. Shalhout^{†,7} V. Shary^{†,71} S. Shaw^{†,33} A.A. Shchukin^{†,88} T. Shears^{†,27} R. Shekhar^{†,14}
 P.F. Shepard^{†,43} M. Shimojima^{†u,50} M. Shochet^{†,11} V. Simak^{†,63} A. Simonenko^{†,13} P. Skubic^{†,116} P. Slattery^{†,45}
 K. Sliwa^{†,51} D. Smirnov^{†,103} J.R. Smith^{†,7} F.D. Snider^{†,15} G.R. Snow^{†,109} J. Snow^{†,115} S. Snyder^{†,114}
 S. Sldner-Rembold^{†,94} H. Song^{†,43} L. Sonnenschein^{†,74} V. Sorin^{†,4} K. Soustruznik^{†,62} R. St. Denis^{†,19}
 M. Stancari^{†,15} J. Stark^{†,67} O. Stelzer-Chilton^{†,15} D. Stentz^{†w,15} D.A. Stoyanova^{†,88} M. Strauss^{†,116} J. Strologas^{†,35}
 Y. Sudo^{†,50} A. Sukhanov^{†,15} I. Suslov^{†,13} L. Suter^{†,94} P. Svoisky^{†,116} K. Takemasa^{†,50} Y. Takeuchi^{†,50} J. Tang^{†,11}
 M. Tecchio^{†,32} I. Shreyber-Tecker^{†,34} P.K. Teng^{†,1} J. Thom^{†f,15} D. S. Thompson^{†,11} E. Thomson^{†,41} V. Thukral^{†,48}
 M. Titov^{†,71} D. Toback^{†,48} S. Tokar^{†,12} V.V. Tokmenin^{†,13} K. Tollefson^{†,33} T. Tomura^{†,50} D. Tonelli^{†e,15}
 S. Torre^{†,17} D. Torretta^{†,15} P. Totaro^{†,40} M. Trovato^{†yy,42} Y.-T. Tsai^{†,45} D. Tsybychev^{†,113} B. Tuchming^{†,71}
 C. Tully^{†,111} F. Ukegawa^{†,50} S. Uozumi^{†,25} L. Uvarov^{†,89} S. Uvarov^{†,89} S. Uzunyan^{†,99} R. Van Kooten^{†,101}
 W.M. van Leeuwen^{†,85} N. Varelas^{†,98} E.W. Varnes^{†,95} I.A. Vasilyev^{†,88} F. Vázquez^{†l,16} G. Velev^{†,15}
 C. Vellidis^{†,15} A.Y. Verkheev^{†,13} C. Vernieri^{†yy,42} L.S. Vertogradov^{†,13} M. Verzocchi^{†,15} M. Vesterinen^{†,94}
 M. Vidal^{†,44} D. Vilanova^{†,71} R. Vilar^{†,9} J. Vizán^{†cc,9} M. Vogel^{†,35} P. Vokac^{†,63} G. Volpi^{†,17} P. Wagner^{†,41}
 H.D. Wahl^{†,97} R. Wallny^{†j,15} M.H.L.S. Wang^{†,15} S.M. Wang^{†,1} J. Warchol^{†,103} D. Waters^{†,28} G. Watts^{†,123}
 M. Wayne^{†,103} J. Weichert^{†,77} L. Welty-Rieger^{†,100} W.C. Wester III^{†,15} D. Whiteson^{†c,41} A.B. Wicklund^{†,2}
 S. Wilbur^{†,7} H.H. Williams^{†,41} M.R.J. Williams^{†,101} G.W. Wilson^{†,105} J.S. Wilson^{†,32} P. Wilson^{†,15}

B.L. Winer^{†,36} P. Wittich^{†f,15} M. Wobisch^{‡,106} S. Wolbers^{†,15} H. Wolfe^{†,36} D.R. Wood^{‡,107} T. Wright^{†,32}
 X. Wu^{†,18} Z. Wu^{†,5} T.R. Wyatt^{‡,94} Y. Xie^{‡,15} S. Yacoob^{‡,100} R. Yamada^{‡,15} K. Yamamoto^{†,38} D. Yamato^{†,38}
 S. Yang^{‡,60} T. Yang^{†,15} U.K. Yang^{†,25} Y.C. Yang^{†,25} W.-M. Yao^{†,26} T. Yasuda^{‡,15} Y.A. Yatsunenko^{‡,13}
 W. Ye^{‡,113} Z. Ye^{‡,15} G.P. Yeh^{†,15} K. Yi^{†m,15} H. Yin^{‡,15} K. Yip^{‡,114} J. Yoh^{†,15} K. Yorita^{†,53} T. Yoshida^{†k,38}
 S.W. Youn^{‡,15} G.B. Yu^{†,14} I. Yu^{†,25} J.M. Yu^{‡,32} A.M. Zanetti^{†,49} Y. Zeng^{†,14} J. Zennamo^{‡,112} T.G. Zhao^{‡,94}
 B. Zhou^{‡,32} C. Zhou^{†,14} J. Zhu^{†,32} M. Zielinski^{‡,45} D. Zieminska^{‡,101} L. Zivkovic^{‡,70} and S. Zucchelli^{†uu6}

(CDF Collaboration)[†]

(D0 Collaboration)[‡]

¹*Institute of Physics, Academia Sinica, Taipei, Taiwan 11529, Republic of China*

²*Argonne National Laboratory, Argonne, Illinois 60439, USA*

³*University of Athens, 157 71 Athens, Greece*

⁴*Institut de Fisica d'Altes Energies, ICREA, Universitat Autònoma de Barcelona, E-08193, Bellaterra (Barcelona), Spain*

⁵*Baylor University, Waco, Texas 76798, USA*

⁶*Istituto Nazionale di Fisica Nucleare Bologna, ^{uu}University of Bologna, I-40127 Bologna, Italy*

⁷*University of California, Davis, Davis, California 95616, USA*

⁸*University of California, Los Angeles, Los Angeles, California 90024, USA*

⁹*Instituto de Fisica de Cantabria, CSIC-University of Cantabria, 39005 Santander, Spain*

¹⁰*Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA*

¹¹*Enrico Fermi Institute, University of Chicago, Chicago, Illinois 60637, USA*

¹²*Comenius University, 842 48 Bratislava, Slovakia; Institute of Experimental Physics, 040 01 Kosice, Slovakia*

¹³*Joint Institute for Nuclear Research, RU-141980 Dubna, Russia*

¹⁴*Duke University, Durham, North Carolina 27708, USA*

¹⁵*Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA*

¹⁶*University of Florida, Gainesville, Florida 32611, USA*

¹⁷*Laboratori Nazionali di Frascati, Istituto Nazionale di Fisica Nucleare, I-00044 Frascati, Italy*

¹⁸*University of Geneva, CH-1211 Geneva 4, Switzerland*

¹⁹*Glasgow University, Glasgow G12 8QQ, United Kingdom*

²⁰*Harvard University, Cambridge, Massachusetts 02138, USA*

²¹*Division of High Energy Physics, Department of Physics, University of Helsinki,*

FIN-00014, Helsinki, Finland; Helsinki Institute of Physics, FIN-00014, Helsinki, Finland

²²*University of Illinois, Urbana, Illinois 61801, USA*

²³*The Johns Hopkins University, Baltimore, Maryland 21218, USA*

²⁴*Institut für Experimentelle Kernphysik, Karlsruhe Institute of Technology, D-76131 Karlsruhe, Germany*

²⁵*Center for High Energy Physics: Kyungpook National University,*

Daegu 702-701, Korea; Seoul National University, Seoul 151-742,

Korea; Sungkyunkwan University, Suwon 440-746,

Korea; Korea Institute of Science and Technology Information,

Daejeon 305-806, Korea; Chonnam National University,

Gwangju 500-757, Korea; Chonbuk National University, Jeonju 561-756,

Korea; Ewha Womans University, Seoul, 120-750, Korea

²⁶*Ernest Orlando Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

²⁷*University of Liverpool, Liverpool L69 7ZE, United Kingdom*

²⁸*University College London, London WC1E 6BT, United Kingdom*

²⁹*Centro de Investigaciones Energeticas Medioambientales y Tecnológicas, E-28040 Madrid, Spain*

³⁰*Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA*

³¹*Institute of Particle Physics: McGill University, Montréal, Québec H3A 2T8,*

Canada; Simon Fraser University, Burnaby, British Columbia V5A 1S6,

Canada; University of Toronto, Toronto, Ontario M5S 1A7,

Canada; TRIUMF, Vancouver, British Columbia V6T 2A3, Canada

³²*University of Michigan, Ann Arbor, Michigan 48109, USA*

³³*Michigan State University, East Lansing, Michigan 48824, USA*

³⁴*Institution for Theoretical and Experimental Physics, ITEP, Moscow 117259, Russia*

³⁵*University of New Mexico, Albuquerque, New Mexico 87131, USA*

³⁶*The Ohio State University, Columbus, Ohio 43210, USA*

³⁷*Okayama University, Okayama 700-8530, Japan*

³⁸*Osaka City University, Osaka 558-8585, Japan*

³⁹*University of Oxford, Oxford OX1 3RH, United Kingdom*

⁴⁰*Istituto Nazionale di Fisica Nucleare, Sezione di Padova, ^{vv}University of Padova, I-35131 Padova, Italy*

⁴¹*University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*

⁴²*Istituto Nazionale di Fisica Nucleare Pisa, ^{ww}University of Pisa,*

^{xx}University of Siena, ^{yy}Scuola Normale Superiore,

- I-56127 Pisa, Italy, ^{zz} INFN Pavia, I-27100 Pavia, Italy, ^{aaa} University of Pavia, I-27100 Pavia, Italy
- ⁴³ University of Pittsburgh, Pittsburgh, Pennsylvania 15260, USA
- ⁴⁴ Purdue University, West Lafayette, Indiana 47907, USA
- ⁴⁵ University of Rochester, Rochester, New York 14627, USA
- ⁴⁶ The Rockefeller University, New York, New York 10065, USA
- ⁴⁷ Istituto Nazionale di Fisica Nucleare, Sezione di Roma 1, ^{bbb} Sapienza Università di Roma, I-00185 Roma, Italy
- ⁴⁸ Mitchell Institute for Fundamental Physics and Astronomy, Texas A&M University, College Station, Texas 77843, USA
- ⁴⁹ Istituto Nazionale di Fisica Nucleare Trieste, ^{ccc} Gruppo Collegato di Udine, ^{ddd} University of Udine, I-33100 Udine, Italy, ^{eee} University of Trieste, I-34127 Trieste, Italy
- ⁵⁰ University of Tsukuba, Tsukuba, Ibaraki 305, Japan
- ⁵¹ Tufts University, Medford, Massachusetts 02155, USA
- ⁵² University of Virginia, Charlottesville, Virginia 22906, USA
- ⁵³ Waseda University, Tokyo 169, Japan
- ⁵⁴ Wayne State University, Detroit, Michigan 48201, USA
- ⁵⁵ University of Wisconsin, Madison, Wisconsin 53706, USA
- ⁵⁶ Yale University, New Haven, Connecticut 06520, USA
- ⁵⁷ LAFEX, Centro Brasileiro de Pesquisas Físicas, Rio de Janeiro, Brazil
- ⁵⁸ Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil
- ⁵⁹ Universidade Federal do ABC, Santo André, Brazil
- ⁶⁰ University of Science and Technology of China, Hefei, People's Republic of China
- ⁶¹ Universidad de los Andes, Bogotá, Colombia
- ⁶² Charles University, Faculty of Mathematics and Physics, Center for Particle Physics, Prague, Czech Republic
- ⁶³ Czech Technical University in Prague, Prague, Czech Republic
- ⁶⁴ Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
- ⁶⁵ Universidad San Francisco de Quito, Quito, Ecuador
- ⁶⁶ LPC, Université Blaise Pascal, CNRS/IN2P3, Clermont, France
- ⁶⁷ LPSC, Université Joseph Fourier Grenoble 1, CNRS/IN2P3, Institut National Polytechnique de Grenoble, Grenoble, France
- ⁶⁸ CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
- ⁶⁹ LAL, Université Paris-Sud, CNRS/IN2P3, Orsay, France
- ⁷⁰ LPNHE, Universités Paris VI and VII, CNRS/IN2P3, Paris, France
- ⁷¹ CEA, Irfu, SPP, Saclay, France
- ⁷² IPHC, Université de Strasbourg, CNRS/IN2P3, Strasbourg, France
- ⁷³ IPNL, Université Lyon 1, CNRS/IN2P3, Villeurbanne, France and Université de Lyon, Lyon, France
- ⁷⁴ III. Physikalisches Institut A, RWTH Aachen University, Aachen, Germany
- ⁷⁵ Physikalisches Institut, Universität Freiburg, Freiburg, Germany
- ⁷⁶ II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen, Germany
- ⁷⁷ Institut für Physik, Universität Mainz, Mainz, Germany
- ⁷⁸ Ludwig-Maximilians-Universität München, München, Germany
- ⁷⁹ Panjab University, Chandigarh, India
- ⁸⁰ Delhi University, Delhi, India
- ⁸¹ Tata Institute of Fundamental Research, Mumbai, India
- ⁸² University College Dublin, Dublin, Ireland
- ⁸³ Korea Detector Laboratory, Korea University, Seoul, Korea
- ⁸⁴ CINVESTAV, Mexico City, Mexico
- ⁸⁵ Nikhef, Science Park, Amsterdam, the Netherlands
- ⁸⁶ Radboud University Nijmegen, Nijmegen, the Netherlands
- ⁸⁷ Moscow State University, Moscow, Russia
- ⁸⁸ Institute for High Energy Physics, Protvino, Russia
- ⁸⁹ Petersburg Nuclear Physics Institute, St. Petersburg, Russia
- ⁹⁰ Institució Catalana de Recerca i Estudis Avançats (ICREA) and Institut de Física d'Altes Energies (IFAE), Barcelona, Spain
- ⁹¹ Uppsala University, Uppsala, Sweden
- ⁹² Lancaster University, Lancaster LA1 4YB, United Kingdom
- ⁹³ Imperial College London, London SW7 2AZ, United Kingdom
- ⁹⁴ The University of Manchester, Manchester M13 9PL, United Kingdom
- ⁹⁵ University of Arizona, Tucson, Arizona 85721, USA
- ⁹⁶ University of California Riverside, Riverside, California 92521, USA
- ⁹⁷ Florida State University, Tallahassee, Florida 32306, USA
- ⁹⁸ University of Illinois at Chicago, Chicago, Illinois 60607, USA
- ⁹⁹ Northern Illinois University, DeKalb, Illinois 60115, USA

- ¹⁰⁰Northwestern University, Evanston, Illinois 60208, USA
¹⁰¹Indiana University, Bloomington, Indiana 47405, USA
¹⁰²Purdue University Calumet, Hammond, Indiana 46323, USA
¹⁰³University of Notre Dame, Notre Dame, Indiana 46556, USA
¹⁰⁴Iowa State University, Ames, Iowa 50011, USA
¹⁰⁵University of Kansas, Lawrence, Kansas 66045, USA
¹⁰⁶Louisiana Tech University, Ruston, Louisiana 71272, USA
¹⁰⁷Northeastern University, Boston, Massachusetts 02115, USA
¹⁰⁸University of Mississippi, University, Mississippi 38677, USA
¹⁰⁹University of Nebraska, Lincoln, Nebraska 68588, USA
¹¹⁰Rutgers University, Piscataway, New Jersey 08855, USA
¹¹¹Princeton University, Princeton, New Jersey 08544, USA
¹¹²State University of New York, Buffalo, New York 14260, USA
¹¹³State University of New York, Stony Brook, New York 11794, USA
¹¹⁴Brookhaven National Laboratory, Upton, New York 11973, USA
¹¹⁵Langston University, Langston, Oklahoma 73050, USA
¹¹⁶University of Oklahoma, Norman, Oklahoma 73019, USA
¹¹⁷Oklahoma State University, Stillwater, Oklahoma 74078, USA
¹¹⁸Brown University, Providence, Rhode Island 02912, USA
¹¹⁹University of Texas, Arlington, Texas 76019, USA
¹²⁰Southern Methodist University, Dallas, Texas 75275, USA
¹²¹Rice University, Houston, Texas 77005, USA
¹²²University of Virginia, Charlottesville, Virginia 22904, USA
¹²³University of Washington, Seattle, Washington 98195, USA

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We summarize and combine direct measurements of the mass of the W boson in $\sqrt{s} = 1.96$ TeV proton-antiproton collision data collected by CDF and D0 experiments at the Fermilab Tevatron Collider. Earlier measurements from CDF and D0 are combined with the two latest, more precise measurements: a CDF measurement in the electron and muon channels using data corresponding to 2.2 fb^{-1} of integrated luminosity, and a D0 measurement in the electron channel using data corresponding to 4.3 fb^{-1} of integrated luminosity. The resulting Tevatron average for the mass of the W boson is $M_W = 80387 \pm 16$ MeV. Including measurements obtained in electron-positron collisions at LEP yields the most precise value of $M_W = 80385 \pm 15$ MeV.

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I. INTRODUCTION

In the standard model (SM), quantum corrections to the mass of the W boson (M_W) are dominated by con-

*Deceased

†With visitors from ^aUniversity of British Columbia, Vancouver, BC V6T 1Z1, Canada, ^bIstituto Nazionale di Fisica Nucleare, Sezione di Cagliari, 09042 Monserrato (Cagliari), Italy, ^cUniversity of California Irvine, Irvine, CA 92697, USA, ^dInstitute of Physics, Academy of Sciences of the Czech Republic, 182 21, Czech Republic, ^eCERN, CH-1211 Geneva, Switzerland, ^fCornell University, Ithaca, NY 14853, USA, ^gUniversity of Cyprus, Nicosia CY-1678, Cyprus, ^hOffice of Science, U.S. Department of Energy, Washington, DC 20585, USA, ⁱUniversity College Dublin, Dublin 4, Ireland, ^jETH, 8092 Zürich, Switzerland, ^kUniversity of Fukui, Fukui City, Fukui Prefecture, Japan 910-0017, ^lUniversidad Iberoamericana, Lomas de Santa Fe, México, C.P. 01219, Distrito Federal, ^mUniversity of Iowa, Iowa City, IA 52242, USA, ⁿKinki University, Higashi-Osaka City, Japan 577-8502, ^oKansas State University, Manhattan, KS 66506, USA, ^pBrookhaven National Laboratory, Upton, NY 11973, USA, ^qUniversity of Manchester, Manchester M13 9PL, United Kingdom, ^rQueen Mary, University of London, London, E1 4NS, United Kingdom, ^sUniversity of Melbourne, Victoria 3010, Australia, ^tMuons, Inc., Batavia, IL 60510, USA, ^uNagasaki Institute of Applied Science, Nagasaki 851-0193, Japan, ^vNational Research Nuclear University, Moscow 115409, Russia, ^wNorthwestern University, Evanston, IL 60208, USA, ^xUniversity of Notre Dame, Notre Dame, IN 46556, USA, ^yUniversidad de Oviedo, E-33007 Oviedo, Spain, ^zCNRS-IN2P3,

Paris, F-75205 France, ^{aa}Universidad Tecnica Federico Santa Maria, 110v Valparaiso, Chile, ^{bb}The University of Jordan, Amman 11942, Jordan, ^{cc}Universite catholique de Louvain, 1348 Louvain-La-Neuve, Belgium, ^{dd}University of Zürich, 8006 Zürich, Switzerland, ^{ee}Massachusetts General Hospital, Boston, MA 02114 USA, ^{ff}Harvard Medical School, Boston, MA 02114 USA, ^{gg}Hampton University, Hampton, VA 23668, USA, ^{hh}Los Alamos National Laboratory, Los Alamos, NM 87544, USA, ⁱⁱUniversità degli Studi di Napoli Federico I, I-80138 Napoli, Italy

[‡]With visitors from ^{jj}Augustana College, Sioux Falls, SD, USA, ^{kk}The University of Liverpool, Liverpool, UK, ^{ll}DESY, Hamburg, Germany, ^{mm}Universidad Michoacana de San Nicolas de Hidalgo, Morelia, Mexico, ⁿⁿSLAC, Menlo Park, CA, USA, ^{oo}University College London, London, UK, ^{pp}Centro de Investigacion en Computacion - IPN, Mexico City, Mexico, ^{qq}Universidade Estadual Paulista, So Paulo, Brazil, ^{rr}Karlsruher Institut für Technologie (KIT) - Steinbuch Centre for Computing (SCC), ^{ss}Office of Science, U.S. Department of Energy, Washington, D.C. 20585, USA, ^{tt}Thomas Jefferson National Accelerator Facility, Newport News, VA 23606, USA

tributions dependent on the mass of the top quark (m_t), the mass of the Higgs boson (M_H), and the fine-structure constant α . A precise measurement of M_W and m_t therefore constrains M_H . Comparing this constraint with the mass of the Higgs boson recently discovered at the LHC [1] is a critical test of its nature and the consistency of the SM. Details of the experimental methods used in measurements of M_W are discussed in Ref. [2]. Prior to the combination reported here, the uncertainty on the world average M_W was 23 MeV [3, 4]. Direct measurements of m_t at the Fermilab Tevatron collider have a combined uncertainty of 0.94 GeV [5], and the uncertainty on M_W would have to be 6 MeV [6] to provide equally constraining information on M_H . The experimental precision on the measured M_W is therefore currently the limiting factor on the constraints.

The CDF and D0 experiments at the Fermilab Tevatron proton-antiproton collider reported several direct measurements of the natural width [7] and mass [8–18] of the W boson, using the $e\nu_e$ and $\mu\nu_\mu$ decay modes of the W boson. Measurements of M_W have been reported by CDF with data sets collected during 1988–1989 [8], 1992–1993 [9], 1994–1995 [10], and 2001–2004 [11] and by D0 using data taken during 1992–1995 [12–15] and 2002–2006 [16].

This article describes a combination of M_W measurements including recent measurements from CDF using the 2002–2007 dataset [17] and D0 using the 2006–2009 dataset [18] denoted below as CDF (2012) and D0 (2012), respectively. The recent CDF (2012) measurement supersedes the previous measurement [11], which was based on an integrated luminosity of 200 pb⁻¹ and was used in previous combinations [3, 19]. The combination takes into account the statistical and systematic uncertainties as well as correlations among systematic uncertainties and supersedes the previous combinations [3, 19, 20]. All the combinations presented in this article are done using the best linear unbiased estimator (BLUE) method [21], which prescribes the construction of a covariance matrix from partially correlated measurements.

II. W -BOSON MASS MEASUREMENT STRATEGY AT THE TEVATRON

At the Tevatron, W bosons are primarily produced in quark-antiquark annihilation, $q\bar{q}' \rightarrow W + X$, where X can include QCD radiation, such as initial-state gluon radiation, that results in measurable hadronic-recoil energy. The W -boson mass is measured using low-background samples of $W \rightarrow \ell\nu_\ell$ decays ($\ell = e, \mu$ at CDF and $\ell = e$ at D0) that are reconstructed using the CDF [22] and D0 [23] detectors. The mass is determined using three kinematic variables measured in the plane perpendicular to the beam direction: the transverse momentum of the charged lepton (p_T^ℓ), the transverse momentum of the neutrino (p_T^ν), and the transverse mass $m_T^\ell = \sqrt{2p_T^\ell p_T^\nu (1 - \cos \Delta\phi)}$, where $\Delta\phi$ is the open-

ing angle between the lepton and neutrino momenta in the plane transverse to the beam. The magnitude and direction of p_T^ν is inferred from the vector of the missing transverse energy \cancel{E}_T^ℓ [24]. The W -boson mass is extracted from maximum-likelihood fits to the binned distributions of the observed p_T^ℓ , \cancel{E}_T^ℓ , and m_T^ℓ values using a parametrized simulation of these distributions as a function of M_W . These simulations depend on the kinematic distributions of the W -boson decay products and also on detector effects that are constrained using theoretical calculations and control samples. The kinematic distributions are determined by several effects including the W -boson transverse momentum $p_T(W)$ and the parton distribution functions (PDFs) of the interacting protons and antiprotons. Major detector effects include energy response to leptons, hadronic recoil, the response to QED radiation, and multiple-interaction pileup, together with calorimeter acceptance effects and lepton-identification efficiencies. The detailed simulations developed at CDF and D0 enable the study of these effects to better than 1 part in 10⁴ precision on the observed value of M_W .

In the CDF (2012) and D0 (2012) measurements, the kinematic properties of W -boson production and decay are simulated using RESBOS [25], which is a next-to-leading order generator that includes next-to-next-to-leading logarithm resummation of soft gluons at low boson p_T [26]. The momenta of interacting partons in RESBOS are calculated as a fractions of the colliding (anti)proton momenta using the CTEQ6.6 [27] PDFs. The radiation of photons from final-state leptons is simulated using PHOTOS [28].

III. CDF (2012) AND D0 (2012) MEASUREMENTS

A. CDF Measurement

The CDF (2012) measurement uses data corresponding to an integrated luminosity of 2.2 fb⁻¹, collected between 2002 and 2007. Both the muon ($W \rightarrow \mu\nu_\mu$) and electron ($W \rightarrow e\nu_e$) channels are considered. Decays of J/ψ and Υ mesons into muon pairs are reconstructed in a central tracking system to establish the absolute momentum scale. A measurement of the Z -boson mass (M_Z) in $Z \rightarrow \mu\mu$ decays is performed as a consistency check. This measurement, which uses the tracking detector, yields $M_Z = 91\,180 \pm 12$ (stat) ± 10 (syst) MeV, consistent with the world average mass of $91\,188 \pm 2$ MeV [29], and is therefore also used as an additional constraint on the momentum scale. The electromagnetic calorimeter energy scale and nonlinearity are determined by fitting the peak of the E/p distribution of electrons from $W \rightarrow e\nu$ and $Z \rightarrow ee$ decays, where E is the energy measured in the calorimeter and p is the momentum of the associated charged particle. The lower tail of the E/p distribution is used to determine the amount of material in the tracking detector. The Z -boson mass

Source	Uncertainty (MeV)
Lepton energy scale and resolution	7
Recoil energy scale and resolution	6
Lepton removal from recoil	2
Backgrounds	3
Experimental subtotal	10
Parton distribution functions	10
QED radiation	4
$p_T(W)$ model	5
Production subtotal	12
Total systematic uncertainty	15
W -boson event yield	12
Total uncertainty	19

TABLE I: Uncertainties of the CDF (2012) M_W measurement determined from the combination of the six measurements.

measured in $Z \rightarrow ee$ decays is used as a consistency check and to constrain the energy scale. The value of $M_Z = 91\,230 \pm 30$ (stat) ± 14 (syst) MeV from the calorimetric measurement is also consistent with the world average.

The CDF (2012) measurement of M_W is obtained from the combination of six observables: p_T^μ , \cancel{E}_T^μ , m_T^μ , p_T^e , \cancel{E}_T^e and m_T^e . The combined result is $M_W = 80\,387 \pm 12$ (stat) ± 15 (syst) MeV. Table I summarizes the sources of uncertainty in the CDF measurement.

B. D0 Measurement

The D0 (2012) measurement uses data corresponding to 4.3 fb^{-1} of integrated luminosity recorded between 2006 and 2009. D0 calibrates the calorimeter energy scale using $Z \rightarrow ee$ decays. Corrections for energy lost in un-instrumented regions are based on a comparison between the shower-development profiles from data and from a detailed GEANT-based simulation [30] of the D0 detector. The world average value for M_Z [29] is used to determine the absolute energy-scale of the calorimeter, which is thereafter used to correct the measurement of the electron energy from the W -boson decay. This M_W measurement is therefore equivalent to a measurement of the ratio of W - and Z -boson masses. This calibration method eliminates many systematic uncertainties common to the W - and Z -boson mass measurements, but its precision is limited by the size of the available Z -boson data set.

The results obtained with the two most sensitive observables m_T^e and p_T^e are combined to determine the W -boson mass of $M_W = 80\,367 \pm 13$ (stat) ± 22 (syst) MeV. A summary of the uncertainties is presented in Table II. This D0 (2012) measurement is combined with a previous D0 measurement [16] corresponding to an integrated luminosity of 1.0 fb^{-1} , which uses data recorded between 2002 and 2006, to yield $M_W = 80\,375 \pm 11$ (stat) ± 20 (syst) MeV.

Source	Uncertainty (MeV)
Electron energy calibration	16
Electron resolution model	2
Electron shower modeling	4
Electron energy loss model	4
Recoil energy scale and resolution	5
Electron efficiencies	2
Backgrounds	2
Experimental subtotal	18
Parton distribution functions	11
QED radiation	7
$p_T(W)$ model	2
Production subtotal	13
Total systematic uncertainty	22
W -boson event yield	13
Total uncertainty	26

TABLE II: Uncertainties of the D0 (2012) M_W measurement determined from the combination of the two most sensitive observables m_T^e and p_T^e .

IV. COMBINATION WITH PREVIOUS TEVATRON MEASUREMENTS

The CDF measurements from Ref. [8] (1988-1989) and Ref. [9] (1992-1993) were made using superseded PDF sets and have been corrected [19] using recent PDF sets. The previous results are also adjusted to use the same combination technique (the BLUE method) as in later combinations. The templates for fitting M_W assume the Breit-Wigner running-width scheme propagator, $1/(\hat{s} - M_W^2 + i\hat{s}\Gamma_W/M_W)$, which makes the value of M_W determined by the fit dependent on Γ_W . Here, \hat{s} is the square of the center-of-mass energy in the parton reference frame and Γ_W is the total width of the W boson. Different measurements have used different values of Γ_W , yielding a shift in measured values of the W -boson mass [19], $\Delta M_W = -(0.15 \pm 0.05) \Delta\Gamma_W$, where $\Delta\Gamma_W$ is the difference between the value of Γ_W predicted by the SM, $\Gamma_W = 2\,092.2 \pm 1.5$ MeV [31], and that used in a particular analysis. The prediction of Γ_W assumes $M_W = 80\,385 \pm 15$ MeV, which is a preliminary world-average combination result [32] of this article. The impact of the corrections on the final M_W combination reported in this article is found to be less than 0.2 MeV. Table III summarizes all inputs to the combination and the corrections made to ensure consistency across measurements.

V. CORRELATIONS IN THE CDF AND D0 M_W MEASUREMENTS

The increased statistical power of CDF (2012) and D0 (2012) M_W measurements necessitates a more detailed treatment of the systematic uncertainties due to the W -boson production and decay model that are independent

	CDF [8] 4.4 pb ⁻¹ (1988-1989)	CDF [9] 18.2 pb ⁻¹ (1992-1993)	CDF [10] 84 pb ⁻¹ (1994-1995)	D0 [12-15] 95 pb ⁻¹ (1992-1995)	D0 [16] 1.0 fb ⁻¹ (2002-2006)	CDF [17] 2.2 fb ⁻¹ (2002-2007)	D0 [18] 4.3 fb ⁻¹ (2006-2009)
Mass and width							
M_W	79 910	80 410	80 470	80 483	80 400	80 387	80 367
Γ_W	2 100	2 064	2 096	2 062	2 099	2 094	2 100
M_W uncertainties							
PDF	60	50	15	8	10	10	11
Radiative corrections	10	20	5	12	7	4	7
Γ_W	0.5	1.4	0.3	1.5	0.4	0.2	0.5
Total	390	181	89	84	43	19	26
M_W corrections							
$\Delta\Gamma_W$	+1.2	-4.2	+0.6	-4.5	+1.1	+0.3	+1.2
PDF	+20	-25	0	0	0	0	0
Fit method	-3.5	-3.5	-0.1	0	0	0	0
Total	+17.7	-32.7	+0.5	-4.5	+1.1	+0.3	+1.2
M_W corrected	79 927.7	80 377.3	80 470.5	80 478.5	80 401.8	80 387.3	80 368.6

TABLE III: The input data used in the M_W combination. All entries are in units of MeV.

of the data-sample size. We assume that for each uncertainty category, the smallest uncertainty across measurements is fully correlated while excesses above that level are generally assumed to be due to uncorrelated differences between measurements. One exception corresponds to the two D0 measurements that use very similar models and are treated as fully correlated [16, 18].

The experimental systematic uncertainties of the D0 measurement are dominated by the uncertainty in the energy scale for electrons and are nearly purely of statistical origin, as they are derived from the limited sample of $Z \rightarrow ee$ decays. CDF uses independent data from the central tracker to set the muon and electron energy scales. Thus, we assume no correlations between the experimental uncertainties of CDF and D0, or between independent measurements by either experiment.

Three sources of systematic uncertainty due to modeling of the production and decay of W and Z bosons are assumed to be at least partially correlated across all Tevatron measurements: (1) the choice of PDF sets, (2) the assumed Γ_W value, and (3) the electroweak radiative corrections.

A. PDF sets

Both experiments use the CTEQ6.6 [27] PDF set in their W -boson production model. D0 uses the CTEQ6.1 [33] uncertainty set to estimate the PDF uncertainties, while CDF uses MSTW2008 [34] and checks consistency with the CTEQ6.6 uncertainty set. Since these PDF sets are similar and rely on common inputs, the uncertainties introduced by PDFs in the recent measurements are assumed to be correlated and treated using the prescription for partial correlations described above.

B. Assumed Γ_W value

We assume that the small uncertainty due to Γ_W is fully correlated across all measurements.

C. QED radiative corrections

Current estimates of the uncertainties due to electroweak radiative corrections include a significant statistical component due to the size of the simulated data sets used in the uncertainty-propagation studies. The PHOTOS [28] radiative correction model is used in the recent measurements with consistency checks from $w(z)$ GRAD [35] and HORACE [36]. These studies yield model differences consistent within statistical uncertainties. We assume that uncertainties from purely theoretical sources, totaling 3.5 MeV, are correlated while remaining uncertainties, partially dependent on detector geometry, are uncorrelated.

VI. COMBINATION OF TEVATRON M_W MEASUREMENTS

The measurements of M_W obtained at Tevatron experiments included in this combination are given in Table III and include both the latest measurements [17, 18] discussed above, but exclude the superseded 0.2 fb⁻¹ CDF measurement [11]. Table IV shows the relative weight of each measurement in the combination. The combined value of the W -boson mass obtained from measurements performed at Tevatron experiments is

$$M_W = 80\,387 \pm 16 \text{ MeV}. \quad (1)$$

The χ^2 for the combination is 4.2 for 6 degrees of freedom, with a probability of 64%. The global correlation

Measurement	Relative weight in %
CDF [8]	0.1
CDF [9]	0.5
CDF [10]	1.9
D0 [12–15]	2.8
D0 [16]	7.9
CDF [17]	60.3
D0 [18]	26.5

TABLE IV: Relative weights of the contributions to the combined Tevatron measurement of M_W .

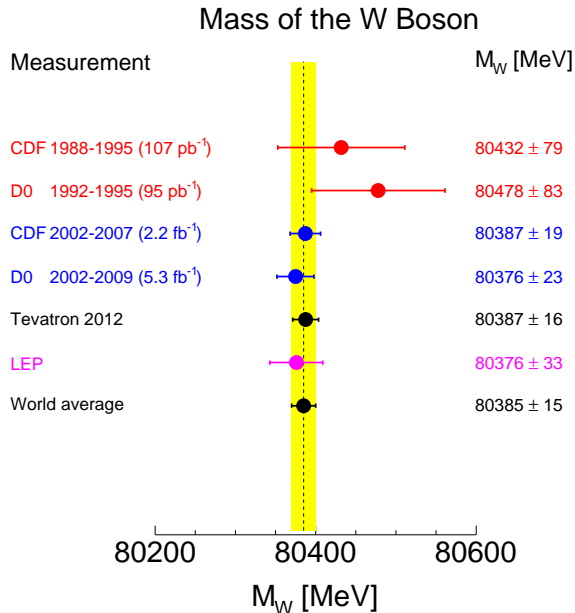


FIG. 1: W -boson mass determinations from the CDF and D0 Run I (1989 to 1996) and Run II (2001 to 2009) measurements, the new Tevatron average, the LEP combined result [29], and the world average obtained by combining the Tevatron and LEP averages assuming no correlations between them. The world-average uncertainty (15 MeV) is indicated by the shaded band.

matrix for the seven measurements is shown in Table V.

VII. WORLD AVERAGE

We also combine the Tevatron measurements with the value $M_W = 80376 \pm 33$ MeV determined from $e^+ e^- \rightarrow W^+ W^-$ production at LEP [29]. Assuming no correlations, this yields the currently most precise value of the W boson mass of

$$M_W = 80385 \pm 15 \text{ MeV.} \quad (2)$$

The combination of the seven statistically independent Tevatron measurements and the LEP measurement yields

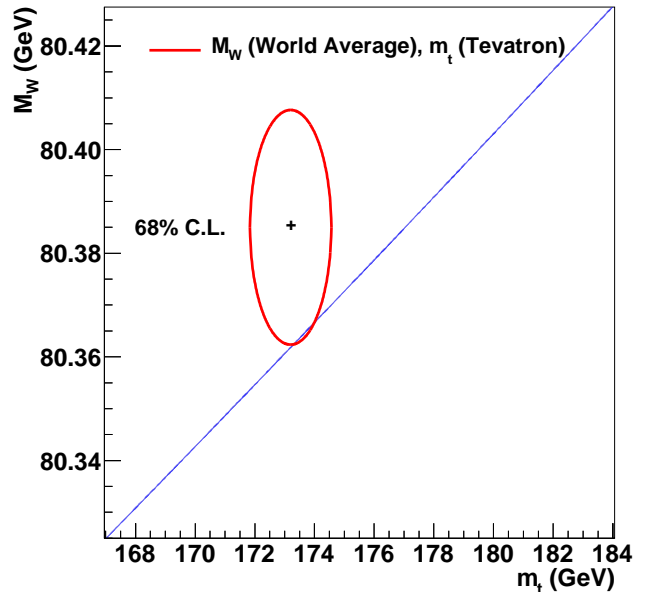


FIG. 2: The most recent world average of M_W is displayed along with the mass of the top quark m_t [5] at 68% C.L. by area. The diagonal line is the indirect prediction of M_W as a function of m_t , in the SM given by Ref. [6], assuming the measurements of the ATLAS and CMS [1] experiments of the candidate Higgs-boson masses of 126.0 GeV and 125.3 GeV respectively.

a χ^2 of 4.3 for 7 degrees of freedom with a probability of 74%. Figure 1 shows the individual measurements and the most recent combined world average of M_W .

VIII. SUMMARY

The latest high-precision measurements of M_W performed at the CDF and D0 experiments, combined with previous measurements by the Tevatron experiments, improve the uncertainty on the combined Tevatron M_W value to 16 MeV. The combination of this measurement with the LEP average for M_W further reduces the uncertainty to 15 MeV. The substantial improvement in the experimental precision on M_W leads to tightened indirect constraints on the mass of the SM Higgs boson. The direct measurements of the mass of the Higgs boson at the LHC [1] agree, at the level of 1.3 standard deviations, with these tightened indirect constraints [37]. This remarkable success of the standard model is also shown in Fig. 2, which includes the new world average W -boson mass, the Tevatron average top-quark mass measurement [5], and shows consistency among these with the calculation of M_W [6], assuming Higgs-boson mass determinations from the ATLAS and CMS experiments [1].

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	CDF [8]	CDF [9]	CDF [10]	D0 [12–15]	D0 [16]	CDF [17]	D0 [18]
CDF [8]	1	0.002	0.003	0.002	0.007	0.015	0.011
CDF [9]		1	0.007	0.005	0.014	0.033	0.024
CDF [10]			1	0.009	0.029	0.066	0.049
D0 [12–15]				1	0.019	0.044	0.032
D0 [16]					1	0.137	0.137
CDF [17]						1	0.230
D0 [18]							1

TABLE V: Correlation coefficients among measurements.

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gram and NRF (Korea); CONACyT (Mexico); FOM (The Netherlands); MON, NRC KI, and RFBR (Russia); the Slovak R&D Agency (Slovakia); the Ministerio de Ciencia e Innovación, and Programa Consolider-Ingenio 2010 (Spain); the Swedish Research Council (Sweden); SNSF (Switzerland); STFC and the Royal Society (United Kingdom); and the A. P. Sloan Foundation (USA).

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